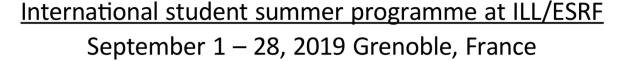


X-RAY AND NEUTRON SCIENCE





Principles of synchrotron radiation

Nicola Carmignani

ESRF, Beam Dynamics group, Accelerator and Source division With help from Boaz Nash, Andrea Franchi and Jean-Luc Revol



The European Synchrotron

OUTLINE

Properties of radiation

Where x-rays come from

review of relativity

bending magnet radiation

undulator radiation

Particle accelerators:

how to store electron beams

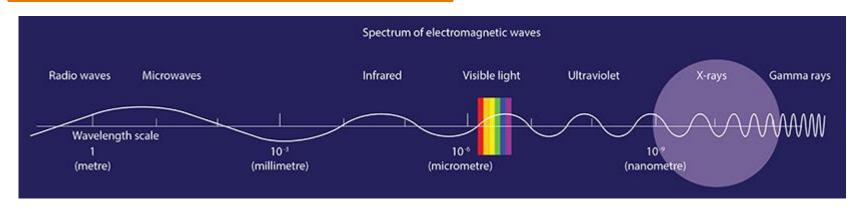
what is the distribution of the electron beam

ESRF upgrade: EBS



PROPERTIES OF RADIATION

Spectrum of electro-magnetic radiation



$$f = c / \lambda$$

 $E_{\gamma} = h f$

c = 299792458 m/s

 $h = 4.135667517 \times 10^{-15} \text{ eV s}$

 λ is the photon wavelength **f** is the photon frequency \mathbf{E}_{γ} is the photon energy c is the speed of light **h** is the Planck constant

 \mathbf{E}_{γ} for X-rays is about 1-100 keV



Nicola Carmignani

PROPERTIES OF RADIATION

Flux of radiation: number of photons per second

Brightness of radiation:

flux at a specific wavelength divided by source size and divergence

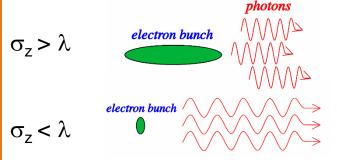


Low brightness: low photon density on sample



High brightness: high photon density on sample

Coherence of radiation: are the photons produced in phase or not?



 σ_{7} is the electron bunch length

Polarization of radiation: directionality of radiation field linear, circular partial/full polarization

Synchrotron light sources give some control over all these properties



WHERE X-RAYS COME FROM

Accelerating charged particles emit electro-magnetic radiation

Maxwell's equations:

$$abla \cdot \mathbf{E} = 4\pi
ho$$

$$\nabla \cdot {\bf B} = 0$$

$$abla extbf{x} extbf{E} = -rac{1}{c} rac{\partial extbf{B}}{\partial t}$$

$$abla extbf{X} extbf{B} = rac{1}{c} \left(4\pi extbf{J} + rac{\partial extbf{E}}{\partial t}
ight)$$

Non-relativistic Larmor formula:

$$P = \frac{q^2 a^2}{6\pi\varepsilon_0 c^3}$$

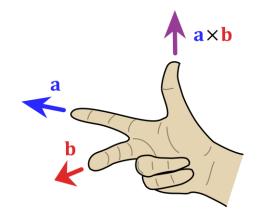
P is the power radiated q is the electric charge a is the acceleration

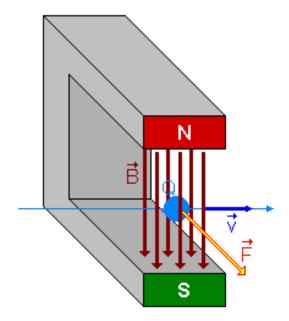


LORENTZ FORCE

We can move charged particles using the Lorentz force

$$ec{F} = q \cdot \left(ec{E} + ec{v} imes ec{B}
ight)$$





The force, and so the acceleration, is perpendicular to the particle velocity. Particles travel on an arc of circumference inside the dipole magnets.

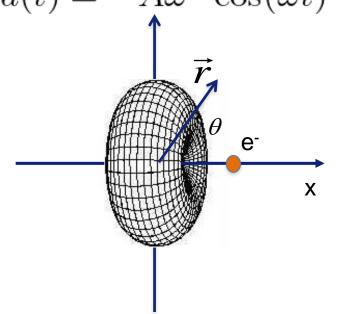
SINGLE PARTICLE IN A HARMONIC OSCILLATOR

Charge oscillating along x

$$x(t) = A\cos(\omega t)$$

$$v(t) = -A\omega \sin(\omega t)$$

$$a(t) = -A\omega^2 \cos(\omega t)$$



It emits radiation in all directions

$$\langle S \rangle \propto \frac{\sin^2 \theta}{r^2} \hat{r}$$

S is the Poynting vector (energy per unit of time and surface)

