

Solution X-ray Scattering from Biological Macromolecules

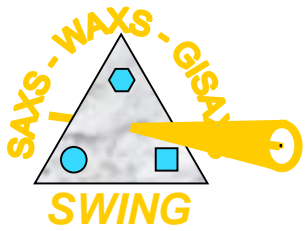
Javier Pérez

Beamline SWING, Synchrotron SOLEIL, Saint-Aubin, France

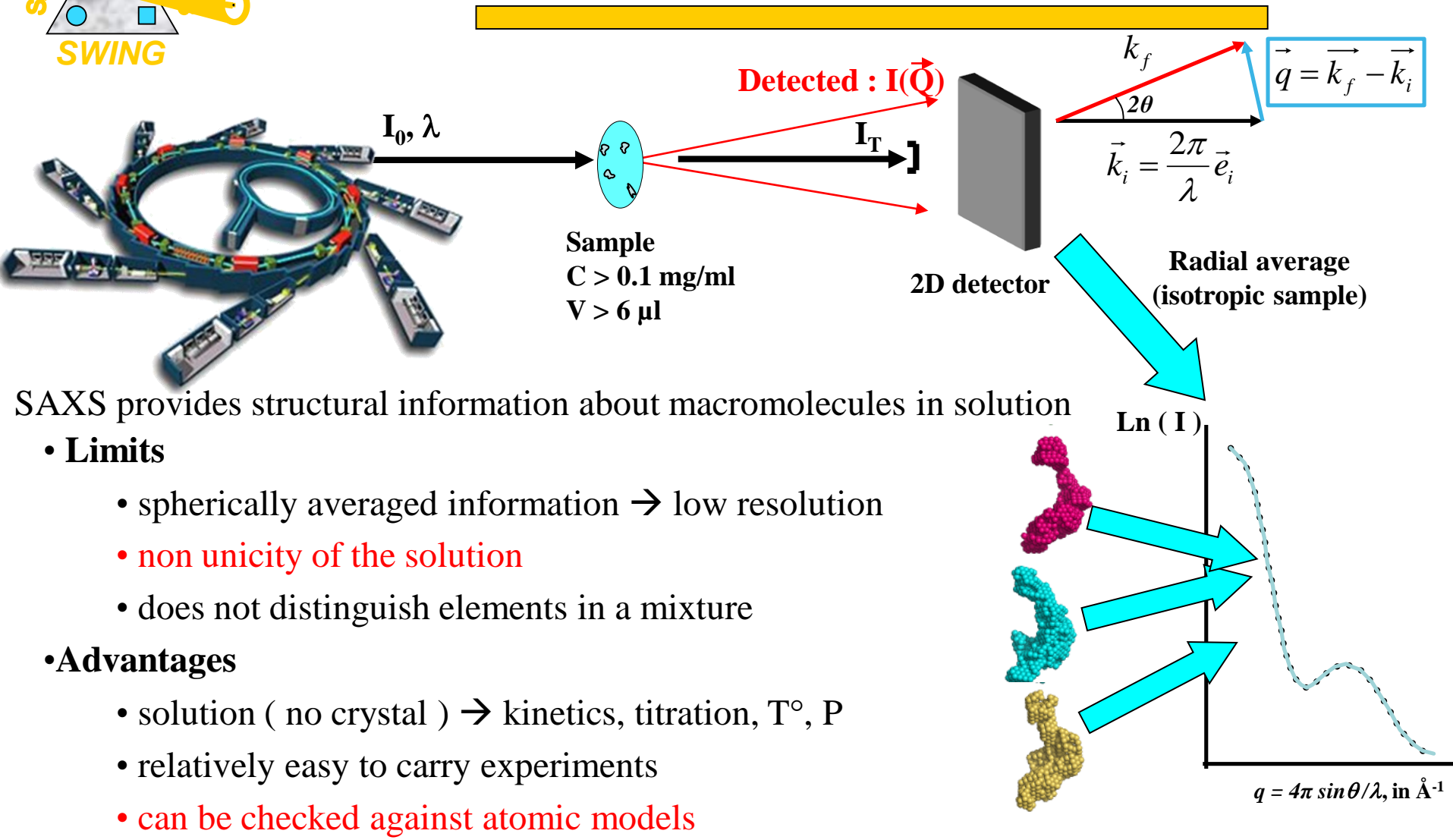
General Outline

- Introduction
- SAXS basics
- Biophysical information
- A few experimental considerations
- Modelling
- Conclusions

