

# Neutron transfer (d,p) reactions to inform neutron capture

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Spectroscopy & Related Topics

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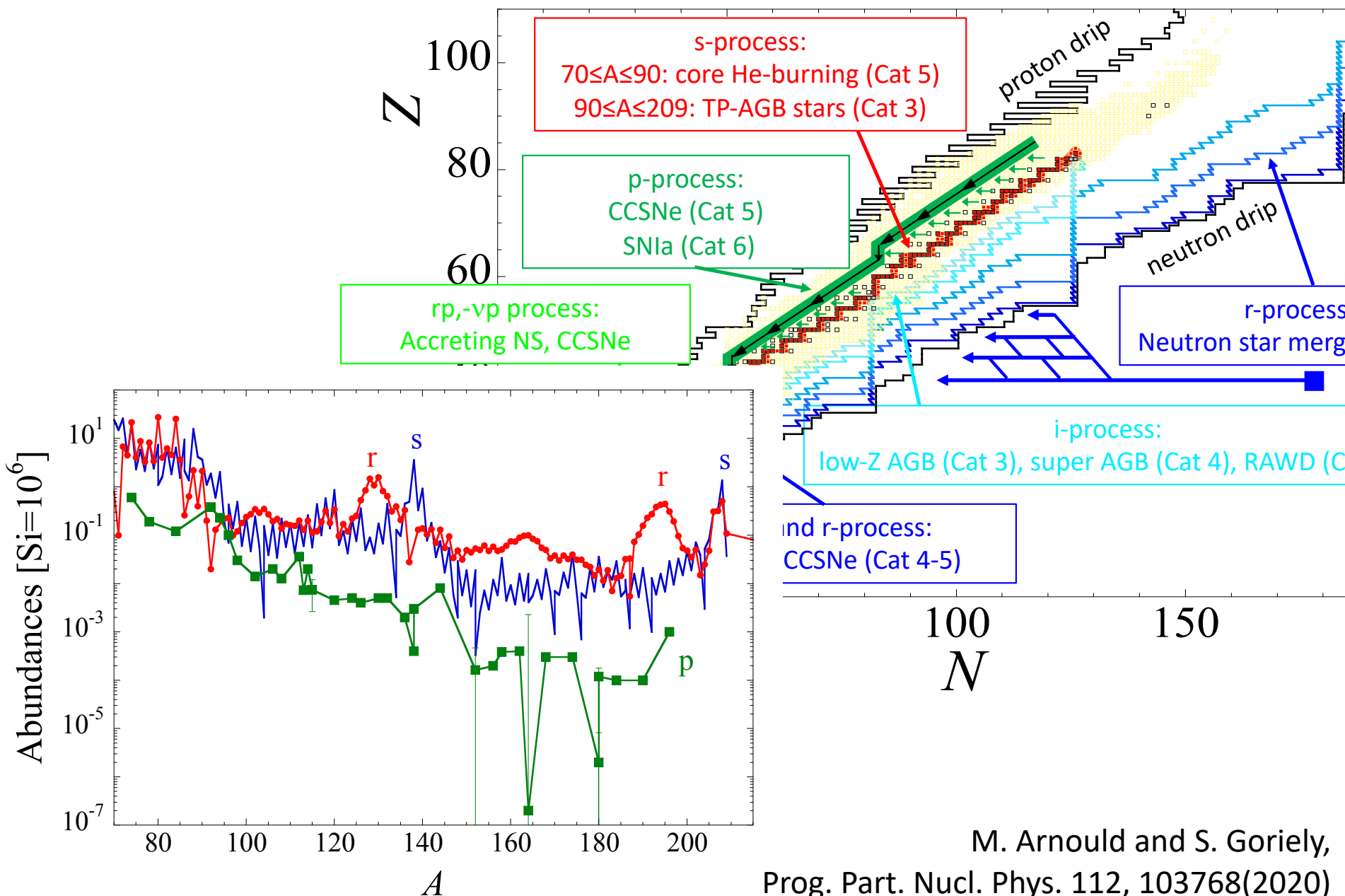
**Karl Smith**

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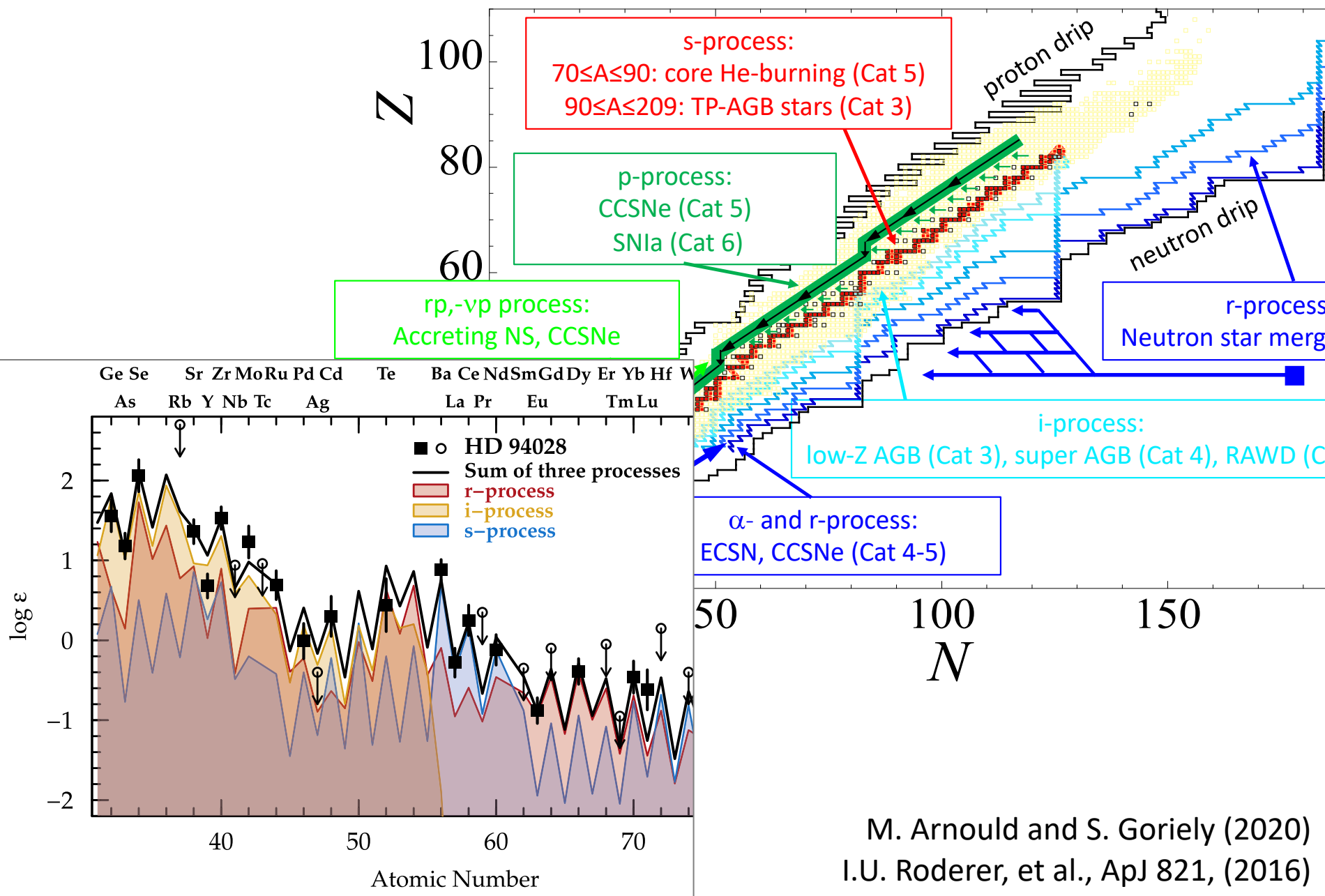


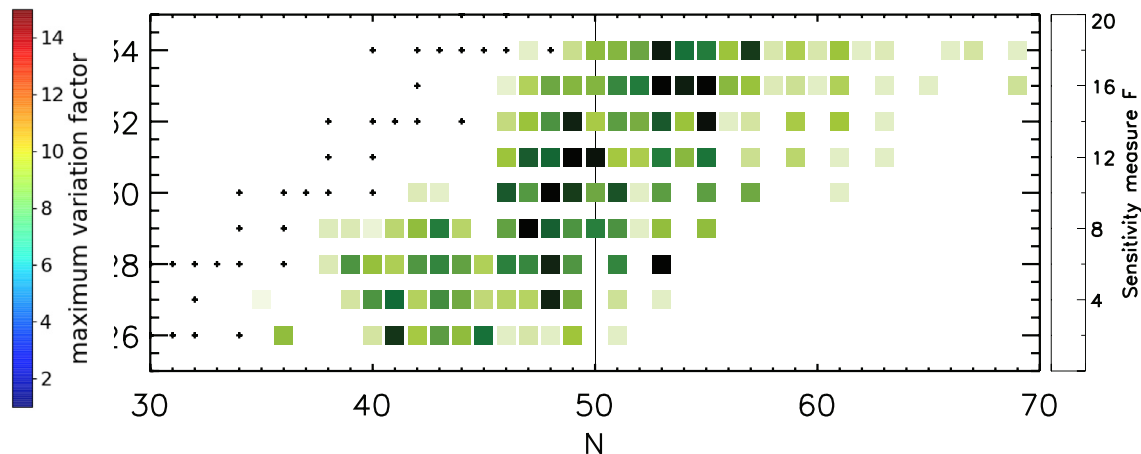
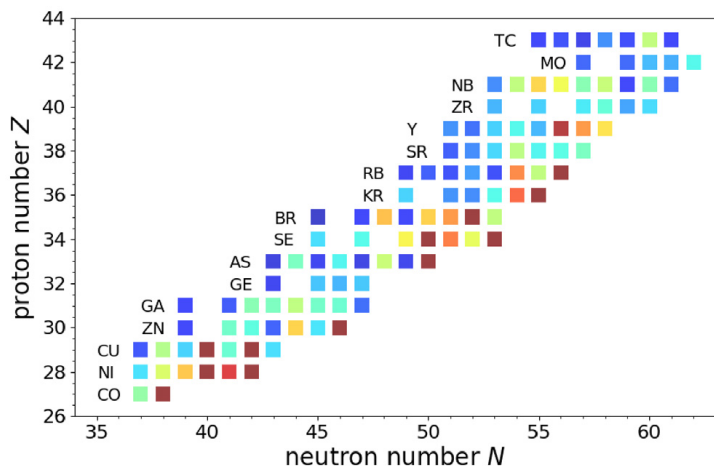
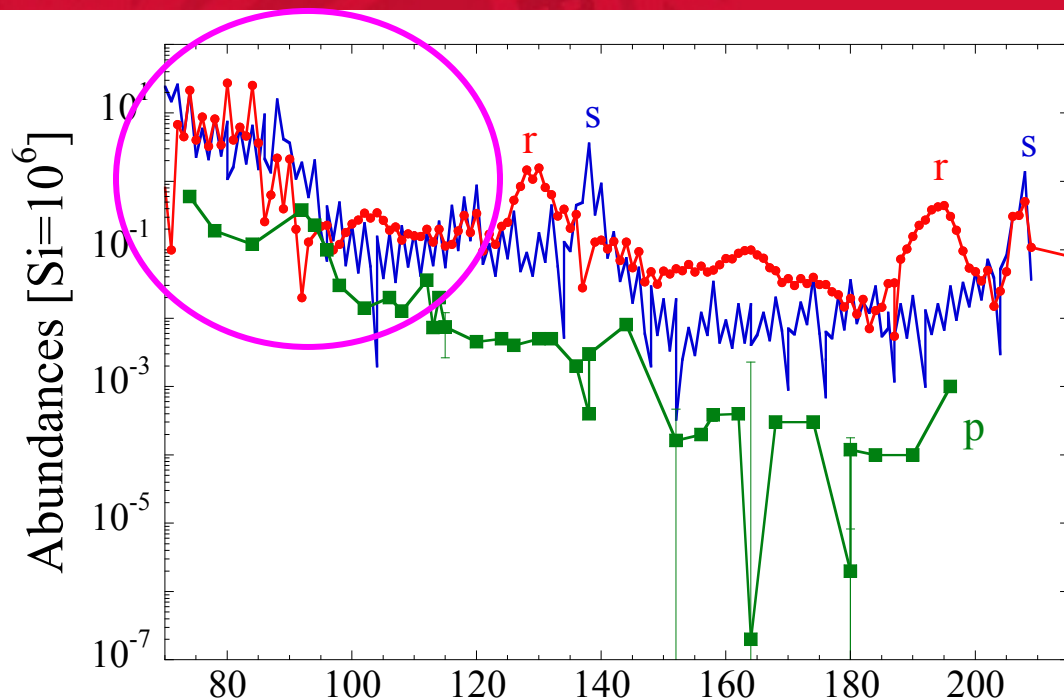






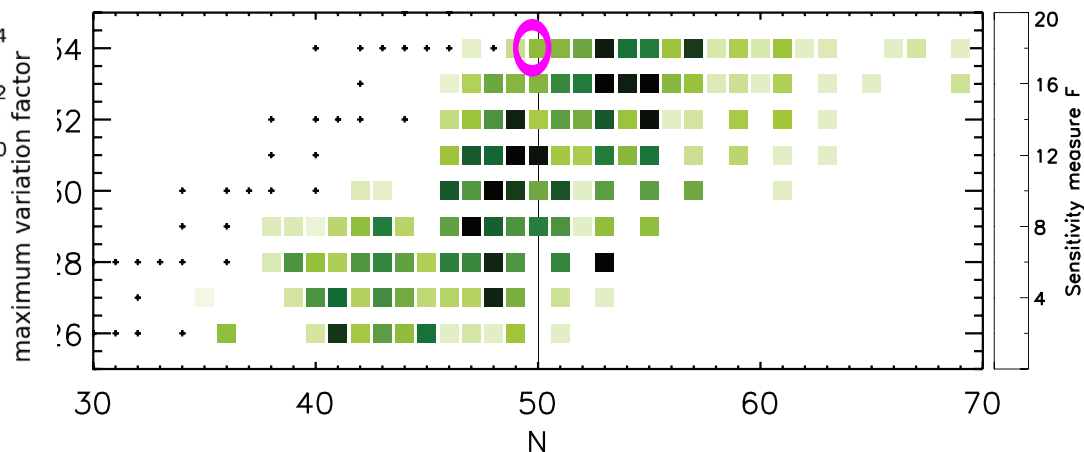
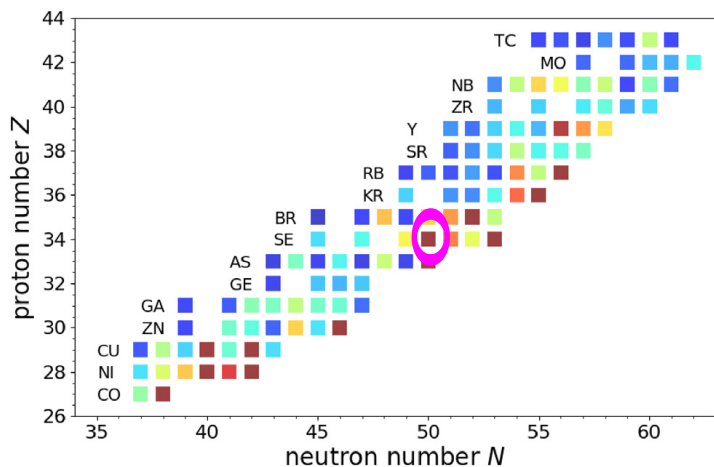
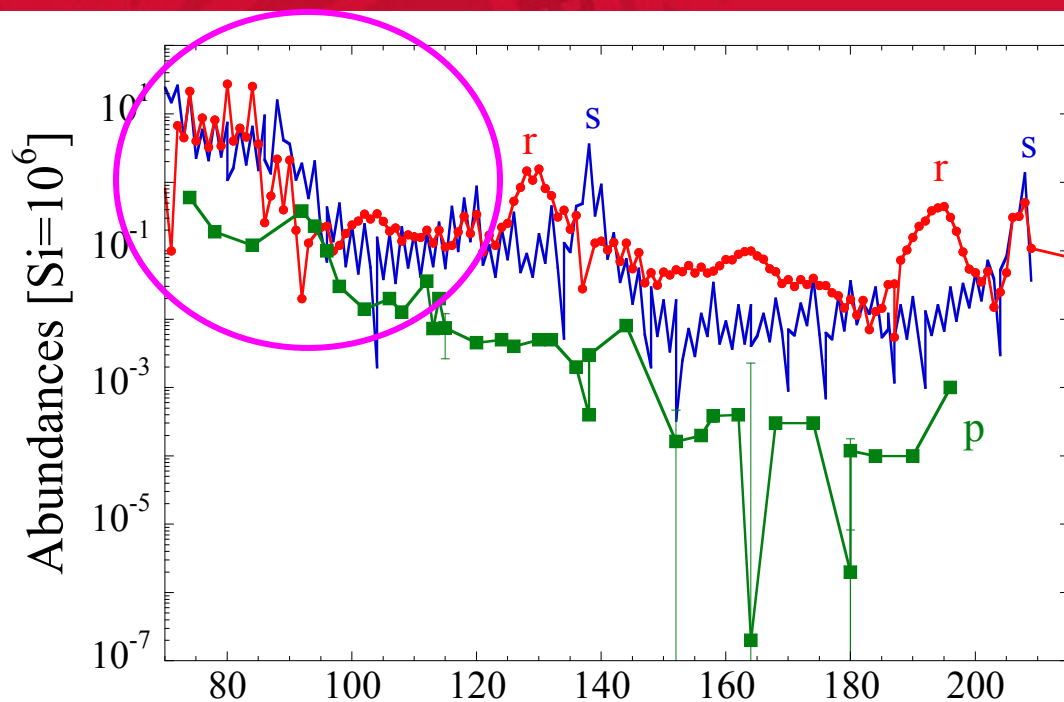
# Single (old) stars: many (n,γ) processes