Frequency of radiation is

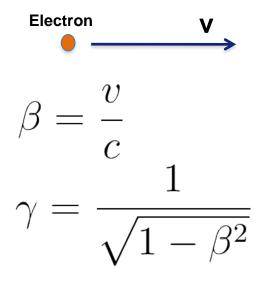
$$f = \frac{\omega}{2\pi}$$

Wavelength of radiation is

$$\lambda = \frac{c}{f}$$



BRIEF REVIEW OF RELATIVITY



c = 299792458 m/s is the speed of light

Total energy
$$E=\gamma mc^2$$

Kinetic energy
$$E_k = (\gamma - 1)mc^2$$

$$E_k \to \frac{1}{2}mv^2, \ for \ \beta << 1$$

For an electron with energy E:

$$\gamma = \frac{E}{mc^2} = \frac{E}{510998.9461 \,\text{eV}}$$

$$m_{e^{-}} = 510998.9461 \frac{\text{eV}}{c^2}$$



BRIEF REVIEW OF RELATIVITY

is always less than 1

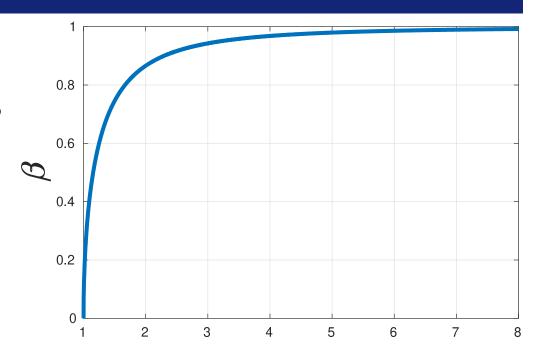
Velocity is almost c already for γ~5

At ESRF γ =11800

Increasing the electron energy from few MeV to 6 GeV does not change much the velocity.

However, there are two effects as gamma gets large:

- length contraction
- time dilation



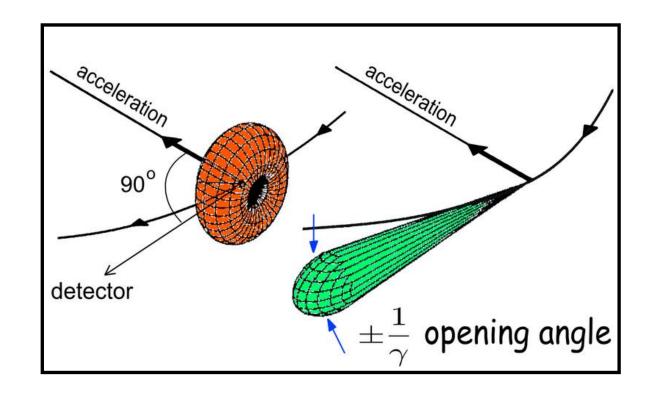
$$L = L_0 \sqrt{1 - \frac{v^2}{c^2}} = \frac{L_0}{\gamma}$$

$$T = \frac{T_0}{\sqrt{1 - \frac{v^2}{c^2}}} = T_0 \gamma$$



RADIATION EMITTED IN RELATIVISTIC CASE

If the charged particle is moving very fast, the radiation is emitted in a cone, with aperture $1/\gamma$

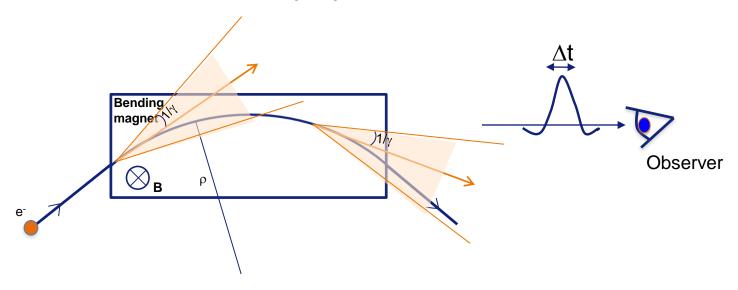




BENDING MAGNET RADIATION: TYPICAL ENERGY

What is the **typical energy** of the radiation from a bending magnet?

First, let's see for how long an observer would see the radiation produced by an ultra-relativistic electron in a bending magnet.



From the length Δt of the pulse we can estimate a typical energy E_{typ} :

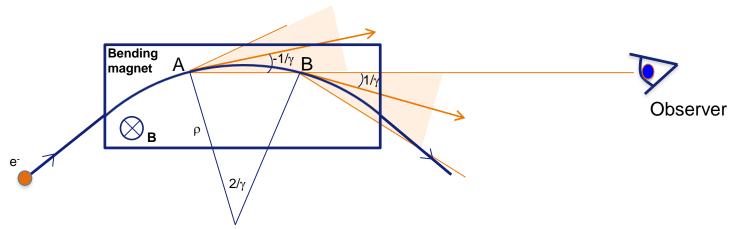
$$f_{typ} = 1/\Delta t$$

$$E_{typ} = h f_{typ}$$



BENDING MAGNET RADIATION: TYPICAL ENERGY

How long does the observer see the radiation produced by an ultra-relativistic electron in a bending magnet?