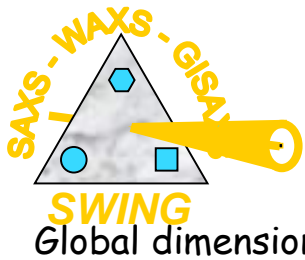
INTRODUCTION



Principles of Small Angle X-ray Scattering in solution

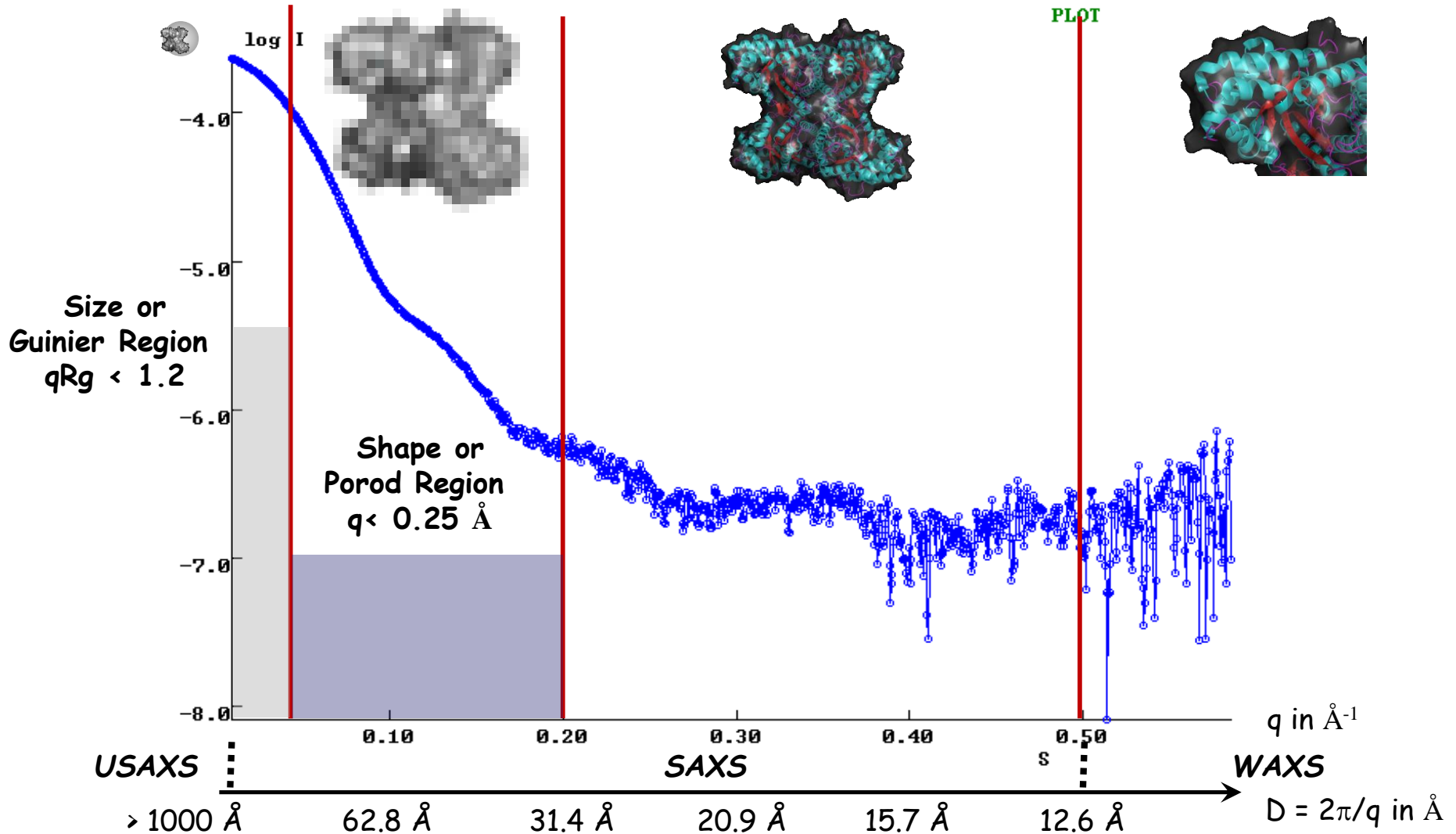


SAXS is at its best when complementary (structural) information is available