McKay et al., (i process)  
MNRAS **491**, 5179–5187 (2020)

R. Surman et al., (weak-r process)  
AIP Advances **4**, 041008 (2014)

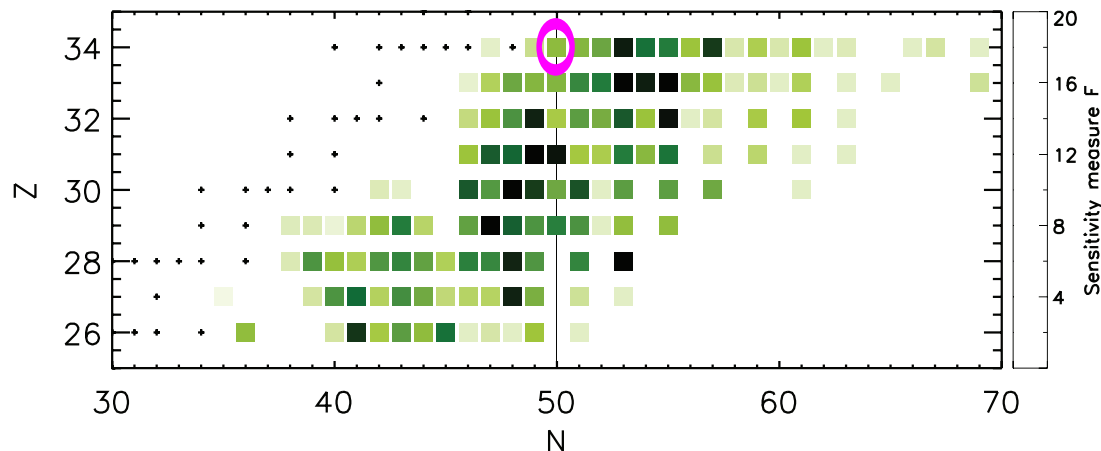
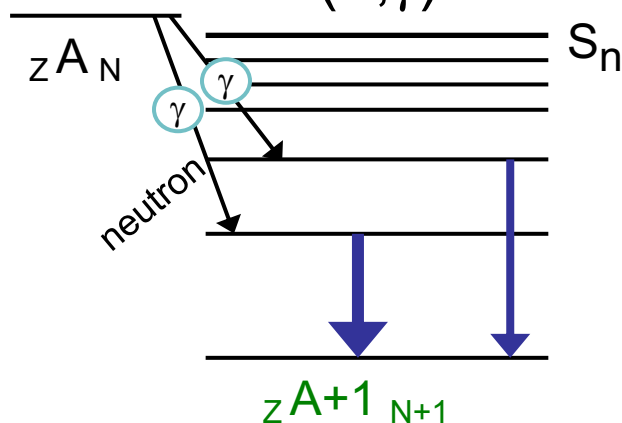


McKay et al., (i process)  
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R. Surman et al., (weak-r process)  
AIP Advances **4**, 041008 (2014)

## Focus: $N=50$ $^{84}\text{Se}$

Direct-semi-direct  
Near N shell closures  
(n, $\gamma$ )



### Near closed shells

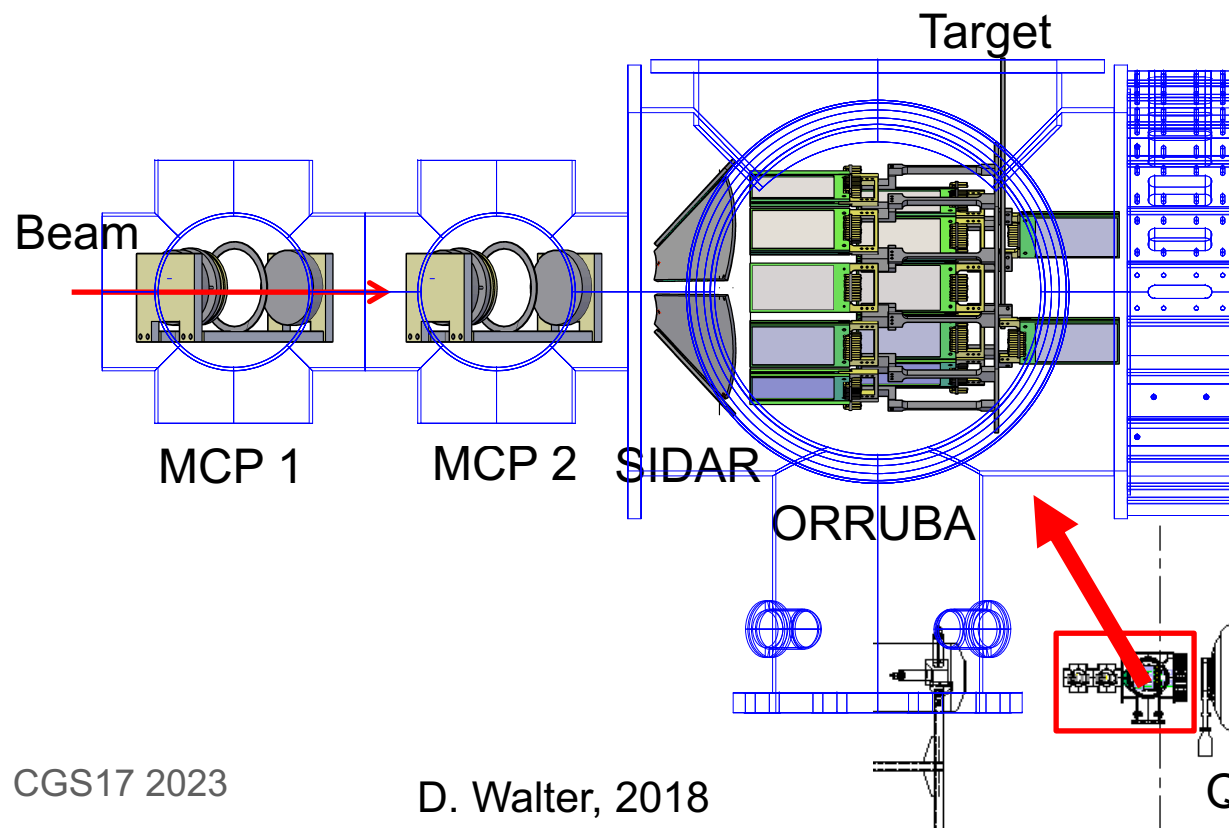
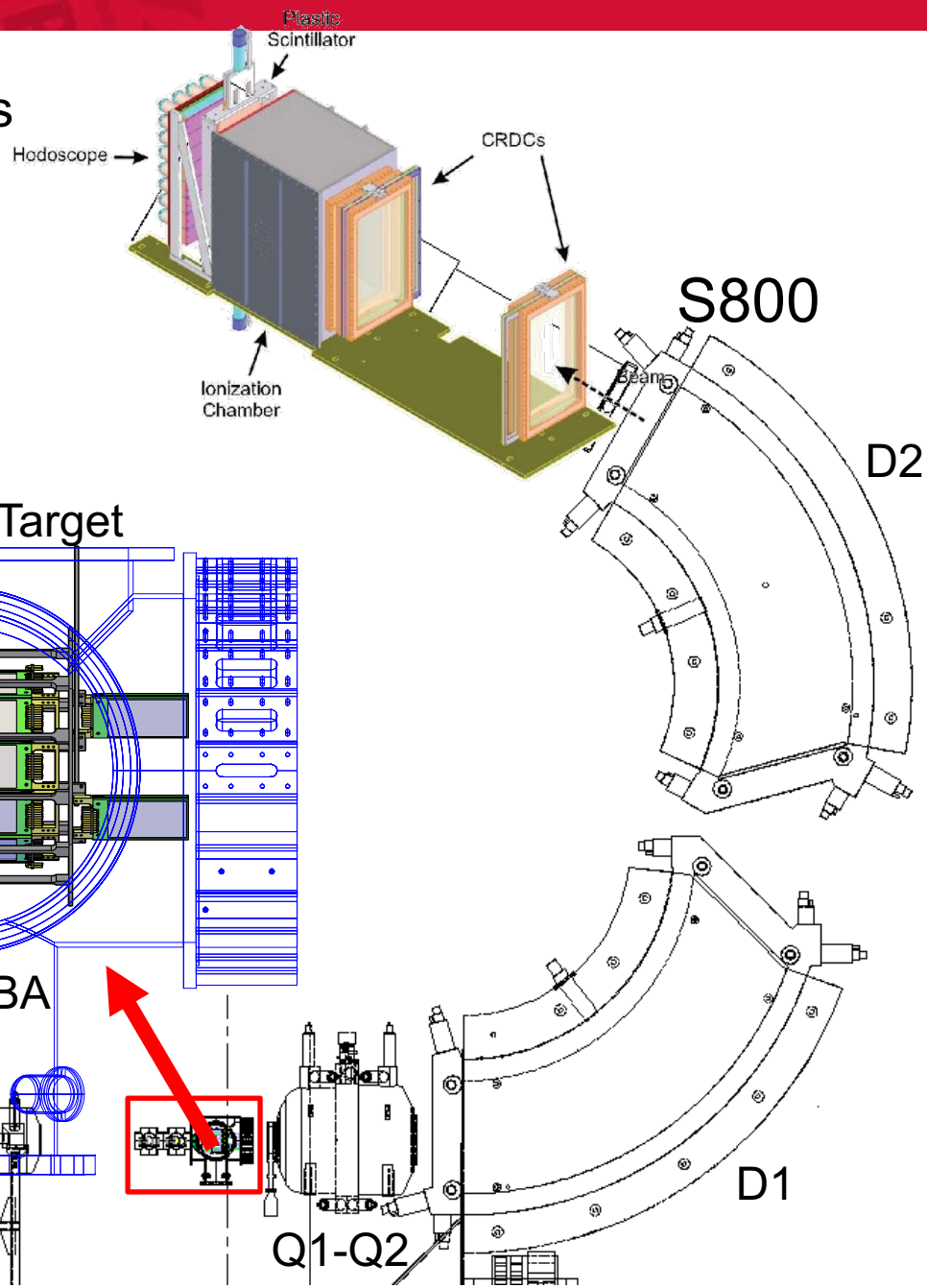
- Level density low near  $S_n$
- Direct neutron capture important
- Depends on
  - $E_x$  of low- $\ell$  single particle states
  - Spectroscopic factor  $S$

$$S = \left( \frac{d\sigma}{d\Omega} \right)_{exp} / \left( \frac{d\sigma}{d\Omega} \right)_{thy}$$

R. Surman et al., (weak-r process)  
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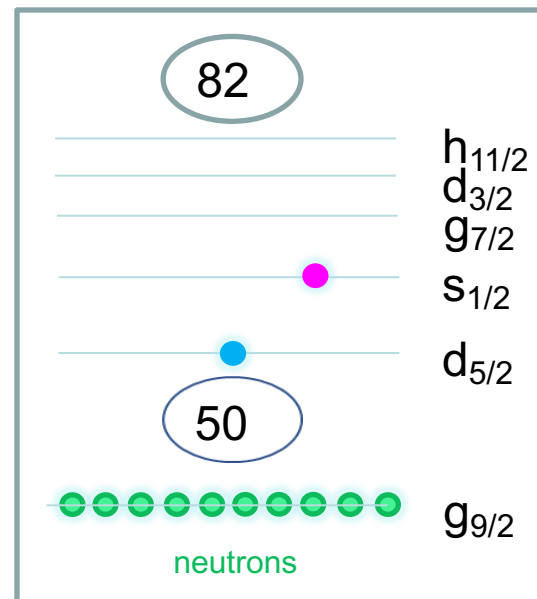
(d,p) with  $^{84}\text{Se}$  45 MeV/u NSCL beams

- $\text{CD}_2$  targets
- Upstream beam tracking
- ORRUBA (Oak Ridge Rutgers university Barrel Array) + SIDAR
- Heavy recoils S800 focal plane



■ 45 MeV/u at NSCL

$$S = \left( \frac{d\sigma}{d\Omega} \right)_{exp} / \left( \frac{d\sigma}{d\Omega} \right)_{thy}$$



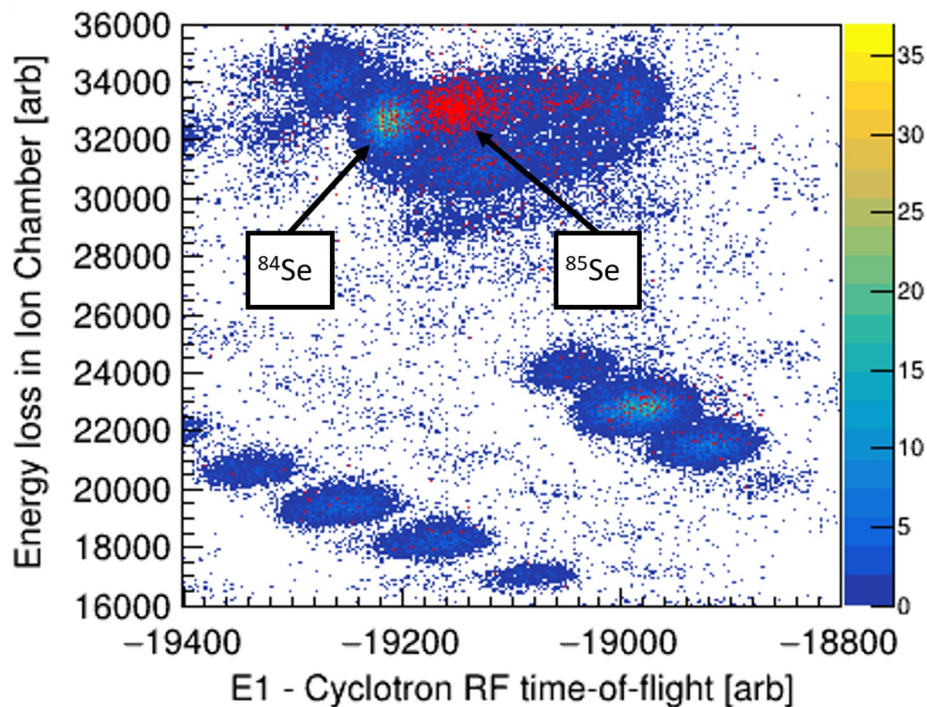
Excitations in  $^{85}\text{Se}$

S800 FP particle ID

- (red) gate on ORRUBA
- Gate on  $^{85}\text{Se}$
- Q-value spectra
- $\left( \frac{d\sigma}{d\Omega} \right)$