The observer will first see the radiation produced in the point A, with an angle -1/ γ . The last photon seen will be the one produced in B with an angle $\pm 1/\gamma$. The observer will therefore see radiation for an amount of time equal to the difference between the time needed by the electron to go from A to B and the time needed by the photon to go from A to B.

$$\Delta t = t_{ABe^{-}} - t_{AB\gamma} = \frac{\overline{AB}}{\beta c} - \frac{\overline{AB}}{c} = \frac{2\rho}{\gamma \beta c} - \frac{2\rho \sin(1/\gamma)}{c}$$

$$\sin(\theta) \simeq \theta - \frac{\theta^{3}}{6}, \quad for \, \theta << 1 \qquad \gamma >> 1, \quad 1 - \beta \simeq \frac{1}{2\gamma^{2}}$$

$$\Delta t \simeq \frac{2\rho}{\beta \gamma c} - \frac{2\rho}{c} \left(\frac{1}{\gamma} - \frac{1}{6\gamma^{3}}\right) = \frac{2\rho}{\beta \gamma c} \left(1 - \beta + \frac{\beta}{6\gamma^{2}}\right) = \frac{4}{3} \frac{\rho}{c\gamma^{3}}$$

$$\omega_{typ} = \frac{2\pi}{\Delta t} = \frac{3\pi c\gamma^3}{2\rho}$$

$$E_{typ} = \hbar \omega_{typ} = \frac{3\pi}{2} \frac{\hbar c}{\rho} \gamma^3$$

BENDING MAGNET RADIATION

Power emitted by an ultra-relativistic particle in a bending magnet

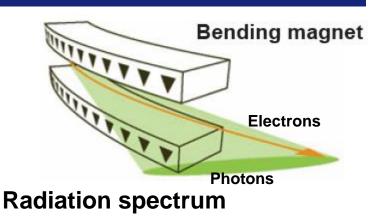
$$P = \frac{1}{6} \frac{e^2 c}{\pi \varepsilon_0} \frac{1}{\rho^2} \left(\frac{E}{mc^2} \right)^4$$

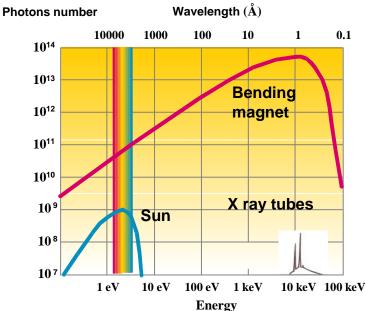
Bending magnet

ρ is the radius of curvature

Critical photon energy: half of the power is radiated at larger energy and half at lower

$$\varepsilon_c = \frac{3}{2} \frac{\hbar c}{\rho} \gamma^3$$





BENDING MAGNET RADIATION

Parameters for the EBS bending magnet beamlines

F = 6 GeV

 $\gamma = 11800$

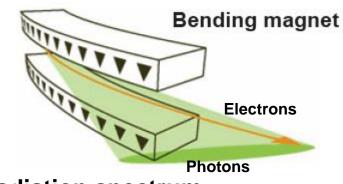
 $\rho = 23.3 \text{ m}$

 $\varepsilon_c = 21 \text{ keV}$

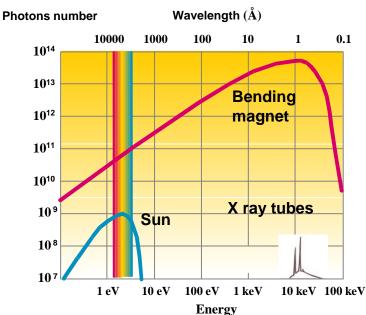
If we put protons instead of electrons?

