Some aspects of the structure of neutron-rich F isotopes in the Particle-Rotor Model

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Summary



Evolution of shell structure and collectivity in exotic nuclei

A driving question in nuclear physics



Evolution of shell structure and collectivity in exotic nuclei

A driving question in nuclear physics



Island of inversion

Article Talk

ŻĄ 1 language ∨

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From Wikipedia, the free encyclopedia

An **island of inversion** is a region of the <u>chart of nuclides</u> where isotopes have enhanced stability in a sea of mostly fleeting and unstable nuclei at the edge of the nuclear map. Each island contains <u>isotopes</u> with a non-standard ordering of single particle levels in the <u>nuclear shell model</u>. Such an area was first described in 1975 by French physicists carrying out spectroscopic mass measurements of exotic isotopes of lithium and sodium.^[1] Since then further studies have shown that five such regions exist within the known table of nuclides. These are centered at neutron-rich isotopes of five elements, namely ¹¹Li, ²⁰C, ³¹Na, ⁴²Si, and ⁶⁴Cr.^[2] Because there are five known islands of inversion, physicists have suggested renaming the phenomenon as an "archipelago of islands of shell breaking".^[2] Studies with the purpose of defining the edges of this region are still ongoing.

Much evidence has been obtained for the existence of deformed ground states, and a good understanding of the physical mechanism behind the inversion.





Evolution of shell structure and collectivity in exotic nuclei

A driving question in nuclear physics

Evolution of shell structure in exotic nuclei

Takaharu Otsuka, Alexandra Gade, Olivier Sorlin, Toshio Suzuki, and Yutaka Utsuno Rev. Mod. Phys. **92**, 015002 – Published 27 March 2020



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PHYSICS LETTERS B

5 February 1987

THE ONSET OF DEFORMATION AT THE N = 20 NEUTRON SHELL CLOSURE FAR FROM STABILITY

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N=20 shell gap

 $\Delta \ell = \Delta j = 2$ $\rightarrow \text{Quadrupole}$ Correlations

Role of the π d_{5/2}- v d_{3/2} interaction

$$H = Esp + GP^+P + xQ.Q$$

A. Poves and J. Retamosa, Phys. Lett. B 184, 311 (1987).
E. K. Warburton, J. A. Becker, and B. A. Brown, Phys. Rev. C41, 1147 (1990).



The cases of ^{25,29}F



PHYSICAL REVIEW C 95, 041301(R) (2017)

Low-Z shore of the "island of inversion" and the reduced neutron magicity toward ²⁸O

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The two odd-even fluorine isotopes 27,29 F were studied via in-beam γ -ray spectroscopy at the RIKEN Radioactive Isotope Beam Factory. A secondary beam of 30 Ne was used to induce one-proton and one-proton-two-neutron removal reactions on carbon and polyethylene targets at midtarget energies of 228 MeV/*u*. Excited states were observed at 915(12) keV for 27 F and at 1080(18) keV for 29 F. Both were assigned a 1/2⁺ spin and parity. The low transition energy for 29 F largely disagrees with shell model predictions restricted to the *sd* model space. Calculations using effective interactions that include the neutron *pf* shell indicate that the *N* = 20 gap is quenched for 29 F, thus extending the "island of inversion" to isotopes with proton number *Z* = 9. Variations of the *N* = 20 gap further reveal a strong correlation to the $1/2^+_1$ level energy in 29 F and suggest a persistent reduced neutron gap for 28 O.

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Much evidence has been obtained for the existence of deformed ground states



How about the deformed shell model?



Particle plus Rotor Model 101



Nilsson levels





Structure of ²⁹F: Geometry





A.O. Macchiavelli, H. L. Crawford, et al. Physics Letters B 775 (2017) 160–162

Structure of ²⁹F: Geometry



State	Energy [MeV]	$\langle R \rangle$	E _{rot} [MeV]	$\langle I_Z \rangle$	$\langle \vec{I}\cdot \vec{j} angle / I $	$\langle \vec{R} \cdot \vec{j} \rangle / R $
5/2 +	0	0.53	0.34	0.11	2.78	-0.30
1/2+	1.08	1.46	2.00	0.5	1.47	-2.14
3/2 ⁺	2.2	2.01	2.54	-1.12	1.58	-2.18
9/2 ⁺	2.6	2.18	2.91	0.04	2.65	1.76
7/2 ⁺	3.2	2.09	2.71	0.60	2.27	0.12