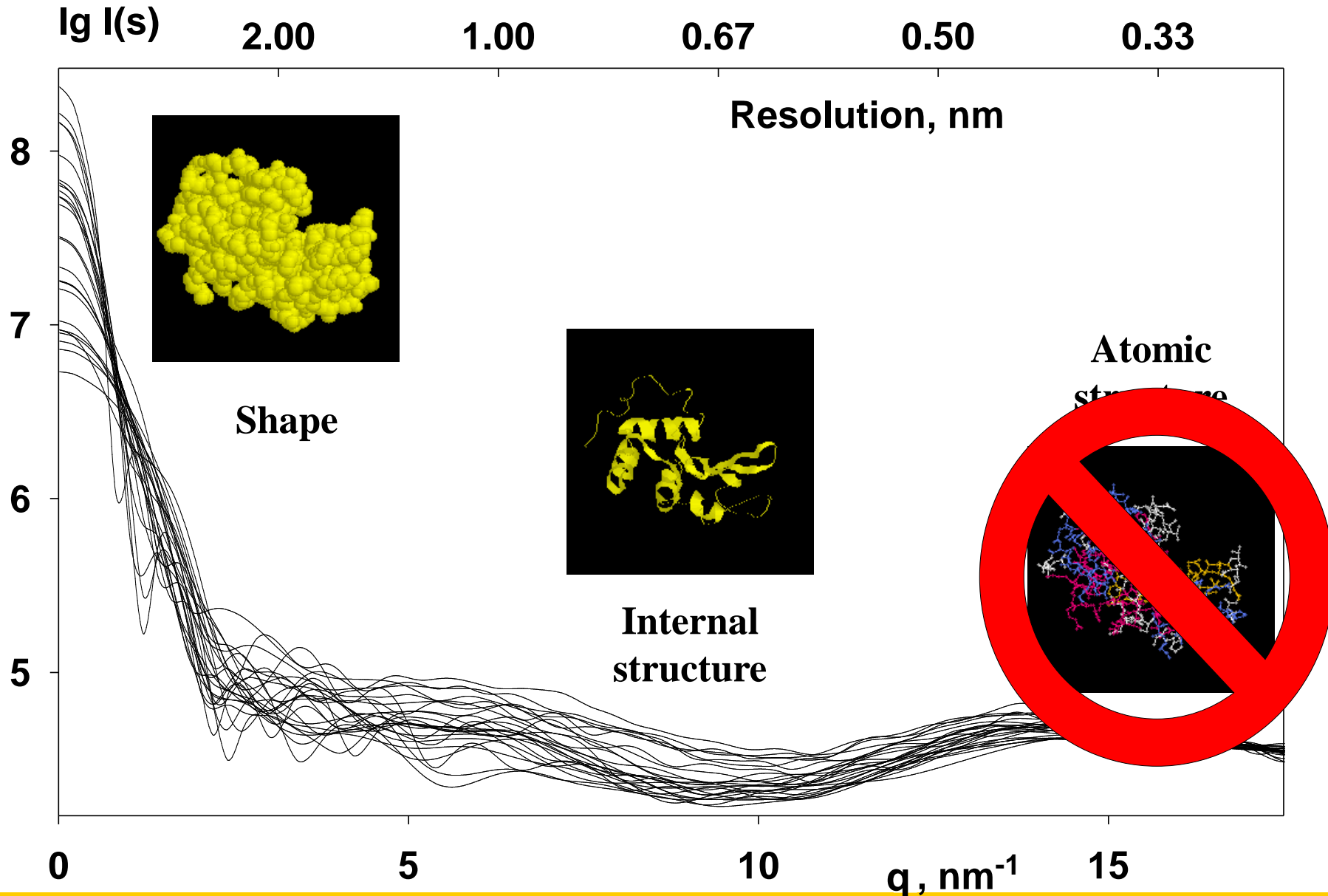


A typical BioSAXS curve from a protein

→ Global shape → Folding + domains → Secondary structures



What may solution scattering yield?