 $m_p = 938 \text{ MeV/c2}$ $m_e = 0.511 \text{ MeV/c2}$

Radiated power would be 10¹³ times smaller $\varepsilon_c = 3.3 \,\mu\text{eV}$

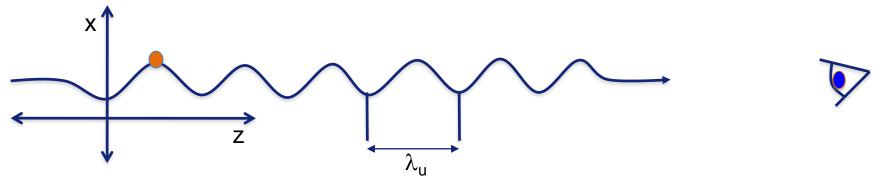


Radiation spectrum

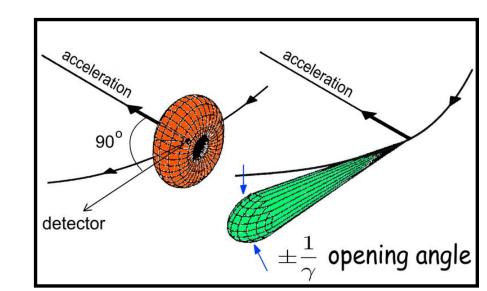


UNDULATOR RADIATION

Consider now a single charge oscillating on x and also moving along z direction

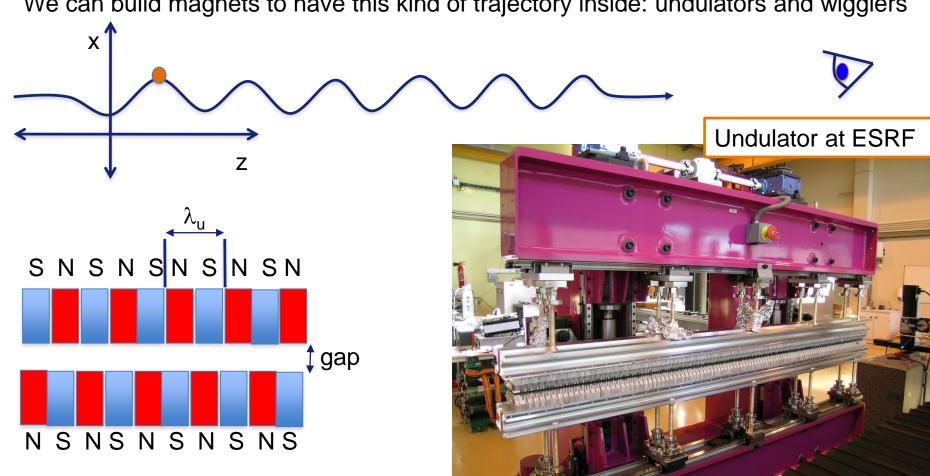


- 1. The wavelength is shifted by Doppler effect and, in case of relativistic speed, there is a time dilation effect. $\lambda = \frac{\lambda_u}{\lambda_u^2}$
- 2. Pattern of radiation gets distorted from motion. For high energy gets bent into cone of angle $1/\gamma$.



UNDULATOR RADIATION

We can build magnets to have this kind of trajectory inside: undulators and wigglers

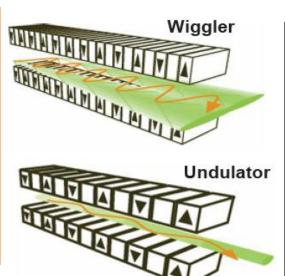


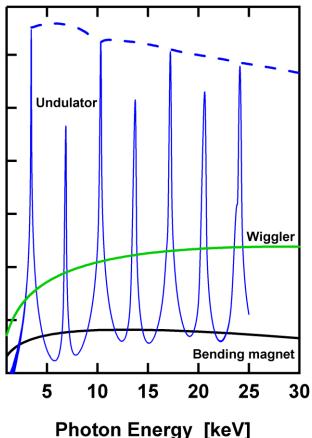
UNDULATOR RADIATION

Undulator parameter K

$$K = \frac{eB_0\lambda_u}{2\pi mc}$$

If K is large (K>>1), the device is no longer called undulator, but wiggler. K can be changed by varying the gap of the undulator.



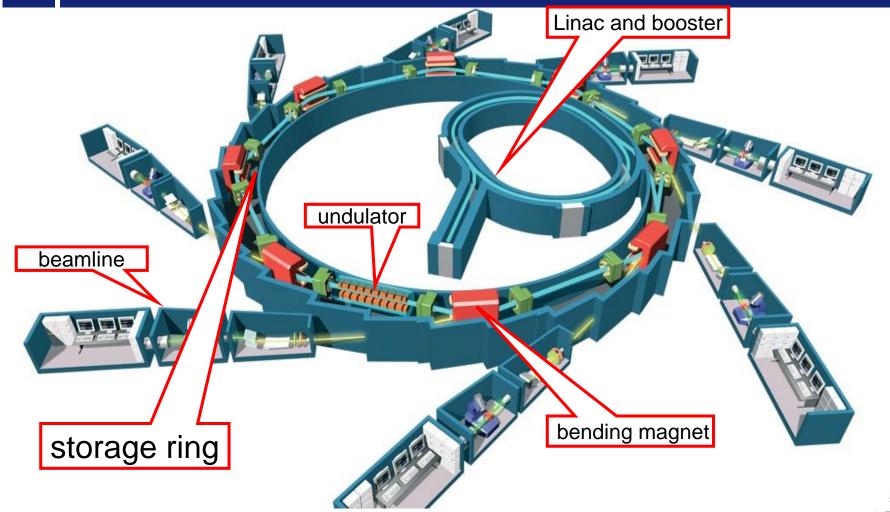


Fundamental undulator wavelength λ₁

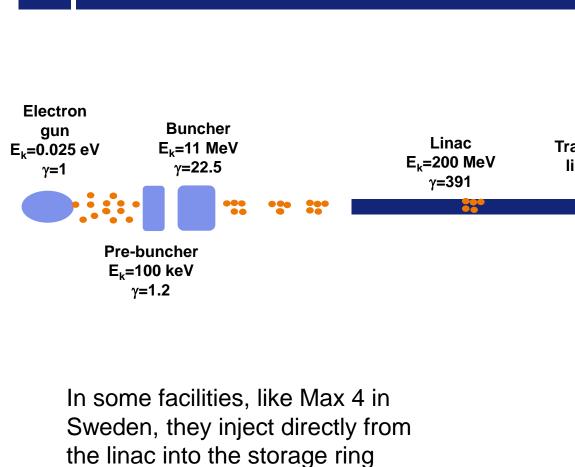
$$\lambda_1(\theta) \approx \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + (\gamma\theta)^2 \right)$$