A.O. Macchiavelli, H. L. Crawford, et al. Physics Letters B 775 (2017) 160–162

Structure of ²⁹F: PRM Solution



NP1712-RIBF164 (H.Crawford, Coulomb excitation of ²⁹F)



Structure of ²⁹F: Decoupled Band



Full *sd* PRM calculation

Pure decoupled $d_{5/2}$ band



Structure of ²⁵F and its effective ²⁴O core

PHYSICAL REVIEW LETTERS 124, 212502 (2020)

How Different is the Core of ²⁵F from ²⁴O_{g.s.}?

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The structure of a neutron-rich ²⁵F nucleus is investigated by a quasifree (p, 2p) knockout reaction at 270A MeV in inverse kinematics. The sum of spectroscopic factors of $\pi 0d_{5/2}$ orbital is found to be 1.0 ± 0.3 . However, the spectroscopic factor with residual ²⁴O nucleus being in the ground state is found to be only 0.36 ± 0.13 , while those in the excited state is 0.65 ± 0.25 . The result shows that the ²⁴O core of ²⁵F nucleus significantly differs from a free ²⁴O nucleus, and the core consists of ~35% ²⁴O_{g.s.} and ~65% excited ²⁴O. The result may infer that the addition of the $0d_{5/2}$ proton considerably changes neutron structure in ²⁵F from that in ²⁴O, which could be a possible mechanism responsible for the oxygen dripline anomaly.





Evidence for a doubly magic ²⁴O

PRL 102, 152501 (2009)

C.R. Hoffman^{a,*}, T. Baumann^b, D. Bazin^b, J. Brown^c, G. Christian^{b,d}, D.H. Denby^e, P.A. DeYoung^e, J.E. Finck^f, N. Frank^{b,d,1}, J. Hinnefeld^g, S. Mosby^h, W.A. Peters^{b,d,2}, W.F. Rogers^h, A. Schiller^{b,3}, A. Spyrou^b, M.J. Scott^f, S.L. Tabor^a, M. Thoennessen^{b,d}, P. Voss^f

PHYSICAL REVIEW LETTERS

One-Neutron Removal Measurement Reveals ²⁴O as a New Doubly Magic Nucleus

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H. Weick² and M. Winkler²

NSCL MoNA

week ending 17 APRIL 2009

²⁶F 85 MeV/A

GSI FRS

RIKEN

²⁴O(p,p')

²⁴O at 920 MeV/A

PRL 109, 022501 (2012)	PHYSICAL	REVIEW	LETTERS	13 JULY 2012

N = 16 Spherical Shell Closure in ²⁴O

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Structure of ²⁵F and its effective ²⁴O core

Channel	Mean [MeV]	Width [MeV]	$\sigma_{\rm exp}$ [µb]	$\sigma_{\mathrm{th}}~[\mu\mathrm{b}]$	$J^{\pi}_{ m th}$	S _{exp}	$S_{\rm th}({\rm USDB})$	$S_{\rm th}({ m SFO})$	$S_{\rm th}({\rm SPDF-MU})$
(²⁵ F, ²⁴ O)	-0.5(1.1)	4.8(1.3)	53(18)	149(24)	5/2+	0.36(13)	1.01	<mark>0.90</mark>	0.95
$({}^{25}F, {}^{23}O)$	6.5(1.4)	6.3(9)	81(26)	125(26)	$5/2^{+}$	0.65(25)	0.01	0.07	0.05
$(^{25}\text{F}, ^{22}\text{O})$	12.7(6)	7.6(6)	274(71)	80(24)	$1/2^{-}$	3.43(1.4)		2.19	

The $0d_{5/2}$ proton knockout from ²⁵F populates the ²⁴O ground state with a smaller probability than the ²⁴O excited states. This result indicates that the oxygen core of ²⁵F is considerably different from ²⁴O_{gs} and has a larger overlap with the excited states of ²⁴O. The change in the neutron-shell structure due to the $0d_{5/2}$ proton may be responsible for the small overlap between ²⁵F and ²⁴O_{gs}

A comparison with the shell model calculations indicates that the USDB, SFO, and SFPD-MU interactions are insufficient to reproduce the present results. A stronger tensor force or other mechanism such as the 3N force effects, or both, might be needed to explain the experimental results. More experimental and theoretical studies are necessary to clarify the mechanism for the change in the core of neutron-rich fluorine from the ground state of oxygen isotopes.