BioSAXS data output

Data analysis

Guinier range analysis

- R_g (size) and $I(0)$ (mass and oligomeric state)

Pair distribution function $p(r)$:

- D_{max} evaluation
- R_g (size) and $I(0)$ compatibility with Guinier approximation
- Global shape of the object

Kratky plot

- type of structure (globular, elongated or unfolded)

Porod law

- molecular volume if globular protein

Molecular modeling

Cristallographic , NMR structures or complete molecular modeling, [AlphaFold models](#)

- theoretical curves calculation and data comparison

Nothing is known

- low resolution shape

Structures of subunits available

- molecular modeling rigid body against SAXS data

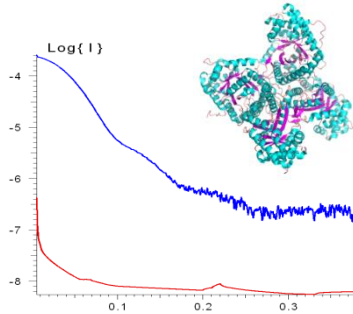
Structures with missing loop or flexible parts

- molecular modeling of missing parts against SAXS data

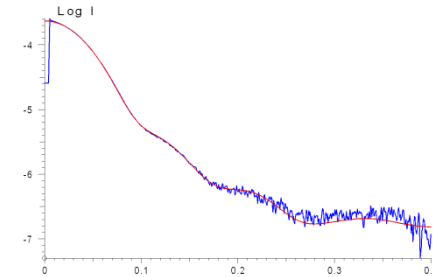
→ « **data compatible** » models: **NOT unique, NOT electronic density map**

BioSAXS data modeling, possible programs

1) Theoretical model or complete atomic structure available

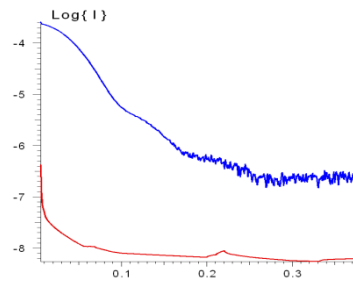


Validation/identification in solution

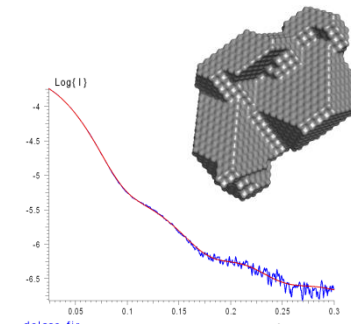


CRY SOL
FOXS
PepsiSAXS
WAXSiS

2) Nothing known (except the curve)

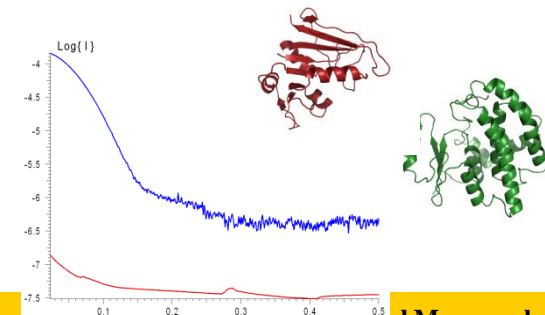


Low resolution model



DAMMIN
DAMMIF
GASBOR
MONSA
DENFERT

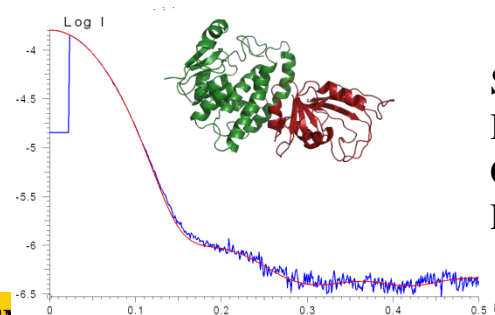
3) Structures of subunits available



Rigid body modeling of the complex and



molecular modeling of the missing part



SASREF
BUNCH
CORAL
DADIMODO

SAXS BASICS

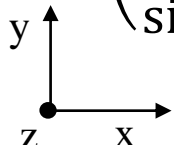
Elastic Thompson scattering by an electron