A TYPICAL SYNCHROTRON LIGHT SOURCE USER FACILITY



ACCELERATION PROCESS

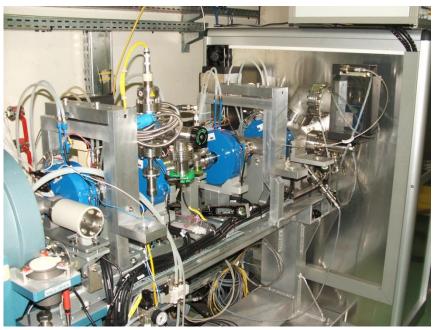


Storage ring E_k=6 GeV $\gamma = 11800$ **Transfer** line 1 **Booster** E_k=6 GeV $\gamma = 11800$ **Transfer** line 2 The European Synchrotron

Nicola Carmignani

INJECTORS

Electron gun, pre-buncher



Linac



Nicola Carmignani

INJECTORS: TRANSFER LINES

Linac to booster transfer line (TL1)





INJECTORS

Booster







128 at ESRF

Bending magnets (dipoles)

have uniform constant vertical magnetic field

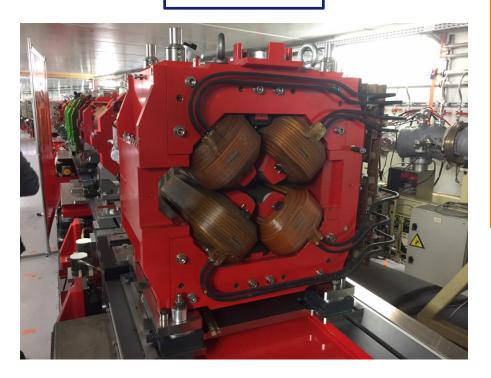
 $B_y = B_0$

They bend the beam and they define the circular trajectory.

At ESRF, each of the 5 pieces have a different magnetic field.



512 at ESRF



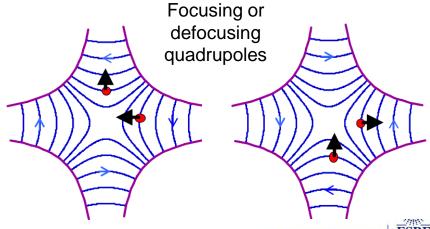
Quadrupoles:

magnetic field is linear with the distance from the center

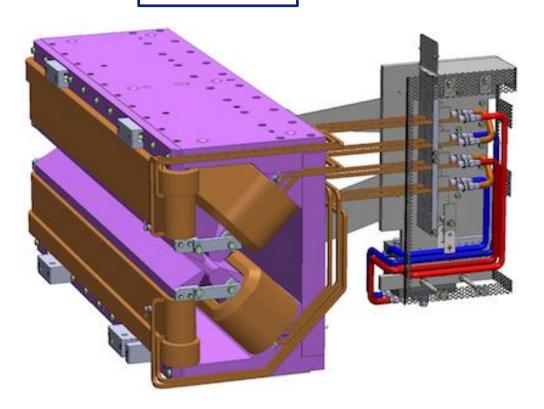
 $B_y = K_1 x$

They are used to focus the beam.

There are focusing or defocusing quadrupoles: if they focus horizontally, they defocus in vertical and vice versa (like astigmatic lenses in optics).



96 at ESRF



Dipole-Quadrupoles (DQ):

magnetic field is the one of a dipole plus the one of a quadrupole:

 $B_y = B_0 + K_1 x$

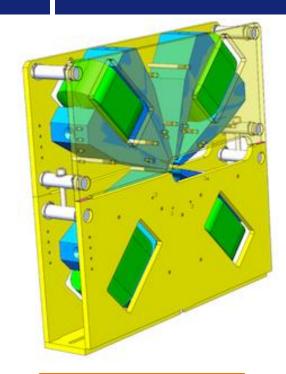
They are used to bend and focus the beam.

192 at ESRF



Sextupoles: magnetic field is quadratic with the distance from the center $B_v = K_2 x^2$ They are used to correct chromatic effects.

If correctly placed, they focus more the particles with higher energy and less the particles with lower energy. They correct chromatic effects.



Octupoles:

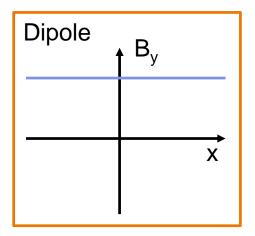
the field is cubic with distance: $By=K_3x^3$

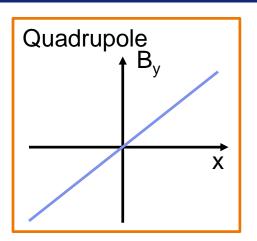


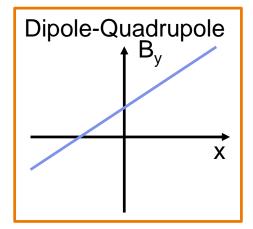
Correctors:

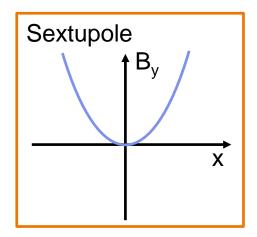
they can provide horizontal dipole field, vertical dipole field, quadrupolar field.