Letter

Core of ²⁵F studied by the ²⁵F(-p) proton-removal reaction

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The ⁹Be(²⁵F(5/2⁺), ²⁴O)X proton-removal reaction was studied at the NSCL using the S800 spectrometer. The experimental spectroscopic factor for the ground-state to ground-state transition indicates a substantial depletion of the proton $d_{5/2}$ strength compared to shell-model expectations, similar to the findings of an inverse-kinematics (p, 2p) measurement performed at RIBF. The ²⁵F to ²⁴O ground-states overlap is considerably less than anticipated if the core nucleons behaved as rigid, doubly-magic ²⁴O within ²⁵F. We interpret the new results within the framework of the Particle-Vibration Coupling (PVC) model, of a $d_{5/2}$ proton coupled to a quadrupole phonon of an effective core. This approach provides a good description of the experimental data, requiring an effective ²⁴O* core with a phonon energy of $\hbar\omega_2$ = 3.2 MeV and a $B(E2) \approx 2.7$ W.u. – softer and more collective than a bare ²⁴O. Both the Nilsson deformed mean field and the PVC models appear to capture the properties of the effective core of ²⁵F, suggesting that the additional proton polarizes ²⁴O in such a way that it becomes either slightly deformed or a quadrupole vibrator.

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PHYSICAL REVIEW C 102, 041301(R) (2020)

Rapid Communications

Core of ²⁵F in the rotational model

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²⁵F in the PRM



The effective ²⁴O core in ²⁵F can be interpreted as a slightly deformed rotor with:

 E_2^+ (core) ≈ 3.2 MeV and $\mathcal{E}_2 \approx 0.15$,



Experiment Zs. Vajta *et al*, heory v. C 89, 054323 (2014).



²⁵F in the PRM



Furthermore, in 26 F * the 1⁺ ground and 4⁺ isomeric states can be associated with the antiparallel and parallel couplings of the odd neutron in the $d_{3/2}$ Nilsson multiplet to the structure of 25 F.

The former, favored by the Gallagher-Moszkowski rule gives 1⁺ as the lowest state and the latter a 4⁺ as the bandhead of a doubly decoupled band.



* A. Lepailleur *et al.*, Phys. Rev. Lett. **110**, 082502 (2013)







Nilsson Spectroscopic Factors

Single nucleon knockout

$$S_{i,f}(j\ell,K) = (g\langle I_i j\Omega_{\nu} K_i | I_f K_f) C_{j,\ell} V_{\nu} | \phi_f | \phi_i \rangle)^2$$

= $\theta_{i,f}(j\ell,K)^2$

Nilsson amplitudes

BCS occupations

Core overlap

$$\frac{d\sigma}{d\Omega} = \sum_{j,\ell} S_{j,\ell} \times \sigma_{\ell}^{sp}$$

B. Elbek and P. Tjom, Advances in Nucl. Phys. Vol 3, 259 (1969)



Spectroscopic factors

$$\mathsf{PRM} \implies \psi_I = \sum_K \mathcal{A}_K | IK \rangle. \implies S_{i,f}(j\ell) = \left(\sum_K \mathcal{A}_K \theta_{i,f}(j\ell, K)\right)^2,$$

Final state in ²⁴ O	S _{exp} Ref. [7]	PRM1	S _{th} PRM2	SDPF-MU
Ground Excited	0.36(13) 0.65(25)	0.85 0.15	0.56 0.44	0.95 0.05
	No Quenching	$\langle \phi_f \phi_i angle pprox 0$.81	

Following: T. Takemasa, M. Sakagami, and M. Sano, Phys. Rev. Lett. **29**, 133 (1979).