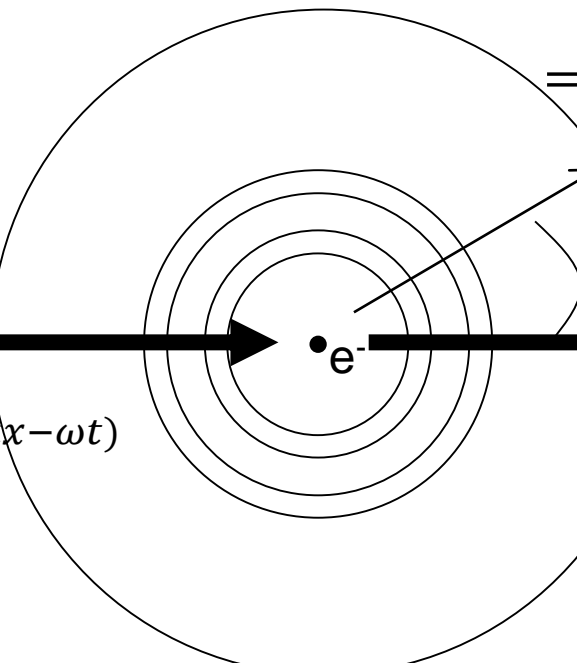
- What scatter X-rays are the electrons

X-ray incident beam

$$k = 2\pi / \lambda$$

$$\text{Flux } I_0 = S \cdot E_0^2$$

$$\vec{E}_i(x, t) = \begin{pmatrix} 0 \\ \cos \alpha \\ \sin \alpha \end{pmatrix} E_0 e^{i(kx - \omega t)}$$


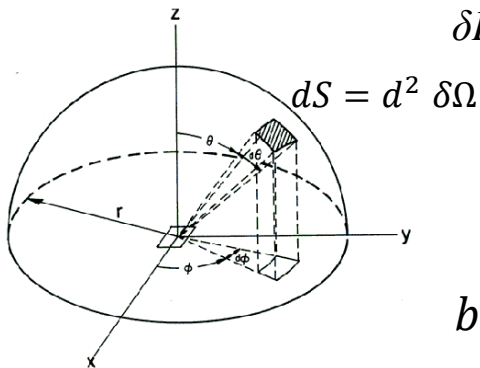


$$\vec{E}_s(d, t) = -\frac{r_0}{d} \begin{pmatrix} 0 \\ \cos \alpha \cos 2\theta \\ \sin \alpha \end{pmatrix} E_0 e^{i(kd - \omega t)}$$

$$r_0 = \frac{e^2}{4\pi\epsilon_0 mc^2} = 2,818 \cdot 10^{-15} m$$

r_0 is the electron classical radius

δI : flux scattered by one e^- through a solid angle $d\Omega$



$$\delta I_s = d^2 \delta\Omega \|\vec{E}_s(r, t)\|^2 = \delta\Omega b_0 E_0^2 * P(\alpha)$$

$$b_0 = \frac{1}{E_0^2} \frac{\delta I_s}{\delta\Omega} = r_0^2 \left(\frac{1 + \cos^2(2\theta)}{2} \right)$$

b_0 is the electron differential scattering cross section

Scattering amplitude by a particle

Coherent scattering : summing up amplitudes

- Waves scattered by two electrons

Electron 1

$$\vec{E}_{s1}(d, t) = -\frac{r_0}{d} \vec{E}_0 e^{i(kd - \omega t)}$$

Electron 2

$$\vec{E}_{s2}(d, t) = -\frac{r_0}{d} \vec{E}_0 e^{i(kd - \omega t + \vec{k}_i \cdot \vec{r} - \vec{k}_s \cdot \vec{r})}$$