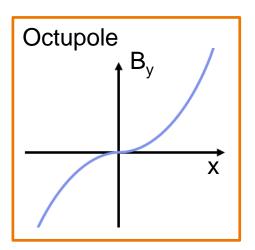




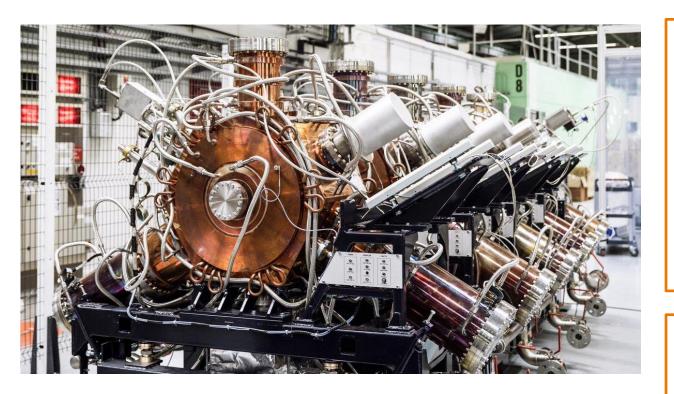








HOW DOES A STORAGE RING WORK: ACCELERATING RF CAVITIES



Most of the ESRF energy use (around 1.5 MW of power) is in these cavities

Radio Frequency cavities

They give to the beam the energy lost in synchrotron radiation and they provide longitudinal focusing.

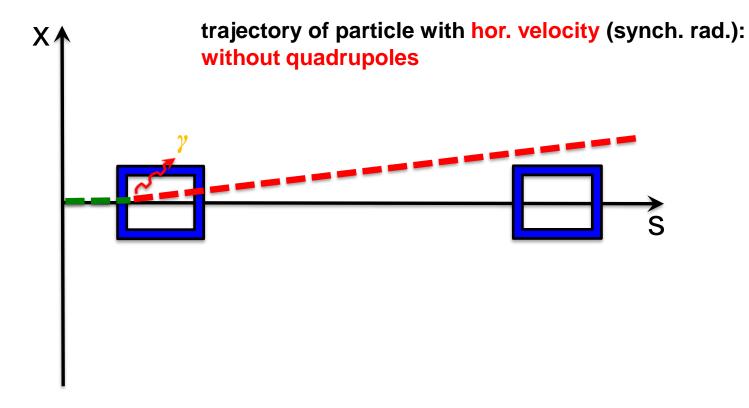
They produce an oscillating longitudinal electric field $V(t)=V_{RF}\sin(h\omega_0 t)$ $\omega_0=2\pi~f_{rev}$

At ESRF:

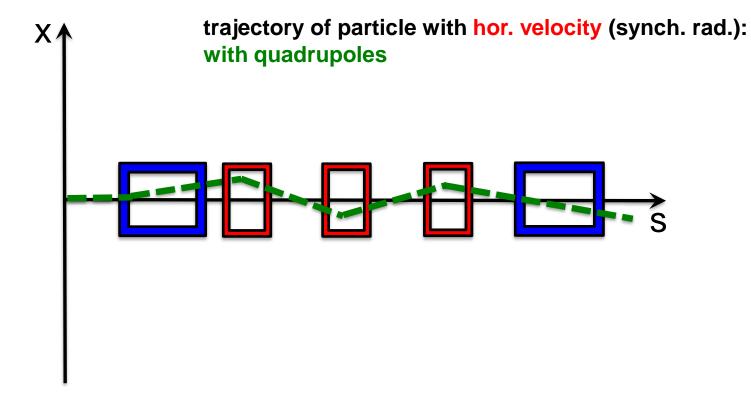
$$f_{rev}$$
 = 355 kHz
h=992
 f_{RF} =h f_{rev} = 352 MHz
 V_{RF} =6.5 MV