H.E. Sims Phd Dissertation (2020)

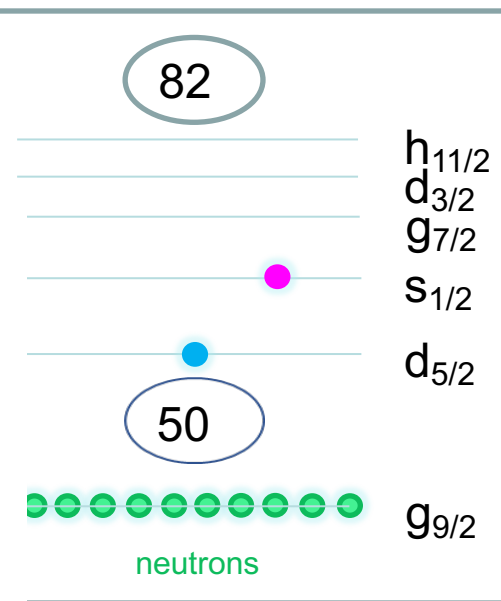
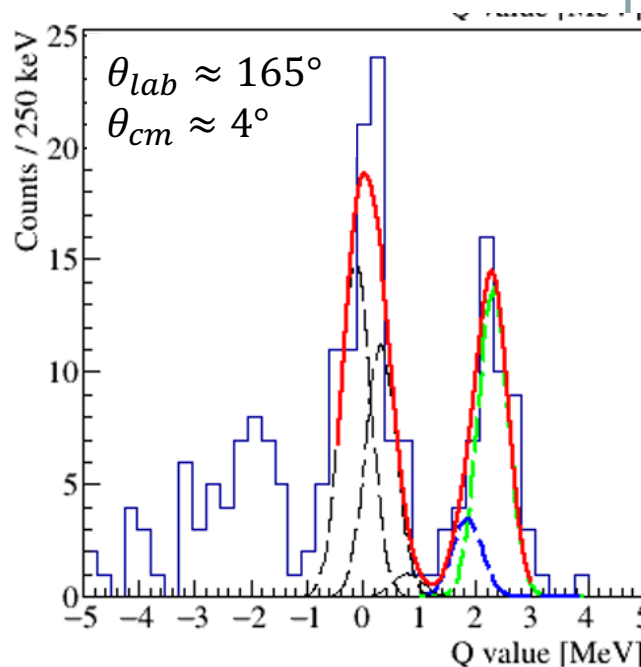
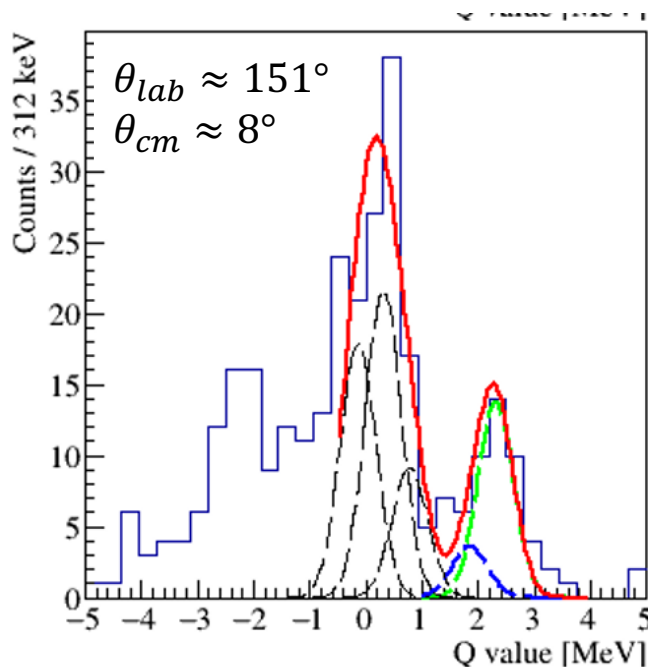
H.E. Sims, D Walter et al.,  
in preparation for PRC (2023)





■ 45 MeV/u at NSCL

$$S = \left( \frac{d\sigma}{d\Omega} \right)_{exp} / \left( \frac{d\sigma}{d\Omega} \right)_{thy}$$



Excitations in  $^{85}\text{Se}$

FR-ADWA w/ KD OPM

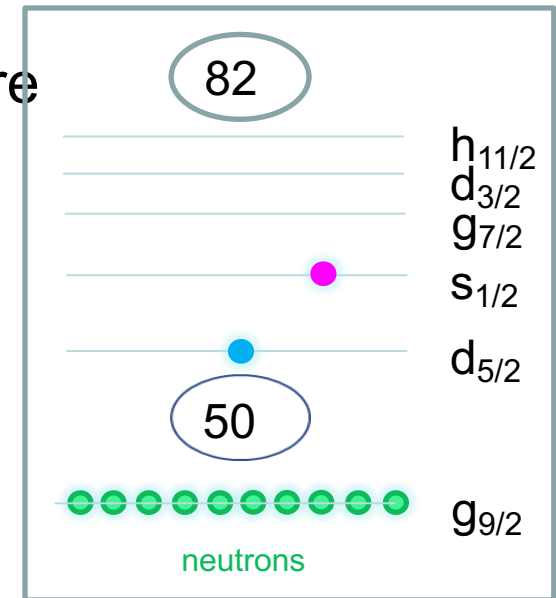
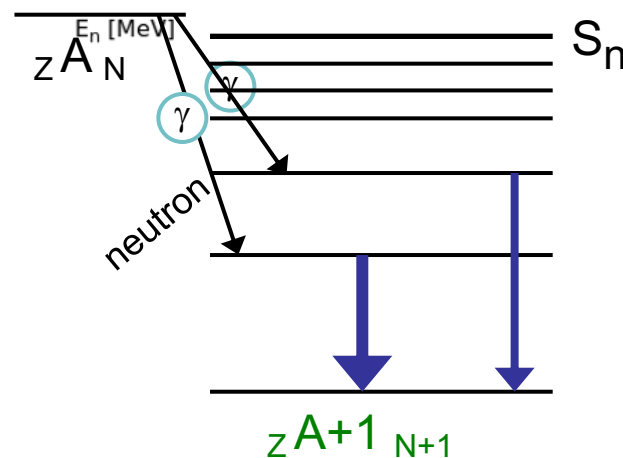
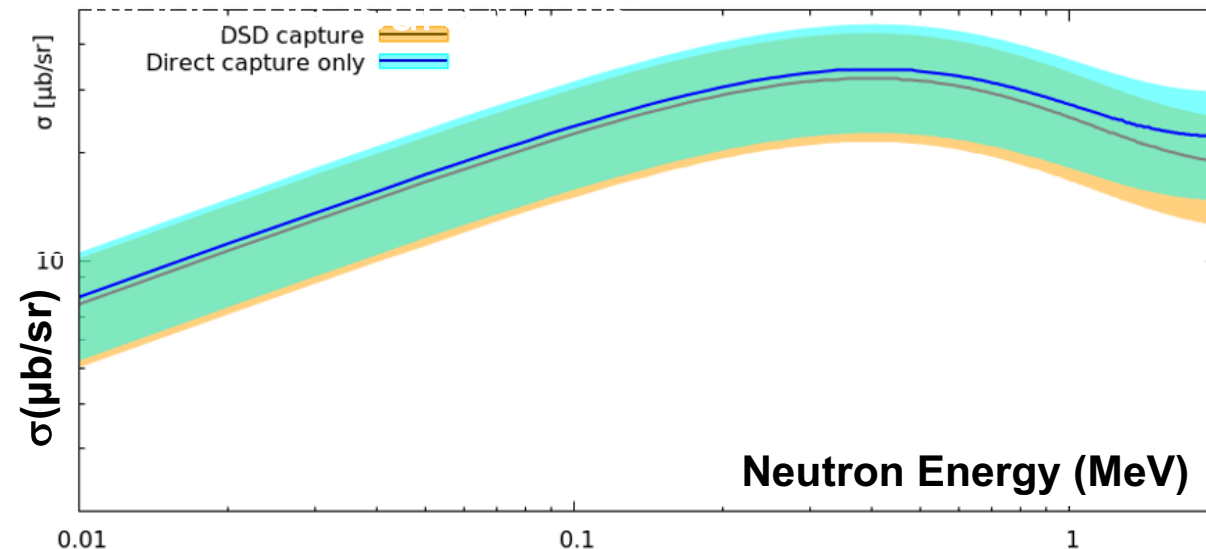
$5/2^+ E_x = 0$	$S = 0.28 (4)$
$1/2^+ E_x = 0.462 \text{ MeV}$	$S = 0.26 (6)$

H.E. Sims Phd Dissertation (2020)

H.E. Sims, D Walter et al.,  
in preparation for PRC (2023)

## Direct-semi-direct (DSD) capture

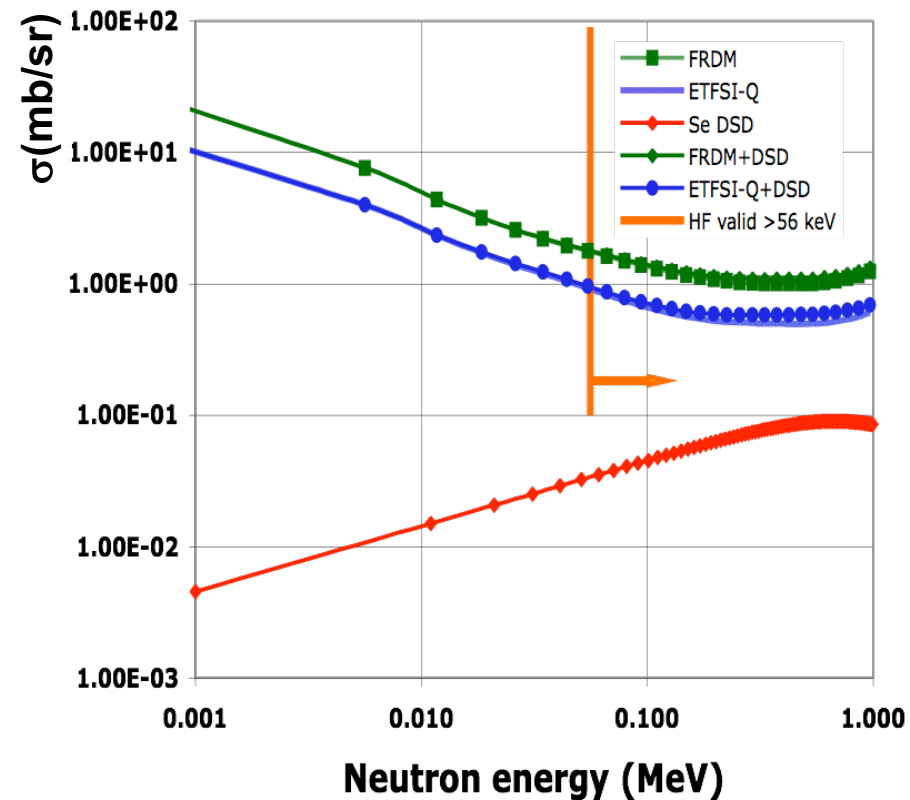
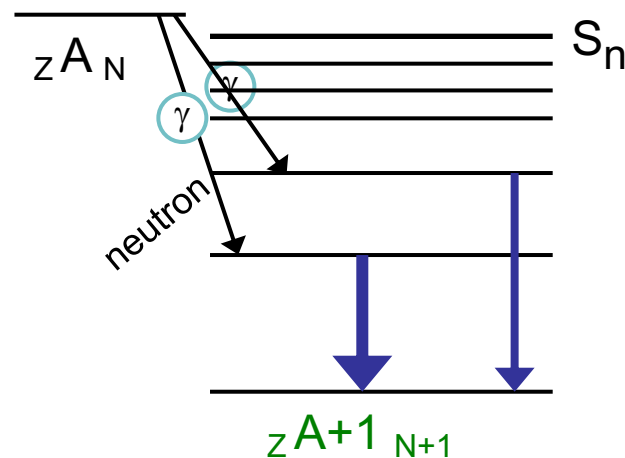
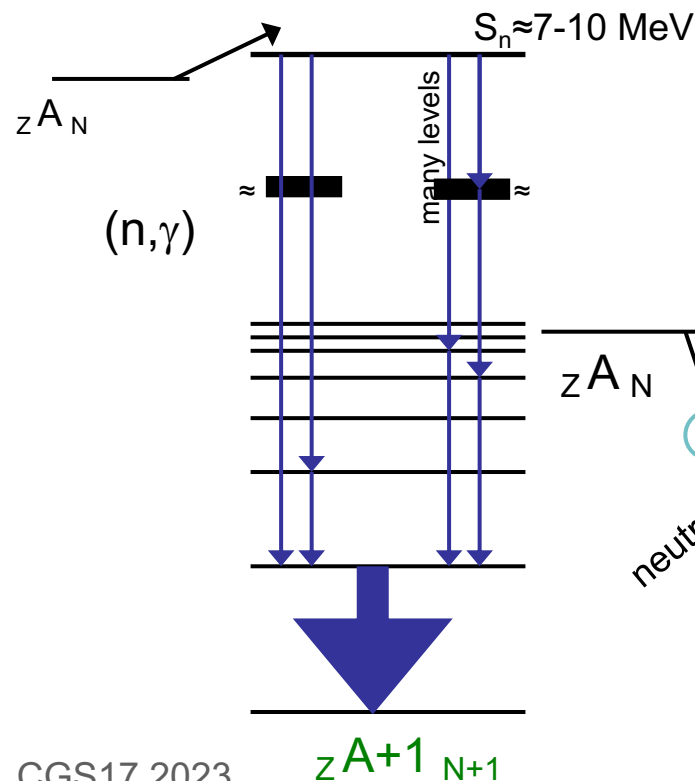
- Cross sections small  $\approx 20 \mu\text{b/sr}$ ;  $p$ -wave capture
- Statistical capture?  $\sigma$  much larger?