T. Takemasa, Comput. Phys. Commun. 36, 79 (1985).



A Particle vibration coupling view



$$\Delta E(j,I) = \frac{h^2(j,j,2)}{\hbar\omega_2} \left(\delta_{jI} + (2j+1)\right) \begin{cases} 2 & j & j \\ 2 & j & I \end{cases}$$

$$h(j, j, 2) = \left(\frac{5}{4\pi}\right)^{1/2} \langle j\frac{1}{2}20|j\frac{1}{2}\rangle \left(\frac{\hbar\omega_2}{2C_2}\right)^{1/2} \langle j|\kappa_2(r)|j\rangle$$

$$B(E2, n = 0 \rightarrow n = 1) = 5\left(\frac{3}{4\pi}ZeR^2\right)^2\left(\frac{\hbar\omega_2}{2C_2}\right)$$

$$\widehat{|j\rangle} \approx a|j,n=0\rangle + b|j,n=1\rangle \qquad b = \frac{h(j,j,2)}{\hbar\omega_2}$$

$$C^2 S_{PVC} = a^2 \langle {}^{24} O | {}^{24} O^* \rangle^2 = 0.6$$

Parameter	$^{24}\mathrm{O}$	${}^{24}\mathrm{O*}$	$^{28}O^{*}$
$\hbar\omega_2$	4.7	3.2	2.0
h(j,j,2)	0.73	1.58	1.75
C_2	204	140	97
B(E2)	0.0012	0.0055	0.0082



A first look at ^{28,30}F

PHYSICAL REVIEW LETTERS 124, 152502 (2020)

Extending the Southern Shore of the Island of Inversion to ²⁸F

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(SAMURAI21 collaboration)

Detailed spectroscopy of the neutron-unbound nucleus ²⁸F has been performed for the first time following proton/neutron removal from ²⁹Ne/²⁹F beams at energies around 230 MeV/nucleon. The invariant-mass spectra were reconstructed for both the ²⁷F^(*) + *n* and ²⁶F^(*) + 2*n* coincidences and revealed a series of well-defined resonances. A near-threshold state was observed in both reactions and is identified as the ²⁸F ground state, with $S_n(^{28}F) = -199(6)$ keV, while analysis of the 2*n* decay channel allowed a considerably improved $S_n(^{27}F) = 1620(60)$ keV to be deduced. Comparison with shell-model predictions and eikonal-model reaction calculations have allowed spin-parity assignments to be proposed for some of the lower-lying levels of ²⁸F. Importantly, in the case of the ground state, the reconstructed ²⁷F + *n* momentum distribution following neutron removal from ²⁹F indicates that it arises mainly from the $1p_{3/2}$ neutron intruder configuration. This demonstrates that the island of inversion around N = 20 includes ²⁸F.

and most probably ²⁹F, and suggests that ²⁸O is not doubly magic.



A first look at ^{28,30}F







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Angular momentum structure



Angular momentum structure



Summary

The low-lying structure of 25,29 F can be understood in terms of the rotation-aligned coupling limit of the PRM. Coriolis coupling on the $d_{5/2}$ proton Nilsson multiplet gives rise to a decoupled band with a $5/2^+$ bandhead.

Calculated proton spectroscopic factors for the ${}^{25}F(5/2^+)(-1p)$ ${}^{24}O$ reaction are in agreement with the experimental data. The observed fragmentation of the $d_{5/2}$ strength is due to both deformation and a core overlap.

The Nilsson plus PRM picture suggests that the extra proton with a dominant component in the downsloping [220] ¹/₂ level polarizes ^{24,28}O and stabilizes its dynamic deformation. Thus, the effective core in ²⁵F (²⁹F) can be interpreted as a slightly deformed rotor with *E*2+ (core) \approx 3.2 MeV (2.5MeV) and $\varepsilon_2 \approx$ 0.15, compared to the real doubly magic ²⁴O (²⁸O) with *E*2+ \approx 4.7 MeV (??) and weak vibrational quadrupole collectivity.

Electromagnetic observables for the three lowest experimental levels, obtained in the PRM, suggest that Coulomb excitation experiments will shed further light on the validity of our interpretation.

Two-qp plus rotor model calculations of 28,30 F are in progress. Inclusion of residual V_{pn} next



Merci !