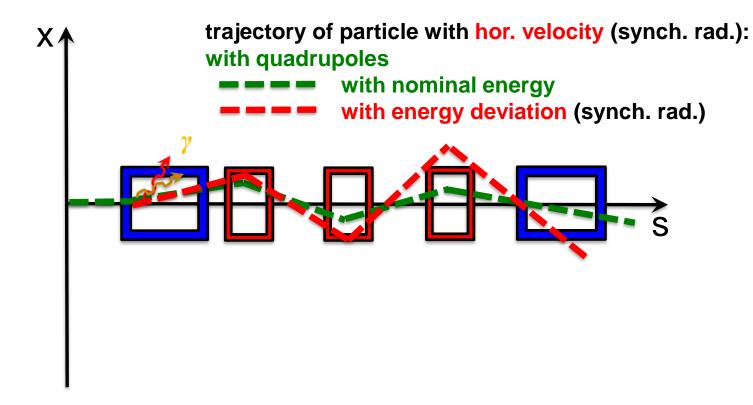
HOW DOES A STORAGE RING WORK: DIPOLES



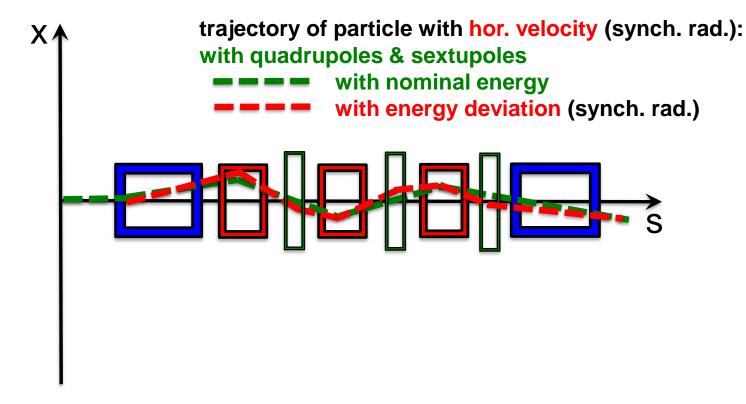
HOW DOES A STORAGE RING WORK: DIPOLES, QUADRUPOLES

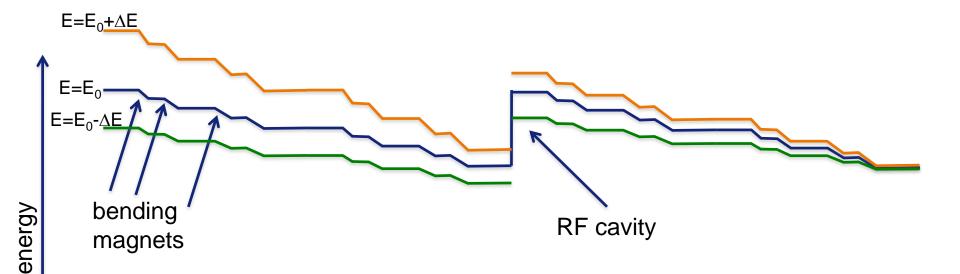


HOW DOES A STORAGE RING WORK: DIPOLES, QUADRUPOLES



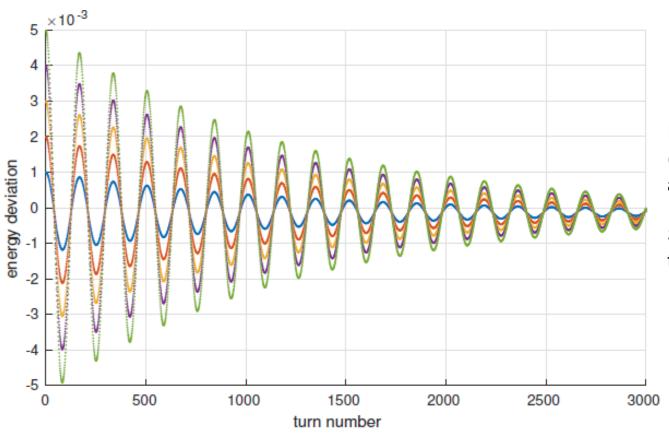
HOW DOES A STORAGE RING WORK: DIPOLES, QUADRUPOLES, SEXTUPOLES





Higher energy radiates more, lower energy less. Causes damping towards $P = \frac{1}{6} \frac{e^2 c}{\pi \varepsilon_0} \frac{1}{\rho^2} \left(\frac{E}{mc^2}\right)^4$ reference energy. reference energy.

$$P = \frac{1}{6} \frac{e^2 c}{\pi \varepsilon_0} \frac{1}{\rho^2} \left(\frac{E}{mc^2} \right)^4$$



Energy oscillations (synchrotron oscillations) and radiation damping.

So what defines the size of the electron beam?



Question: if we have radiation damping, why isn't the electron beam size zero?

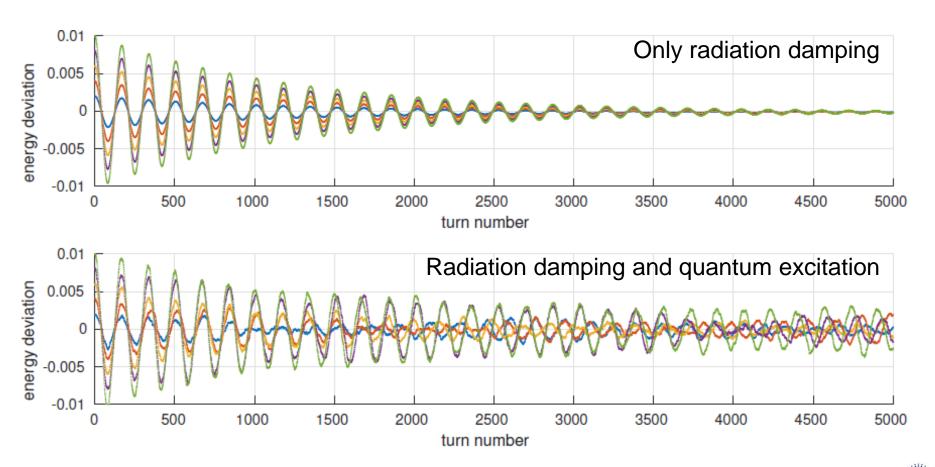
Answer: because of the graininess of the photon emission