H.E. Sims Phd Dissertation (2020)  
 H.E. Sims, D Walter et al.,  
 in preparation for PRC (2023)

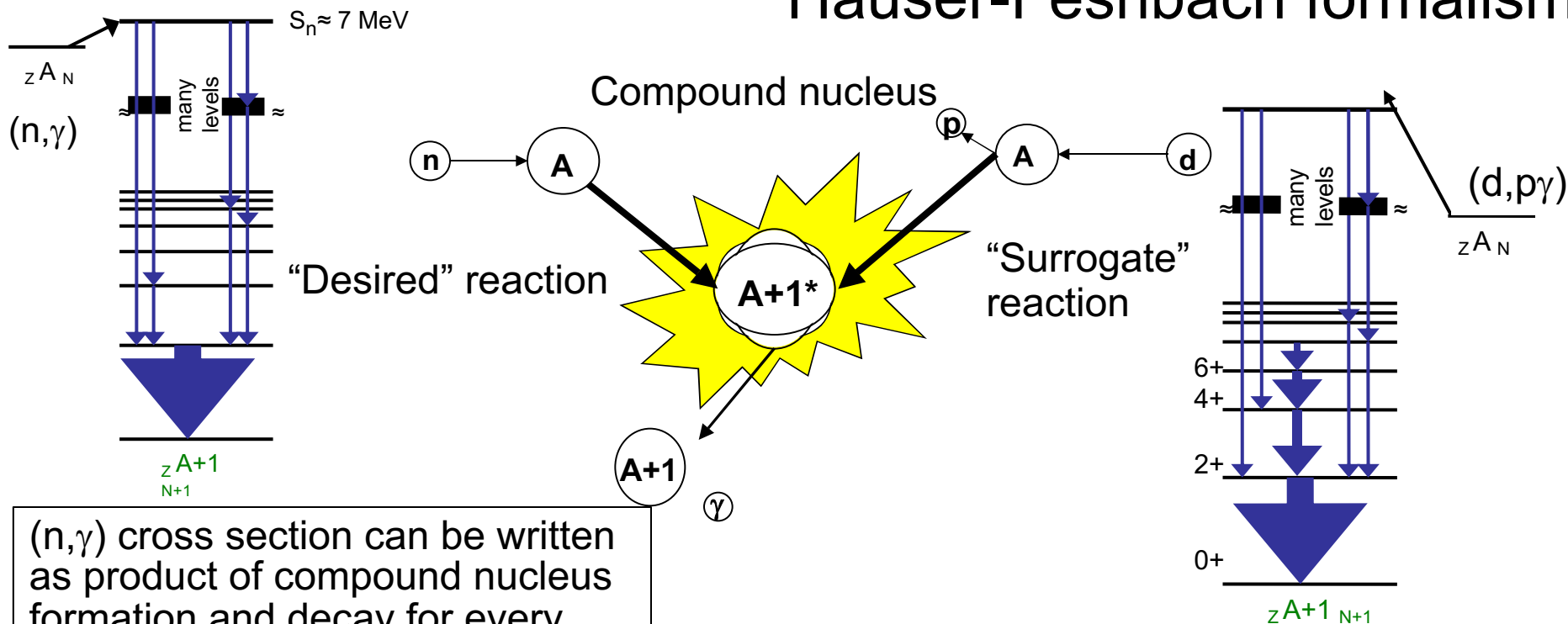
## Direct-semi-direct capture

- Cross sections small  $\approx 20 \mu\text{b/sr}$  for  $p$ -wave capture
- Statistical capture?  $\sigma$  much larger?
- Need valid  $(n,\gamma)$  surrogate reaction



H.E. Sims Phd Dissertation (2020)  
 H.E. Sims, D Walter et al.,  
 in preparation for PRC (2023)  
 J.A. Cizewski et al,  
 AIP CP **1090**, 463 (2009)

# Surrogate reaction concept & Hauser-Feshbach formalism



$(n, \gamma)$  cross section can be written as product of compound nucleus formation and decay for every spin and parity:

$$\sigma_{n\gamma}(E_n) = \sum_{J, \pi} \sigma_n^{CN}(E_x, J, \pi) G_\gamma^{CN}(E_x, J, \pi)$$

Surrogate particle-gamma coincidence can be written as product of compound nucleus formation and decay for every spin and parity:

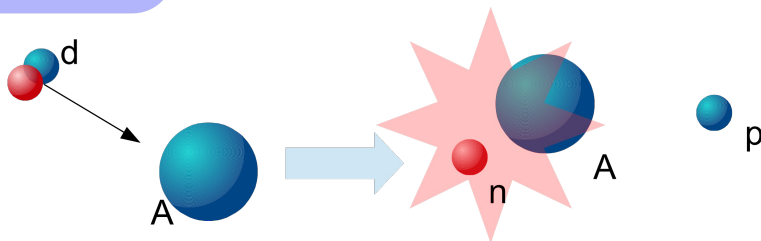
$$P_{p\gamma}(E_x, \theta) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_\gamma^{CN}(E_x, J, \pi)$$

$$P_{p\gamma}(E_x, \theta) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_{\gamma}^{CN}(E_x, J, \pi)$$

# Neutron transfer (d,p) to unbound states, non-elastic breakup and surrogate for (n, $\gamma$ )

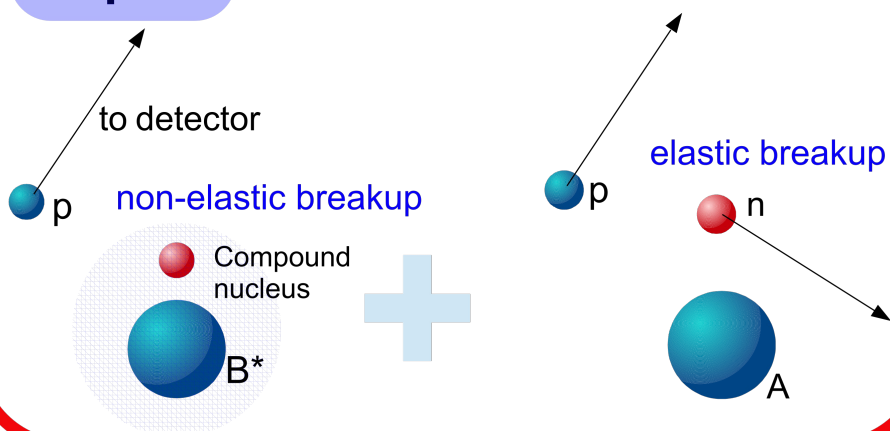
## step 1

separation of the proton



## step 2

propagation of  $n$  in the field of  $B^*$



## Two-step process

- d breakup; B.E. = 2.2 MeV
- n propagation
  - Elastic breakup
  - Non-elastic breakup  $\Rightarrow$  CN and surrogate (n, $\gamma$ )
  - Predicts  $J^\pi$  transfer

Gregory Potel et al. PRC 92, 034611(2015)  $\Rightarrow$  path to CN formation



**$(d,p)$  reaction to forms compound nucleus**

- ❖ Need to measure  $P(d,p\gamma)$
- ❖ Need theory to calculate formation of CN:  $F^{CN}$
- ❖ Need to deduce decay of CN:  $G^{CN}$

$$P_{p\gamma}(E_x, \theta) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_{\gamma}^{CN}(E_x, J, \pi)$$

Validate with  $^{95}\text{Mo}(d,p\gamma)$  reaction &  $^{96}\text{Mo}$  gammas  
 $\ell = 0$  capture on  $5/2^+ \Rightarrow 2^+, 3^+$

$\sigma(n,\gamma)$  was measured and informed

$$P_{pY}(E_x) = \frac{\text{Number of CN decays via channel } Y}{\text{Number of times the CN is formed}}$$

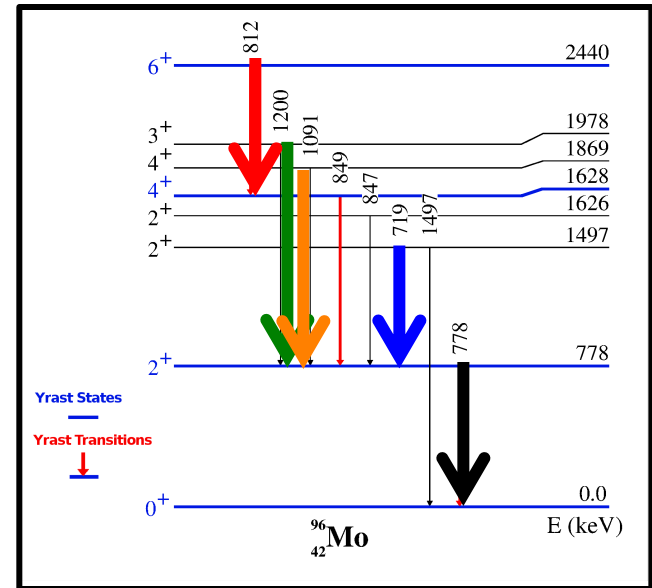
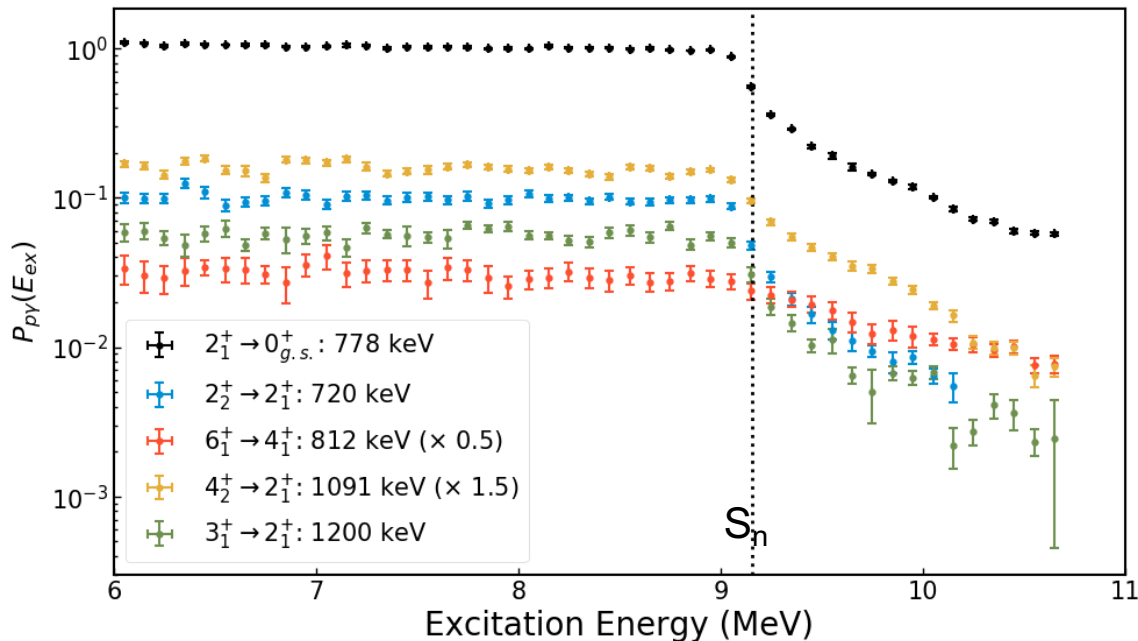
$$P_{pY}(E_x) = \frac{\text{Diagram: A cluster of red and blue spheres (representing a compound nucleus) decaying into a cluster of red and blue spheres (representing a residual nucleus) with a wavy line (representing a particle or photon) emitted. An arrow points from the initial cluster to the final cluster. Below the horizontal line is another cluster of red and blue spheres.}}{\text{Diagram: A cluster of red and blue spheres (representing a compound nucleus).}}$$

$$P_{pY}(E_x) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_Y^{CN}(E_x, J, \pi)$$

$$P_{pY}(E_x) = \frac{\text{Number of CN decays via channel } Y}{\text{Number of times the CN is formed}}$$

- Channel Y: individual discrete  $\gamma$  transitions to low-lying states
  - Intensity (=counts/efficiency) of specific transitions
- Number of times CN is formed
  - Intensity of single protons as a function of  $E_x$

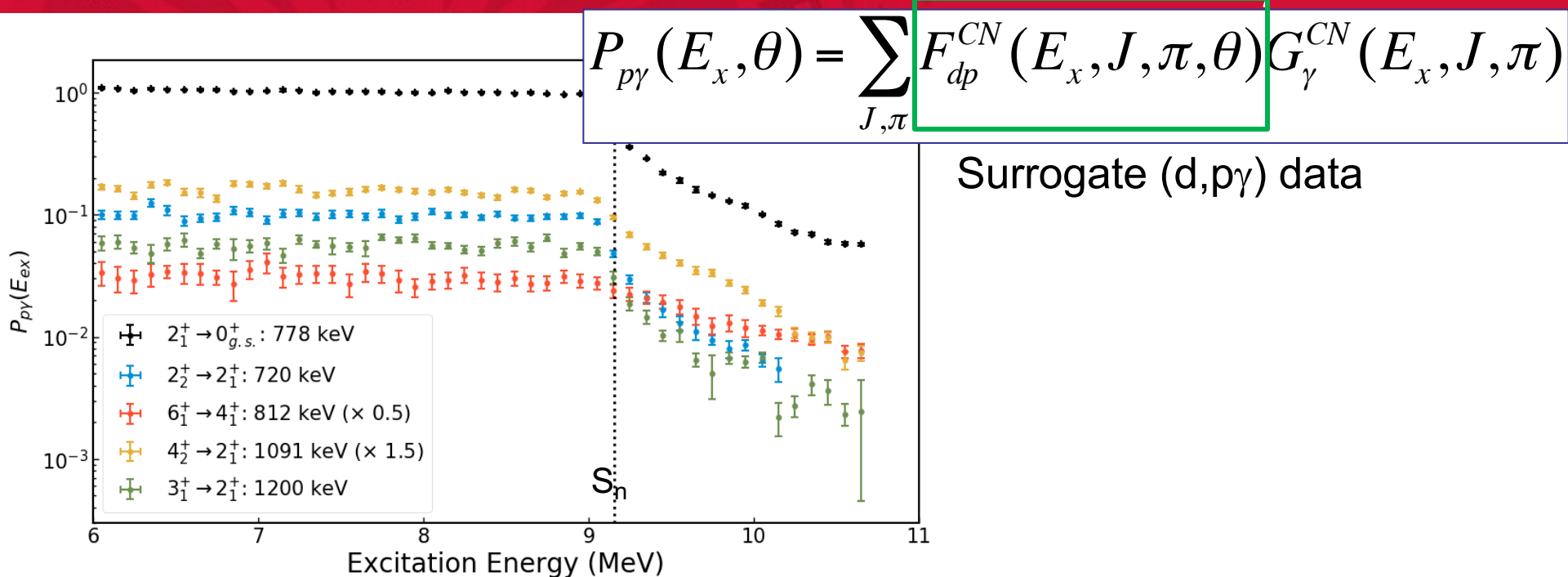
$$P_{pY}(E_x) = \frac{N_{pY}(E_x)}{\varepsilon_Y} \bigg/ N_p(E_x)$$



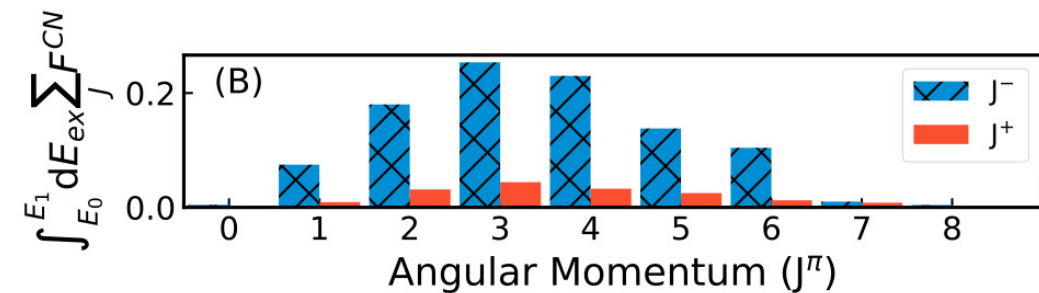
$$P_{pY}(E_x) = \frac{N_{pY}(E_x)/\varepsilon_Y}{N_p(E_x)} = \sum_{J,\pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_{\gamma}^{CN}(E_x, J, \pi)$$