Two sources of randomness:

- emission time of photons: Poisson process
- energy of the photon emitted

This causes a diffusion process in energy, that is converted to a transverse diffusion by the dipoles (higher energy electrons are bended less than lower energy ones). For ESRF, each electron emits only about 800 photons per turn, about 1 photon emitted per meter! (approx. 12 photons per dipole) Energy loss per turn $U_0 = 4.88 \text{ MeV}$



Results of damping and diffusion

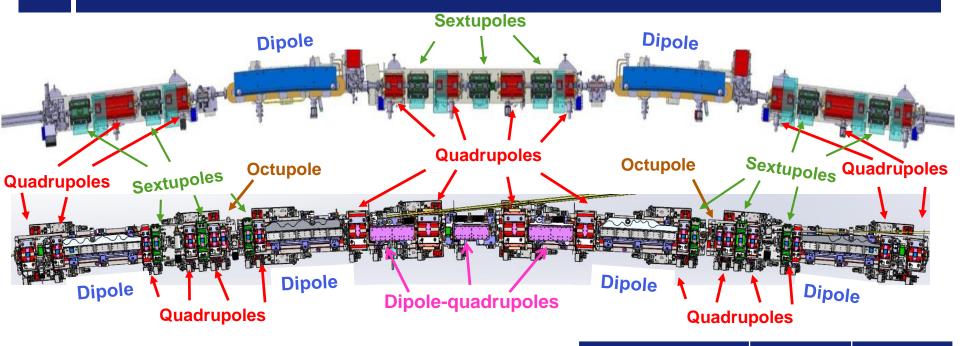
The electron beam reaches a unique Gaussian distribution, independent of how one injects into the ring.

> This is a major difference between electron synchrotrons and proton synchrotrons (e.g. LHC)

By careful choice of where the dipoles and quadrupoles are, one can reduce the size of this equilibrium beam size (emittance = beam size in phase space). So called "low emittance ring design".

> In fact, due to developments in lattice design, ESRF has completely replaced the storage ring in 2019 to reduce the electron beam emittance. 4nm rad -> 130 pm rad

ESRF EBS UPGRADE



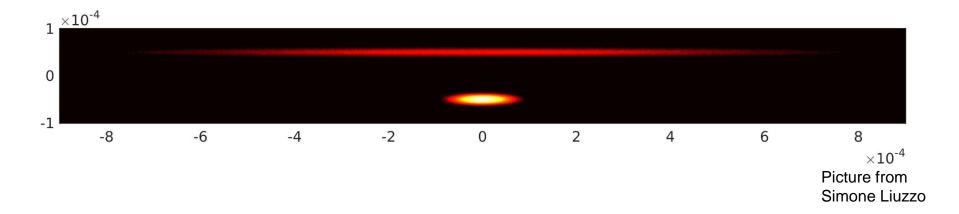
The ESRF Extremely Brilliant Source (EBS) will have 31 magnets per cell instead of 17 Free space between magnets (total for one cell):

3.4m instead of 8m today!

Magnet type	ESRF	ESRF EBS
Dipoles	64	128
Dipole-Quadrupoles	0	96
Quadrupoles	256	512
Sextupoles	224	192
Octupoles	0	64
Correctors	0	96

ESRF EBS UPGRADE

x-y beam distribution for old and new machine.



Electron beam horizontal emittance will be reduced by a factor 30 and brightness will then increase by the same factor.

$$\epsilon_{\text{ESRF}}$$
 = 4 nm rad ϵ_{EBS} = 130 pm rad



REFERENCES

Electromagnetism:

- J. D. Jackson, "Classical Electrodynamics", Wiley.
- R. P. Feynman, "The Feynman Lectures on Physics, Volume 2", www.feynmanlectures.caltech.edu.

Synchrotron radiation:

- A. Hofmann, "Characteristics of Synchrotron Radiation", proceedings of the CERN accelerator school 1998.
- P. Elleaume, "Undulators and wigglers", Taylor & Francis.

Particle accelerator physics:

- H. Wiedemann, "Particle accelerator physics", Springer.
- A. Chao, "Handbook of Accelerator Physics and Engineering", World scientific.
- S. Y. Lee, "Accelerator physics", World scientific.
- M. Sands, "The physics of electron storage rings, an introduction", SLAC report no. 121.

ESRF EBS upgrade:

The orange book, ESRF upgrade programme phase I (2015-2022) Technical Design Study.



MANY THANKS FOR YOUR ATTENTION