Surrogate (d,p $\gamma$ ) data

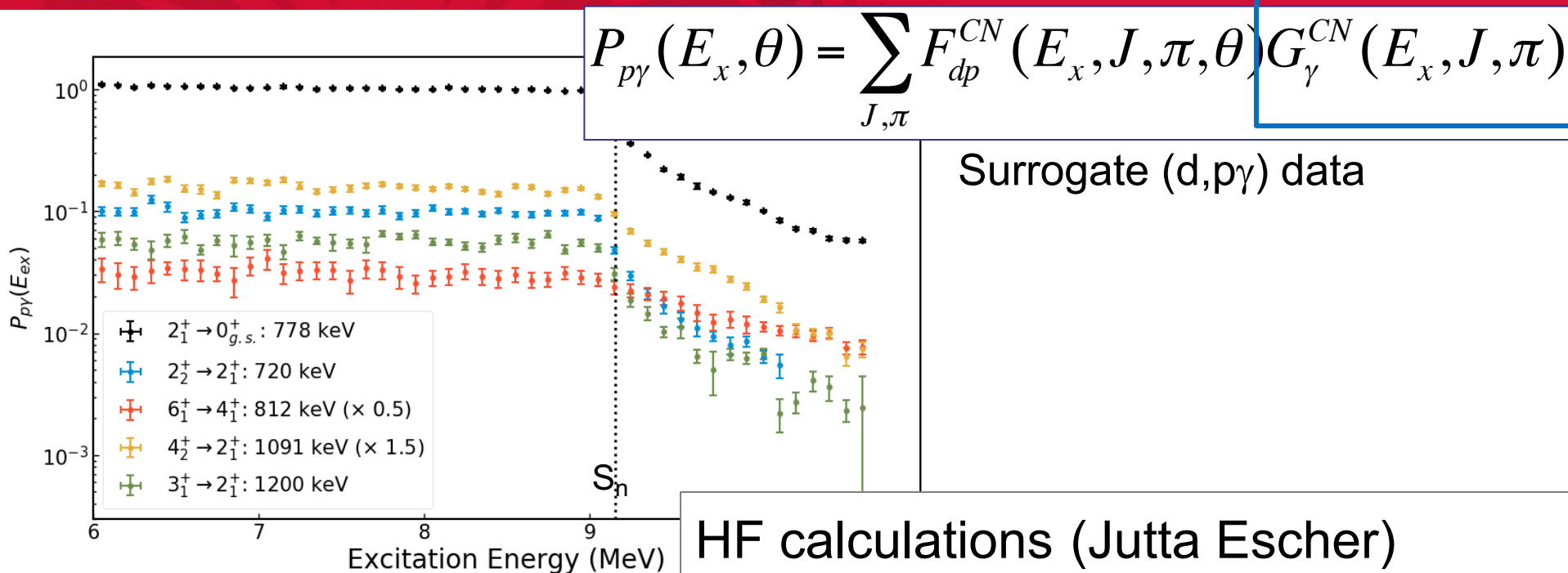
Measure  $P_{pY}(E_x)$  inform  $G_{\gamma}^{CN}(E_x, J, \pi)$



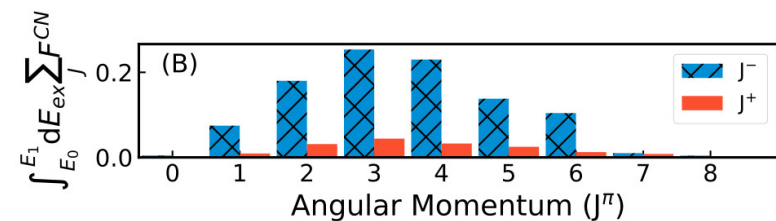
Potel:  $^{96}\text{Mo}$  spin distribution



G. Potel et al, PRC 92, 034611(2015)



Potel:  $^{96}\text{Mo}$  spin distribution



HF calculations (Jutta Escher)


- $F^{\text{CN}}$  from Gregory Potel
- Bayesian fit to observed  $P(d,p\gamma)$ 
  - Level density: Gilbert & Cameron
  - No norm to  $D_0$
  - Lorentzian  $\gamma$  strength function;
  - No  $\langle \Gamma(\gamma) \rangle$

➤  $G^{\text{CN}}(E_x, J, \pi)$

G. Potel et al, PRC 92, 034611(2015)

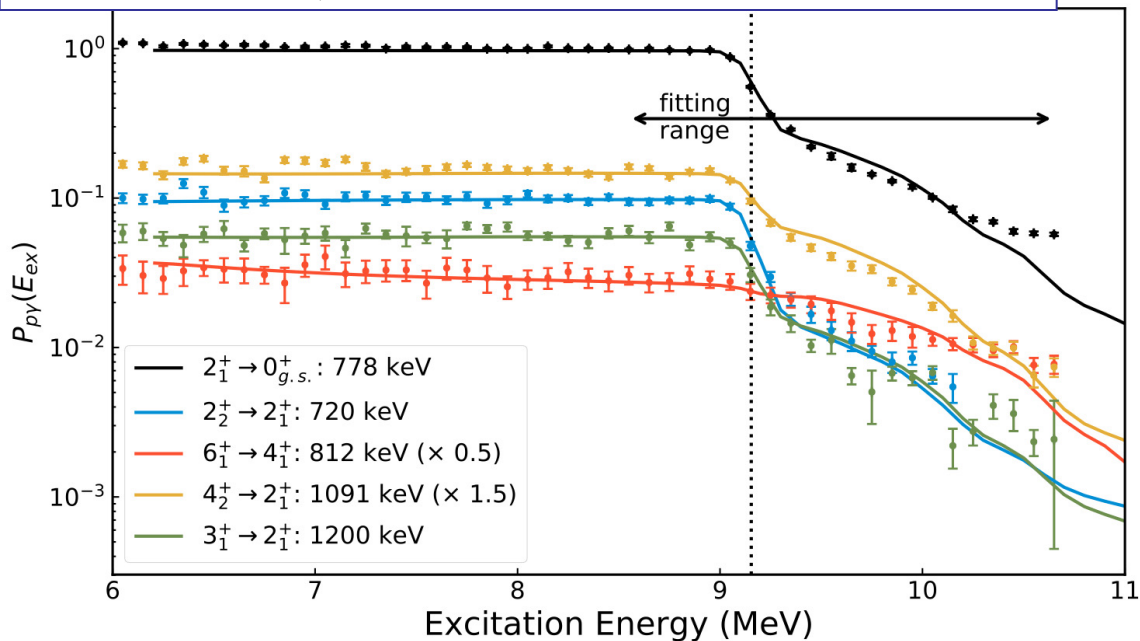


$$P_{p\gamma}(E_x, \theta) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_{\gamma}^{CN}(E_x, J, \pi)$$

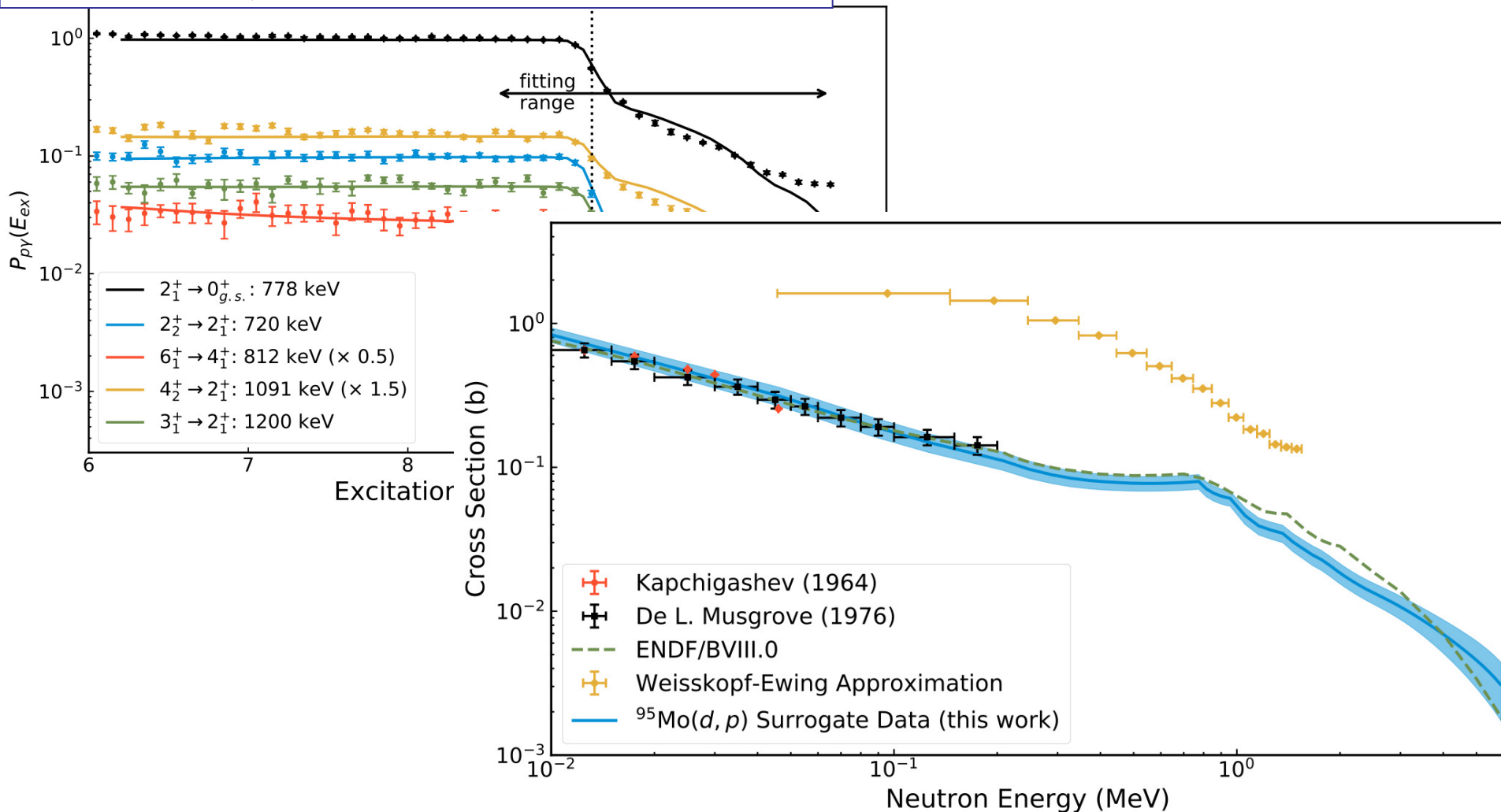
$$\sigma_{n\gamma}(E_n) = \sum_{J, \pi} \boxed{\sigma_n^{CN}(E_x, J, \pi)} G_{\gamma}^{CN}(E_x, J, \pi)$$


- Deduce  $G^{CN}(E_x, J, \pi)$  from fit to data
- Calculate  $\sigma^{CN}$  w/ Koning-Delaroche optical potentials
- Deduce  $\sigma(n, \gamma)$  vs  $E_x$

$$P_{p\gamma}(E_x, \theta) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_{\gamma}^{CN}(E_x, J, \pi)$$



$$P_{p\gamma}(E_x, \theta) = \sum_{J, \pi} F_{dp}^{CN}(E_x, J, \pi, \theta) G_{\gamma}^{CN}(E_x, J, \pi)$$



$$\sigma_{n\gamma}(E_n) = \sum_{J, \pi} \sigma_n^{CN}(E_x, J, \pi) G_{\gamma}^{CN}(E_x, J, \pi)$$

- Heavy beam on light ( $\text{CD}_2$ ) target = inverse kinematics
- Proton detection: good energy and angle resolution: ORRUBA
- Challenge: detecting discrete gammas
  - Relatively low gamma efficiency, especially discrete  $\gamma$
  - Away from even-even closed shells
    - High level density even at low  $E_x$
    - Especially final odd-odd nuclei
- Want  $Y$  – the gamma decay channel:
  - Not dependent on specific gammas

$$P_{pY}(E_x) = \frac{\text{Number of CN decays via channel } Y}{\text{Number of times the CN is formed}}$$

$$P_{pY}(E_x) = \frac{\text{[Diagram: A cluster of blue and red spheres transitions via an arrow to another cluster of blue and red spheres with a wavy line representing a gamma ray emission.]}}{\text{[Diagram: A single cluster of blue and red spheres.]}}$$

- $^{84}\text{Se}(\text{d},\text{p})$  populates  $^{85}\text{Se}^*$  CN
- CN at  $E_x < S_n$ : only decays by gamma emission  $\Rightarrow ^{85}\text{Se}$
- CN at  $E_x > S_n$ : if decays by gamma emission  $\Rightarrow ^{85}\text{Se}$  = channel Y
- CN at  $E_x > S_n$ : if decays by neutron emission  $\Rightarrow ^{84}\text{Se}$

$$P_{p\gamma}(E_x) = \frac{N_{p-^{85}\text{Se}}(E_x)}{\varepsilon} \bigg/ N_p(E_x)$$

$$P_{p\gamma}(E_x) = \frac{\text{Number of CN decays via channel Y}}{\text{Number of times the CN is formed}}$$

$$P_{p\gamma}(E_x) = \frac{\text{[Diagram: A blue and red nucleus decaying into a blue and red nucleus with a wavy line representing gamma emission]}}{\text{[Diagram: A blue and red nucleus]}}$$

- $^{84}\text{Se}(\text{d},\text{p})$  populates  $^{85}\text{Se}^*$  CN
- CN at  $E_x < S_n$ : only decays by gamma emission  $\Rightarrow ^{85}\text{Se}$
- CN at  $E_x > S_n$ : if decays by gamma emission  $\Rightarrow ^{85}\text{Se}$  = channel Y
- CN at  $E_x > S_n$ : if decays by neutron emission  $\Rightarrow ^{84}\text{Se}$

$$P_{p\gamma}(E_x) = \frac{N_{p-^{85}\text{Se}}(E_x)}{\varepsilon} \bigg/ N_p(E_x)$$

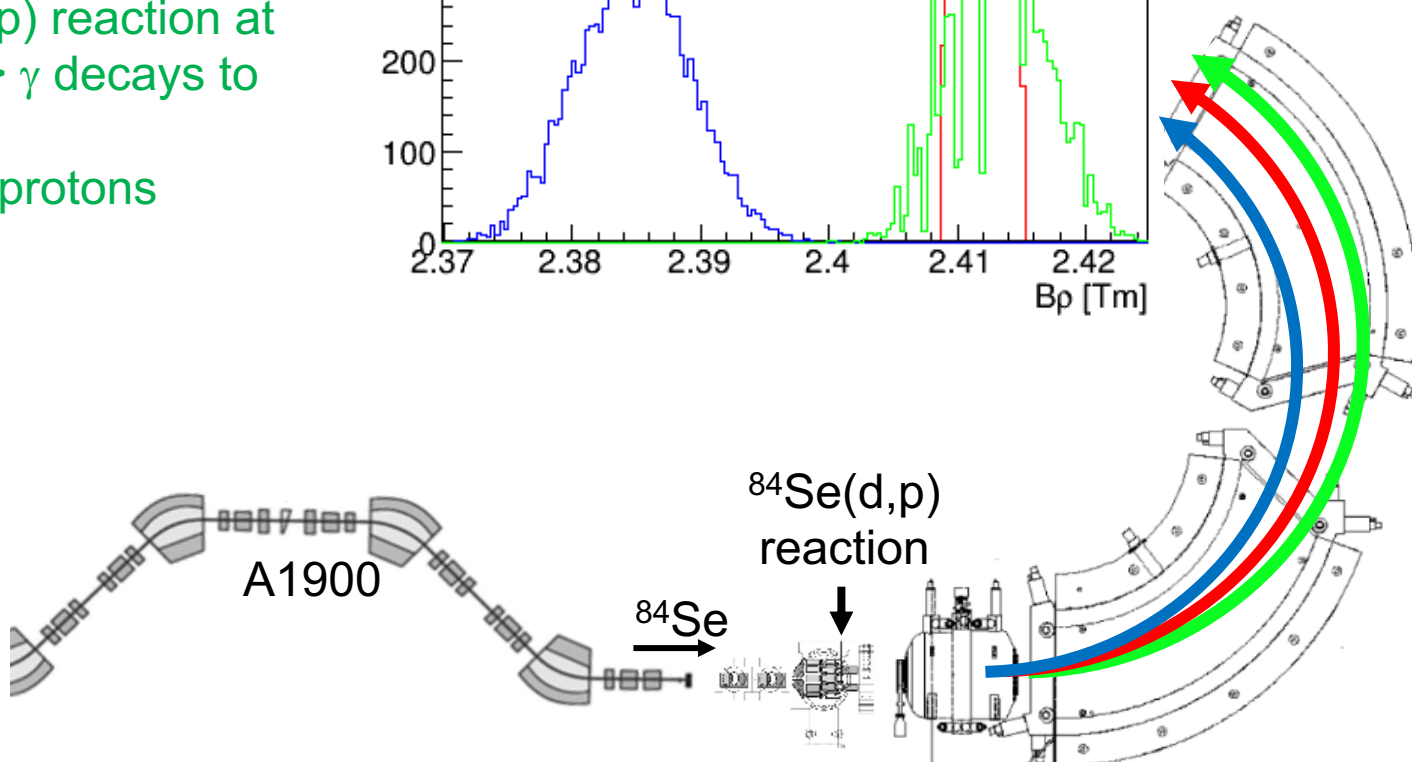
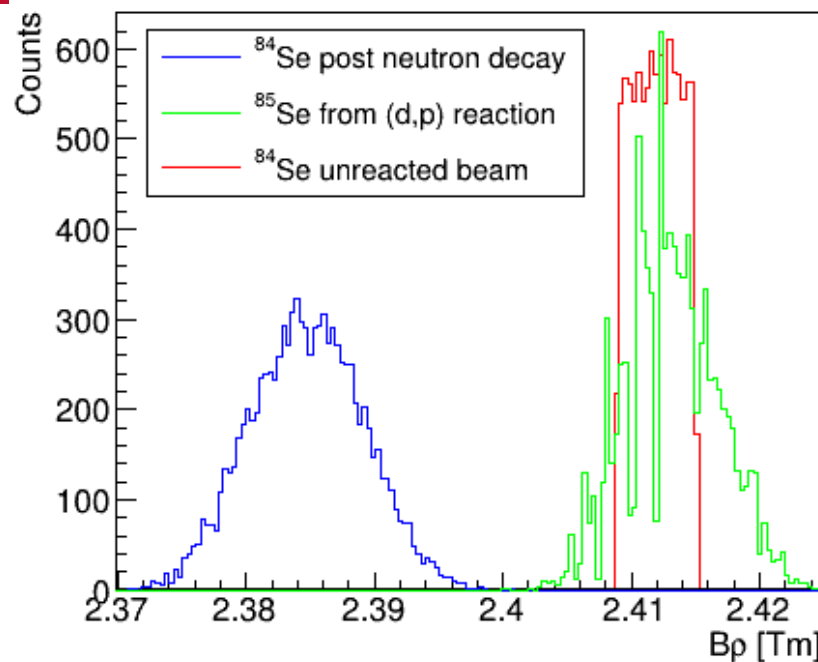
- Need excellent separation of  $^{85}\text{Se}$  and  $^{84}\text{Se}$
- Detection efficiency of heavy recoils  $>$  gammas
- No dependence on details of  $\gamma$ -decay

$$P_{p\gamma}(E_x) = \frac{\text{[Diagram: A nucleus (blue and red spheres) decaying into a nucleus (blue and red spheres) and a wavy line representing a gamma ray]}}{\text{[Diagram: A nucleus (blue and red spheres)]}}$$



Three scenarios:

1.  $^{84}\text{Se}$  does not react with  $\text{CD}_2$  target, continues with same momentum distribution as determined by slits in A1900
2.  $^{84}\text{Se}$  undergoes (d,p) reaction at  $\text{CD}_2 \Rightarrow \text{CN } ^{85}\text{Se} \Rightarrow \gamma$  decays to  $^{85}\text{Se}$  g.s.
  - Know  $E_x$  from protons

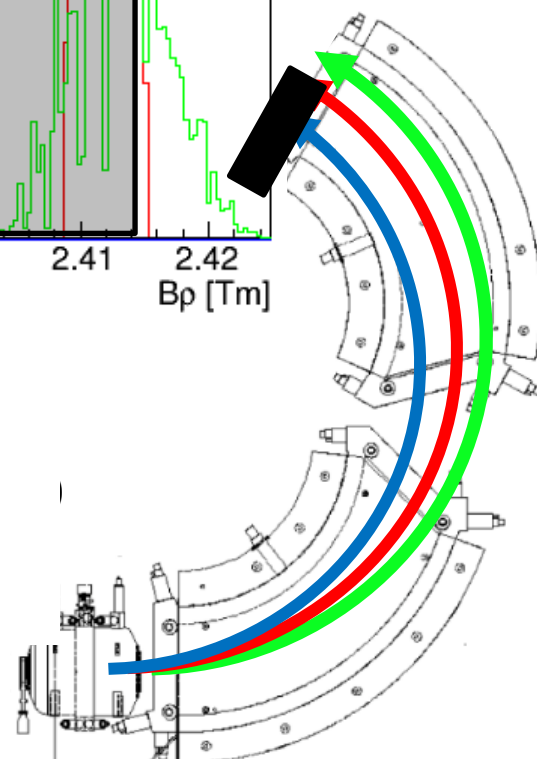
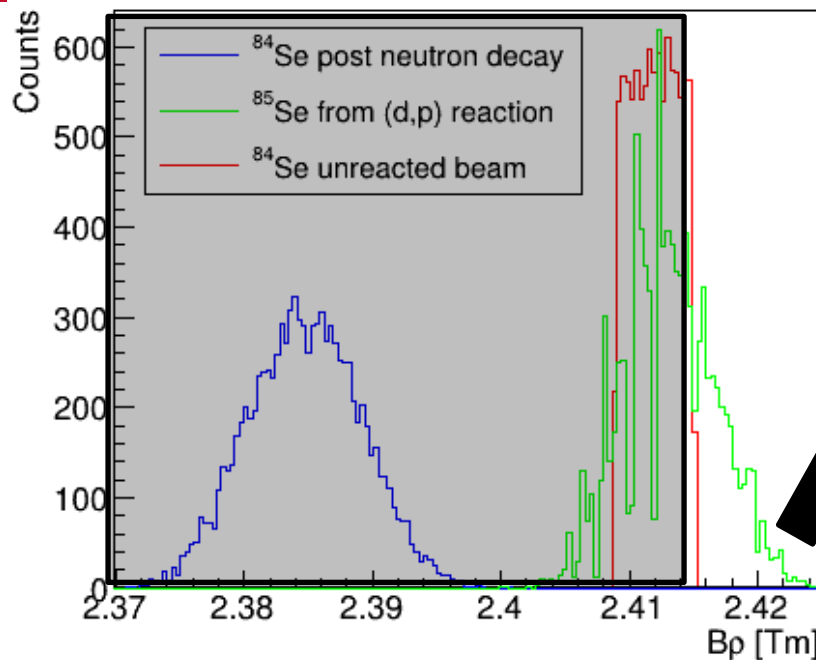


3. Same as point 2, except CN  $^{85}\text{Se}$  emits neutron  $\Rightarrow ^{84}\text{Se}$

Three scenarios:

1.  $^{84}\text{Se}$  does not react with  $\text{CD}_2$  target, continues with same momentum distribution as determined by slits in A1900
2.  $^{84}\text{Se}$  undergoes (d,p) reaction at  $\text{CD}_2 \Rightarrow \text{CN } ^{85}\text{Se} \Rightarrow \gamma$  decays to  $^{85}\text{Se}$  g.s.
  - Know  $E_x$  from protons

- S800 is rate-limited to  $\sim 5$  kHz  $\Rightarrow$  implement blocker
- **Use the recoils to determine whether n/ $\gamma$  decay ( $^{84}\text{Se}/^{85}\text{Se}$ ) – by tagging on the  $^{85}\text{Se}$**



3. Same as point 2, except CN  $^{85}\text{Se}$  emits neutron  $\Rightarrow ^{84}\text{Se}$

## Advantages:

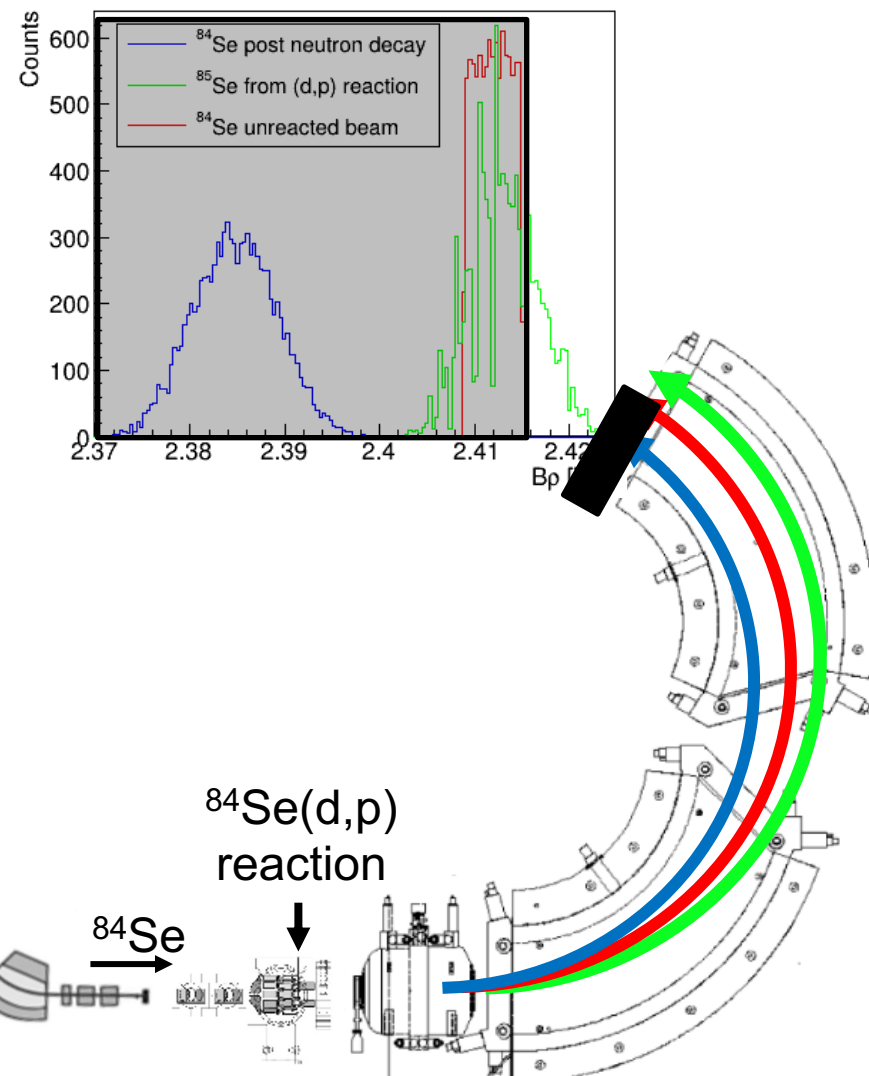
- With (low intensity RIBs) all statistics in single observable
- ~25-30% detection efficiency (much better than  $\gamma$  efficiency  $\approx 13\%$ )
  - Can measure by looking at bound states
  - Not reliant on simulations
  - If can tighten up momentum acceptance, less beam-recoil overlap
- No need for complicated cascade info – get emission probability without knowledge of how gamma decay occurs

## Difference:

- No details or constraint on specific gamma branches or cascade

## Challenges:

- Need significant characterization of background from Carbon in target



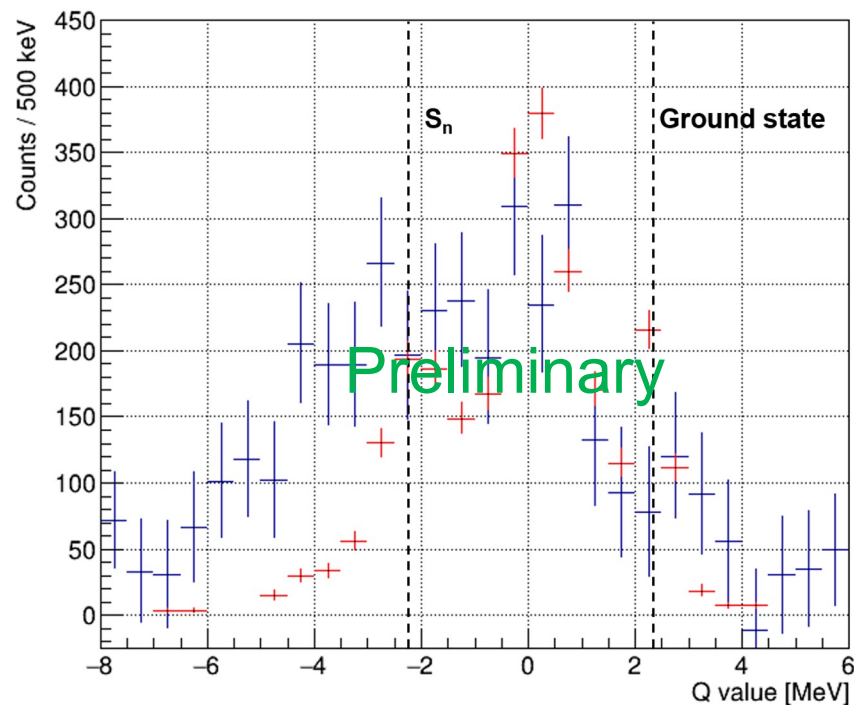
## $P_Y$ from S800 coincidences

$$P_Y(E_x) = \frac{N_{p-^{85}\text{Se}}(E_x) / \epsilon(^{85}\text{Se})}{N_{p\text{-singles}}(E_x)}$$

Proton-S800 coincidences

S800 acceptance

Proton singles (background subtracted)



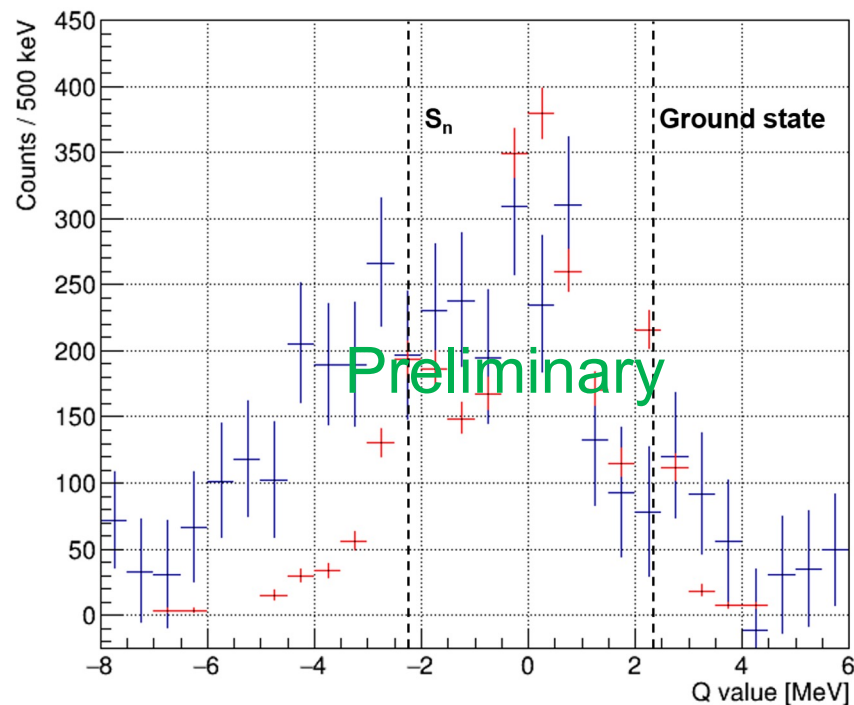
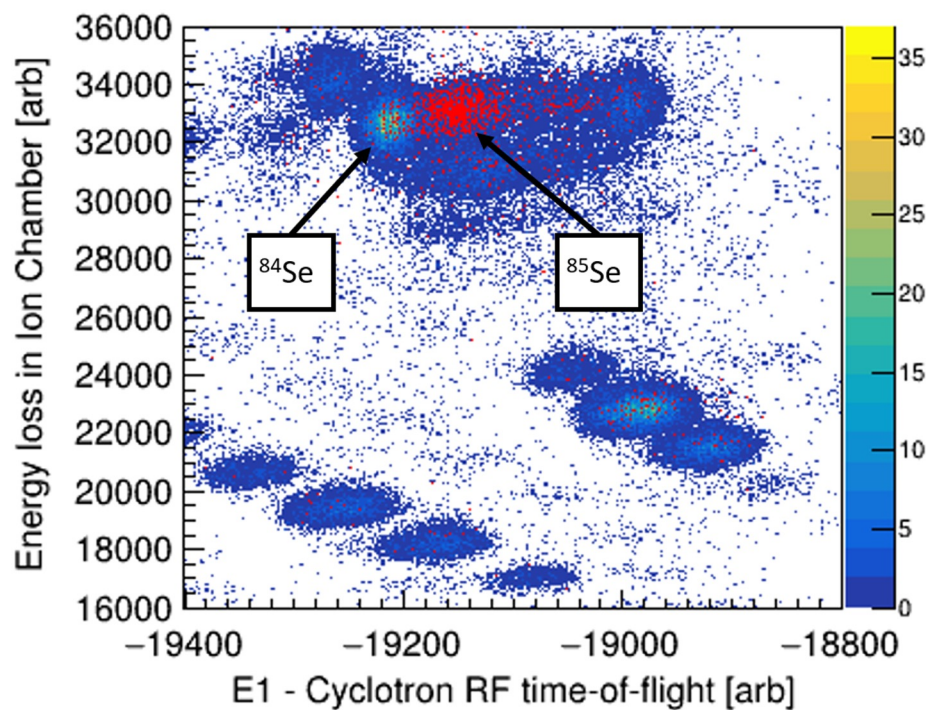
$P_\gamma$  from S800 coincidences

$$P_\gamma(E_x) = \frac{N_{p-^{85}\text{Se}}(E_x) / \epsilon(^{85}\text{Se})}{N_{p\text{-singles}}(E_x)}$$

Proton-S800 coincidences

S800 acceptance

Proton singles (background subtracted)





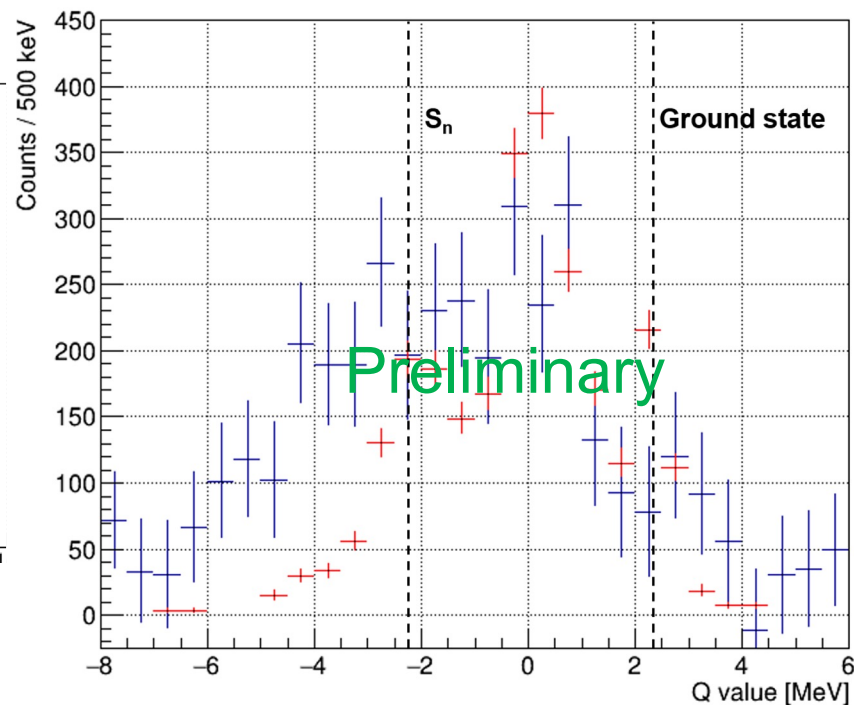
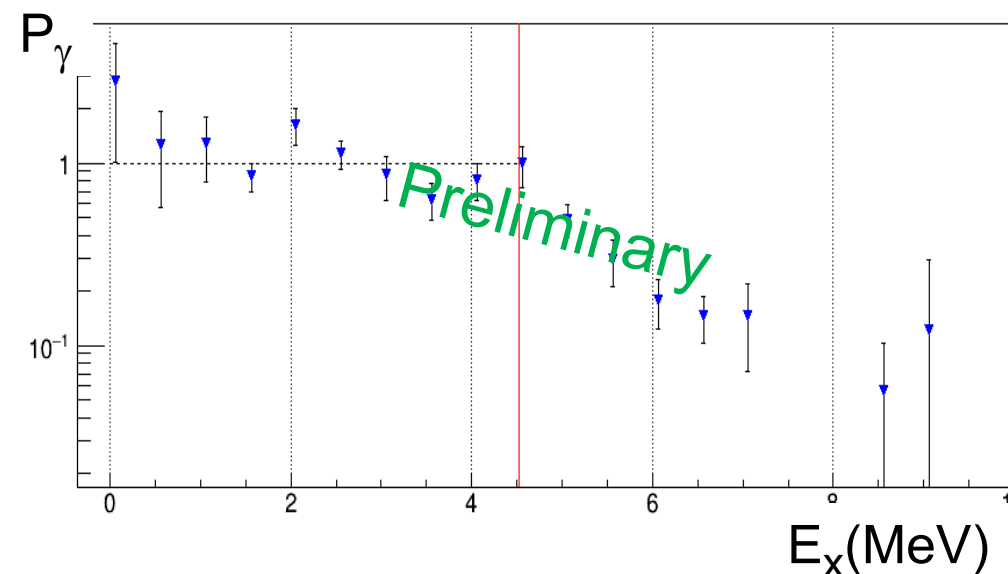
$P_\gamma$  from S800 coincidences

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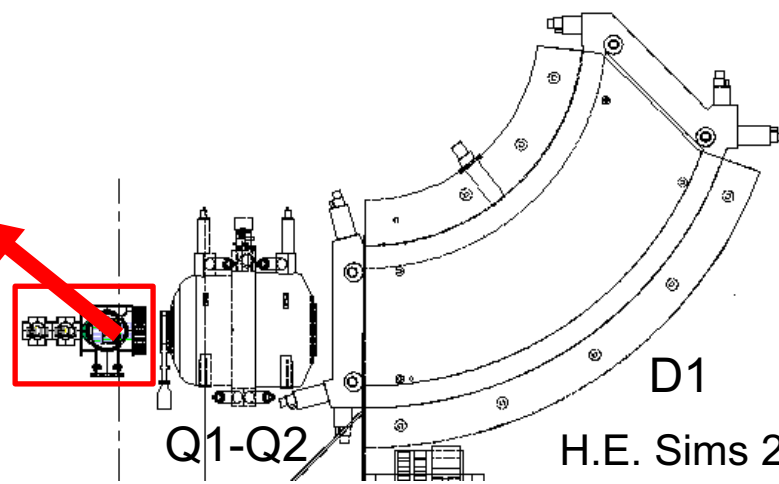
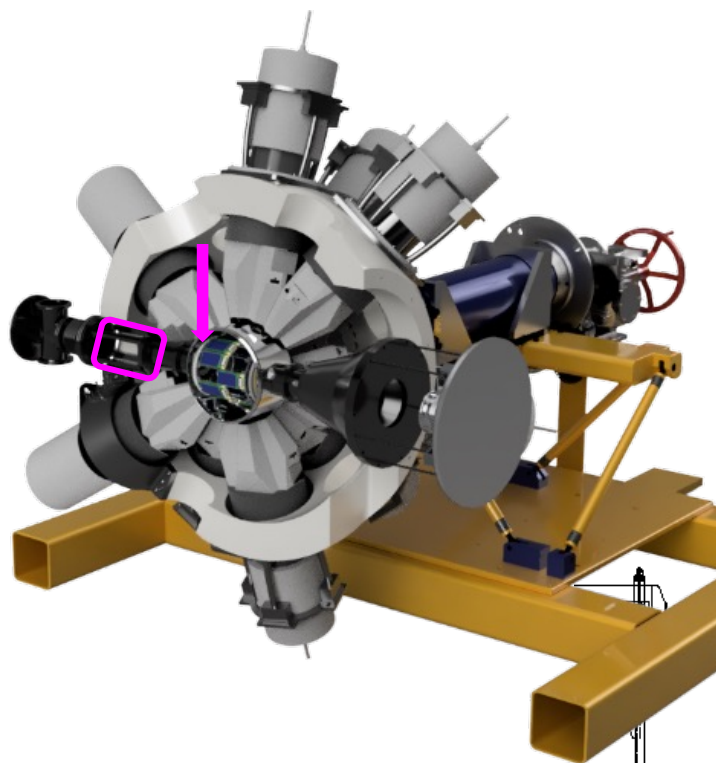
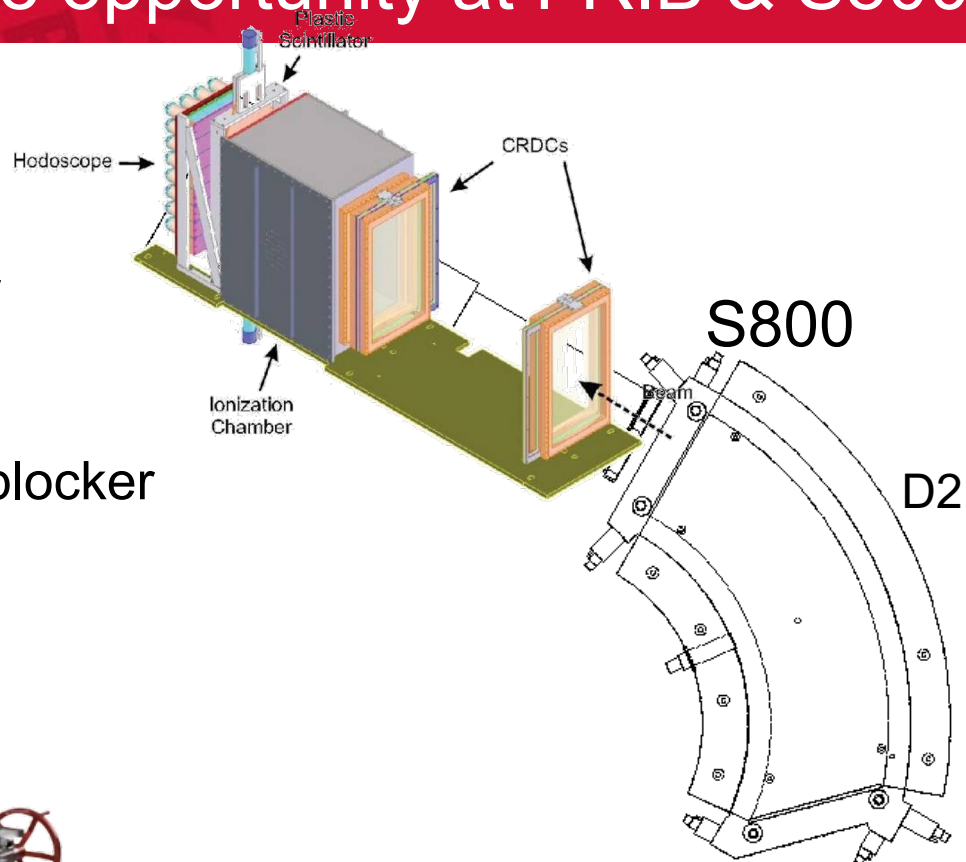
Proton-S800 coincidences

S800 acceptance

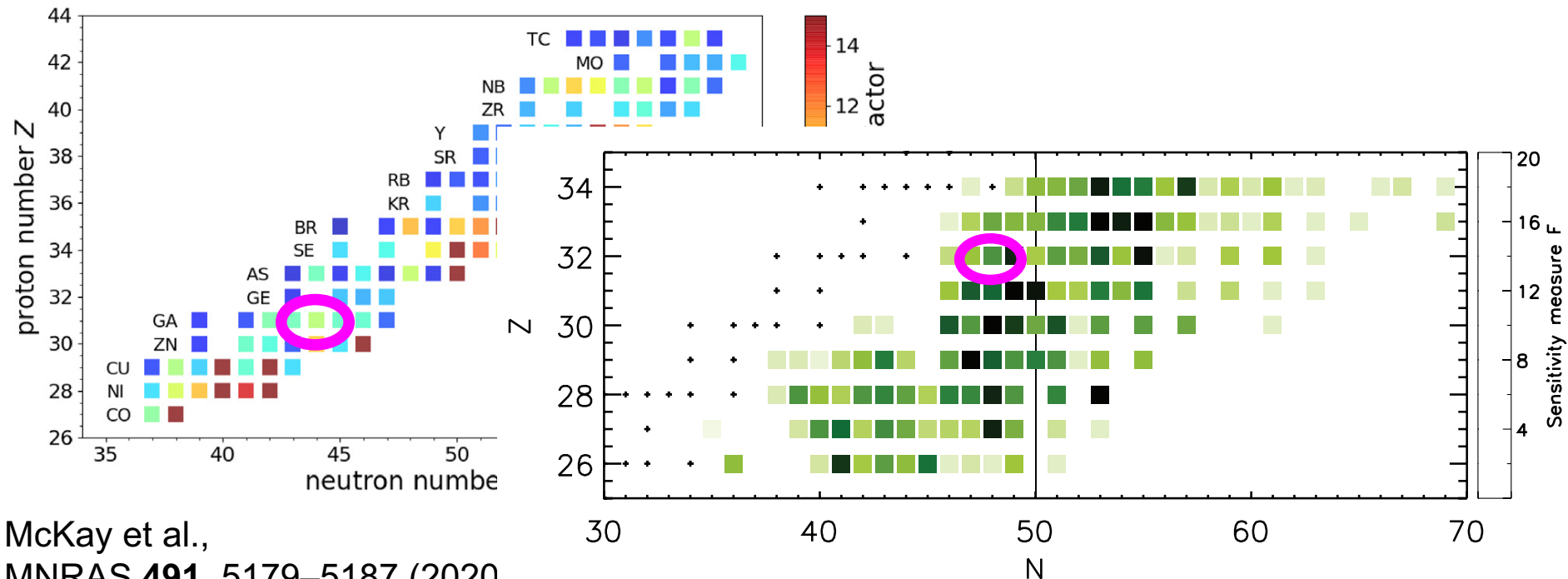
Proton singles (background subtracted)



- ORRUBA + GRETINA at S800
  - Upstream (gas) beam tracking detectors
  - New (smaller footprint) annular QQQ6 detectors
- S800 excellent PID
  - Can separate isotopes; beam blocker
- Requires  $\text{CD}_2$  and  $\text{CH}_2$  data



- Approved:  $(d,p\gamma)$  with  $\approx 45$  MeV/u  $^{80}\text{Ge}$  (N=48) and  $^{75}\text{Ga}$  beams + ORRUBA + GRETINA + S800
- Inform  $\sigma(n,\gamma)$ 
  - No gamma surrogate reaction method
  - Discrete gamma SRM; Gamma rays would confirm isotopics
- Inform i- and weak r-process nucleosynthesis

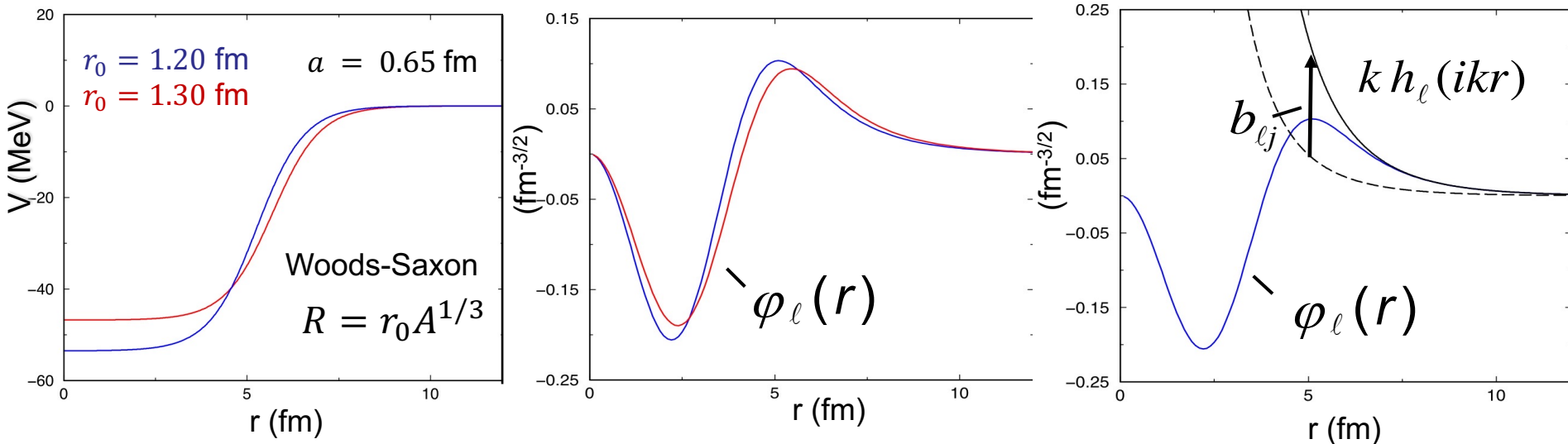




- (d,p) and (d,p $\gamma$ ) reactions inform i- and weak r- process  $A \approx 80$  nucleosynthesis
- Direct-semi-direct capture near neutron closed shells
  - Measure spectroscopic factors with (d,p)
- (d,p $\gamma$ ) validated surrogate reaction method (SRM) for (n, $\gamma$ )
  - Measure discrete gammas
    - Inform LD and  $\gamma$ SF  $\Rightarrow G_{\gamma}^{CN}(E_x, J, \pi) \Rightarrow$  inform  $\sigma(n, \gamma)$
- Prospects for No Gamma Surrogate (NGS) reaction method
  - Measure total population  $A+1$  nucleus
  - Details of gamma decay not needed
    - $G_{\gamma}^{CN}(E_x, J, \pi) \Rightarrow$  inform  $\sigma(n, \gamma)$
- Approved  $^{80}\text{Ge}$ ,  $^{75}\text{Ga}(d, p\gamma)$  at FRIB  
ORRUBA+GRETINA+S800

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ORRUBA+GRETINA+S800

# EXTRA SLIDES



Asymptotically if wave function is pure single-particle, e.g.,  $2d_{5/2}$  neutron:

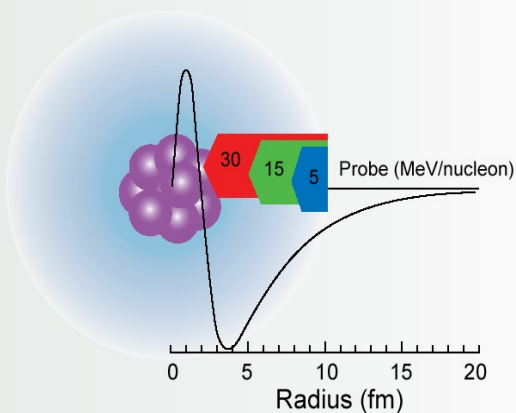
$$\varphi_\ell \rightarrow b_{\ell j} k h_\ell(ikr)$$

Single particle asymptotic normalization coefficient  $b_{\ell j}$  reflects potential's  $(r_0, a)$   
 But usually wave function is not a pure single-particle,  
 Rather overlap with a single particle w.f.  $C_{\ell j} k h_\ell(ikr) = S_{\ell j}^{1/2} b_{\ell j} k h_\ell(ikr)$

Defines the nuclear ANC  $C^2$  (asymptotic normalization coefficient)

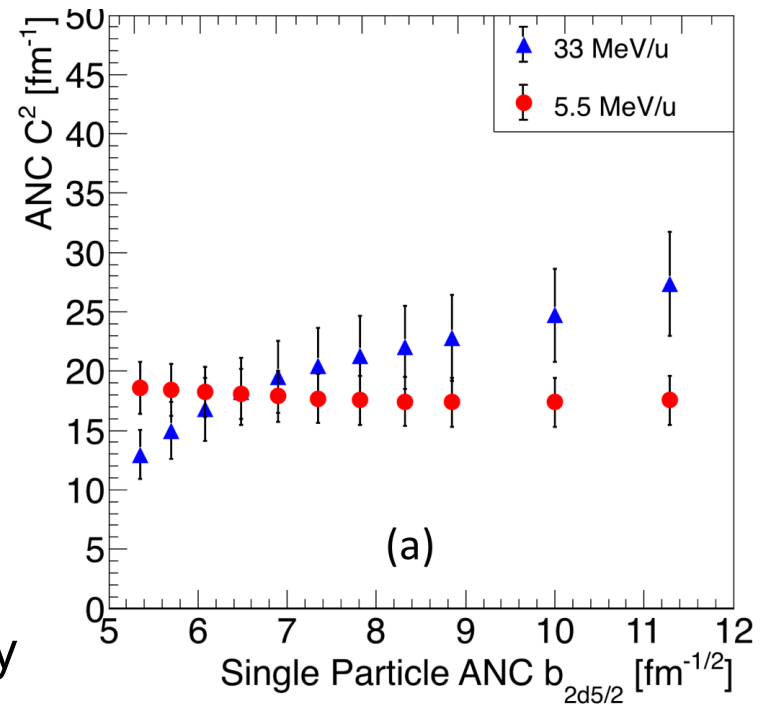
- proportional to the spANC  $b^2$
- proportionality constant  $S$ , the spectroscopic factor

$$C_{\ell j}^2 = S_{\ell j} b_{\ell j}^2$$



$$S = \left( \frac{d\sigma}{d\Omega} \right)^{\text{exp}} / \left( \frac{d\sigma}{d\Omega} \right)^{\text{theory}}$$

$$C_{\ell j}^2 = S_{\ell j} b_{\ell j}^2$$



$^{86}\text{Kr}(d,p)$  g.s.

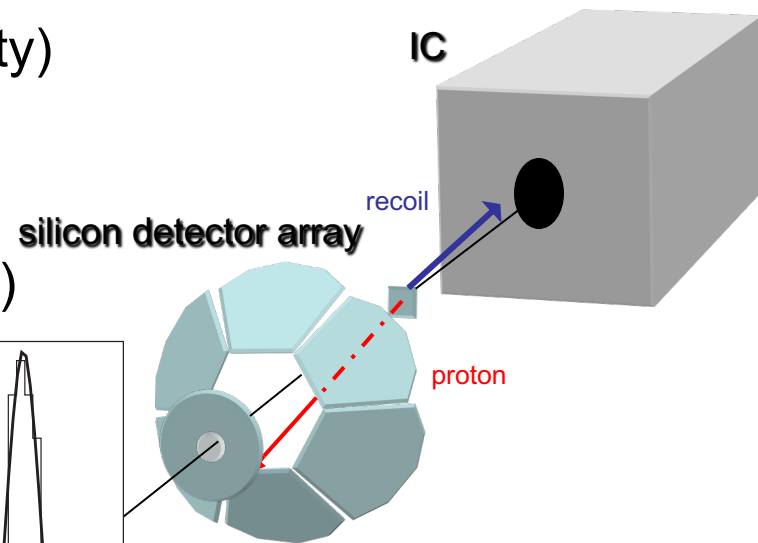
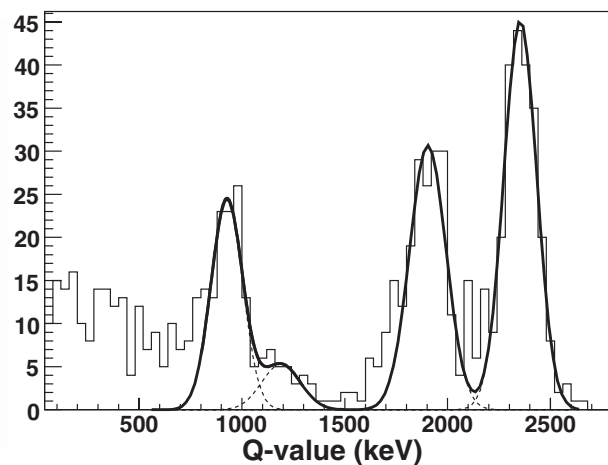
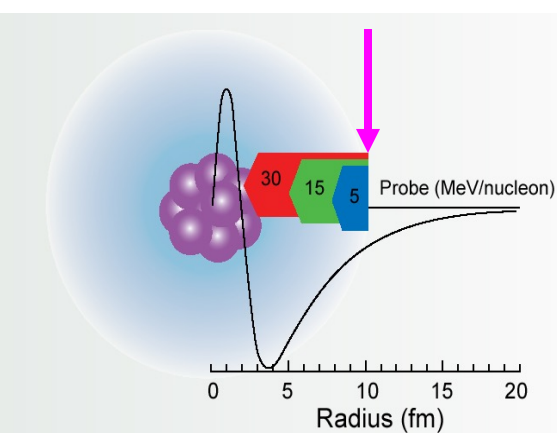
D. Walter et al. PRC **99**, 054625 (2019)

- Fix nuclear ANC ( $C_{\ell j}$ ) using peripheral reaction (lower energy)
- Probe the nuclear interior with higher energy reaction
  - ANC is property of state NOT reaction
- Combine: Constrain single-particle ANC
- $S$  dominated by uncertainties in experimental cross-section measurement rather than uncertainties in bound state potential

**Measure reactions at TWO different energies:**

A. Mukhamedzhanov and F. Nunes, Phys. Rev. C **72**, 017602 (2005)

- 4.5 MeV/u at HRIBF (ORNL former facility)
  - Silicon detector array (SIDAR)
  - Ion chamber recoil detector
  - J.S. Thomas et al. PRC **76** 044302 (2007)

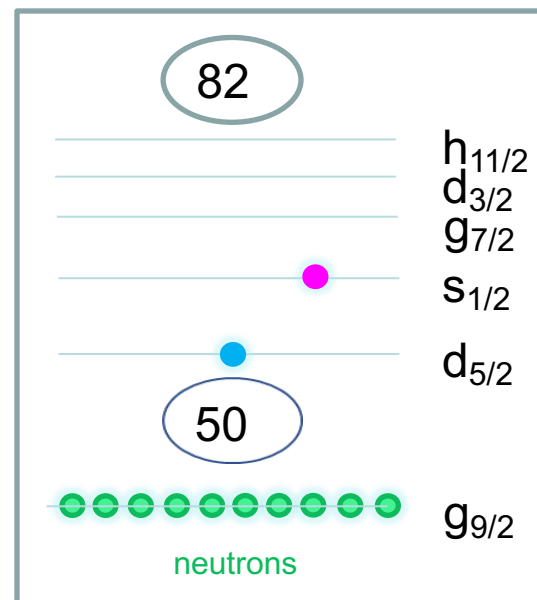


$E_x$ (MeV)	$J^\pi$	$\ell$	$C_{\ell j}^2$
0.000	5/2+	2	$6.11 \pm 1.43$
0.462	1/2+	0	$25.3 \pm 5.9$

- 4.5 MeV/u at HRIBF
- 45 MeV/u at NSCL

$$S = \left( \frac{d\sigma}{d\Omega} \right)_{exp} / \left( \frac{d\sigma}{d\Omega} \right)_{thy}$$

- $^{85}\text{Se}$  states:  $2d_{5/2}$  and  $3s_{1/2}$
- Probes different parts of wave function
  - Low energy = peripheral (only tail)
  - Higher energy = less peripheral (more interior)

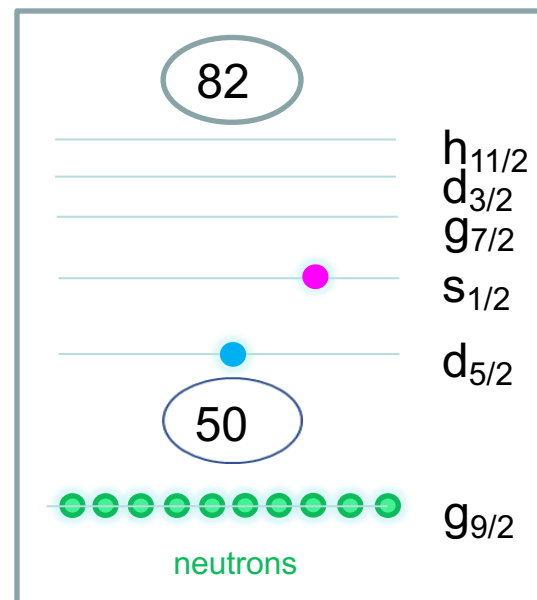
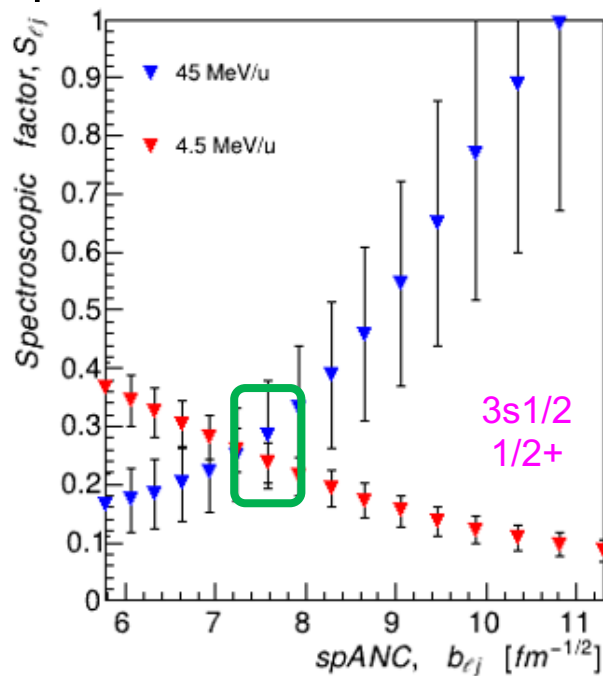
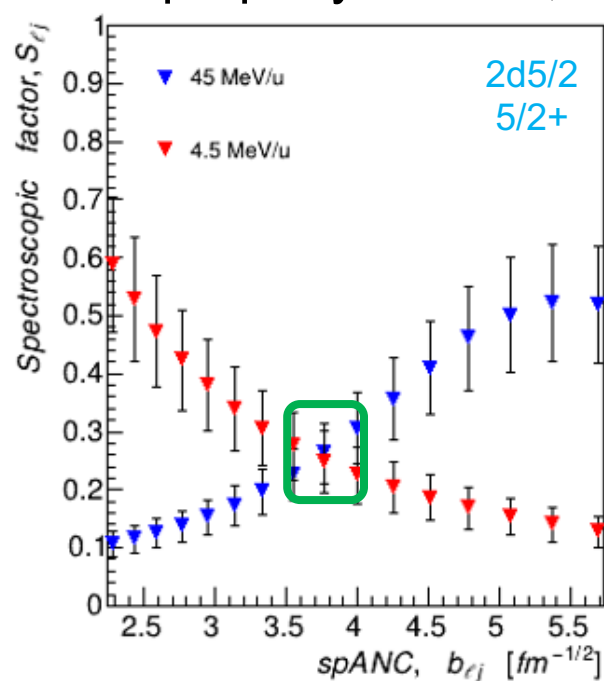


Excitations in  $^{85}\text{Se}$

- FR-ADWA cross sections with FRESKO
  - Includes deuteron b/up (Johnson-Tandy)
  - Global optical model potentials
    - Koning-Delaroche
  - Bound state parameters for the transferred neutron
    - $R = r_0 A^{1/3}$  diffuseness  $a$  Woods-Saxon potential
  - Wave function of transferred particle, e.g.,  $2d_{5/2}$  neutron

# RUTGERS (d,p) studies with 4.5 MeV/u & 45 MeV/u $^{84}\text{Se}$ beams

- 4.5 MeV/u at HRIBF
- 45 MeV/u at NSCL
- $S = \left( \frac{d\sigma}{d\Omega} \right)_{exp} / \left( \frac{d\sigma}{d\Omega} \right)_{thy}$
- spANC  $\leftrightarrow$  unknown Woods-Saxon potential ( $r_0, a$ )
- $S$  is property of state, independent of reaction



Excitations in  $^{85}\text{Se}$

$$C_{\ell j}^2 = S_{\ell j} b_{\ell j}^2$$

$5/2^+ E_x = 0$	$S = 0.28 (4)$	$r_0 = 1.14^{+0.04}_{-0.06}$ $a = 0.59^{+0.02}_{-0.03}$
$1/2^+ E_x = 0.462 \text{ MeV}$	$S = 0.26 (6)$	$r_0 = 1.16^{+0.06}_{-0.09}$ $a = 0.60^{+0.04}_{-0.04}$

H.E. Sims PhD Dissertation (2020)

H.E. Sims, D. Walter et al.,  
in preparation for PRC (2023)