Path shift between waves 1 and 2 :

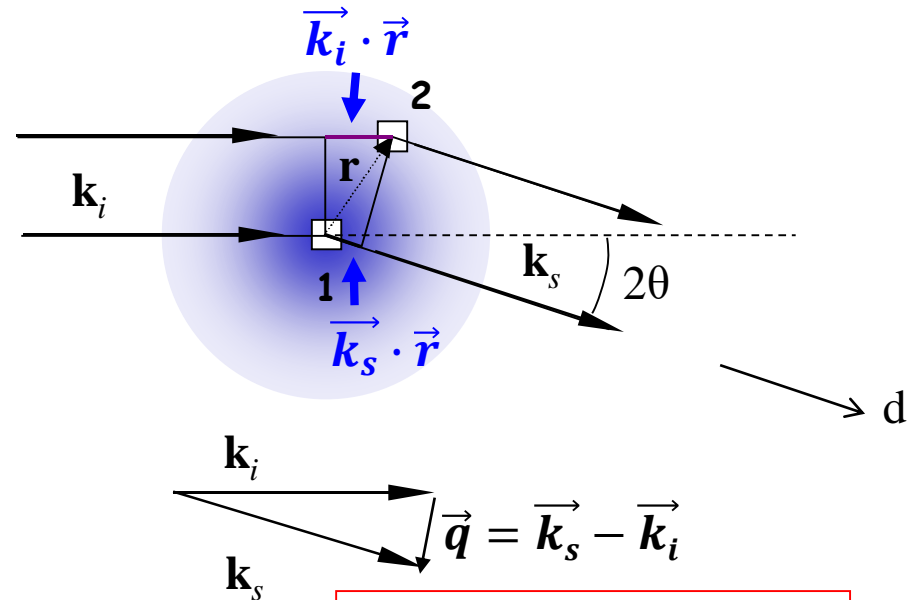
$$\vec{k}_i \cdot \vec{r} - \vec{k}_s \cdot \vec{r} = (\vec{k}_i - \vec{k}_s) \cdot \vec{r} = -\vec{q} \cdot \vec{r}$$

$$\vec{E}_{s2}(d, t) = -r_0 e^{-i\vec{q} \cdot \vec{r}} \frac{1}{d} \vec{E}_0 e^{i(kd - \omega t)}$$

- The scattered wave by a particle is simply the coherent sum of the spherical waves scattered by all the electrons of the particle.

Particle scattering « amplitude » (length)

$$A(\vec{q}) = -r_0 \int_V \rho_e(\vec{r}) e^{-i\vec{q} \cdot \vec{r}} d^3 \mathbf{r}$$



\vec{q} = Momentum transfer

$$q = \|\vec{q}\| = \frac{4\pi \sin(\theta)}{\lambda}$$

Intensity scattered by a sample – Auto-correlation function

Scattering amplitude $A(\vec{q}) = -r_0 \int_V \rho_e(\vec{r}) e^{-i\vec{q} \cdot \vec{r}} d^3\mathbf{r}$

Scattering intensity per unit volume : $I(\mathbf{Q})$, usual unit: cm^{-1} .

$$I(\vec{q}) = \frac{1}{V} A \cdot A^*(\vec{q}) = \frac{r_0^2}{V} \int_V \int_V \rho_e(\vec{r}_1) e^{-i\vec{q} \cdot \vec{r}_1} \rho_e(\vec{r}_2) e^{+i\vec{q} \cdot \vec{r}_2} d^3\mathbf{r}_1 d^3\mathbf{r}_2$$

$$I(\vec{q}) = \frac{r_0^2}{V} \int_V \int_V \rho_e(\vec{r}_1) \rho_e(\vec{r}_2) e^{-i\vec{q} \cdot (\vec{r}_1 - \vec{r}_2)} d^3\mathbf{r}_1 d^3\mathbf{r}_2$$

Auto-correlation function $\gamma_e(\mathbf{r})$:

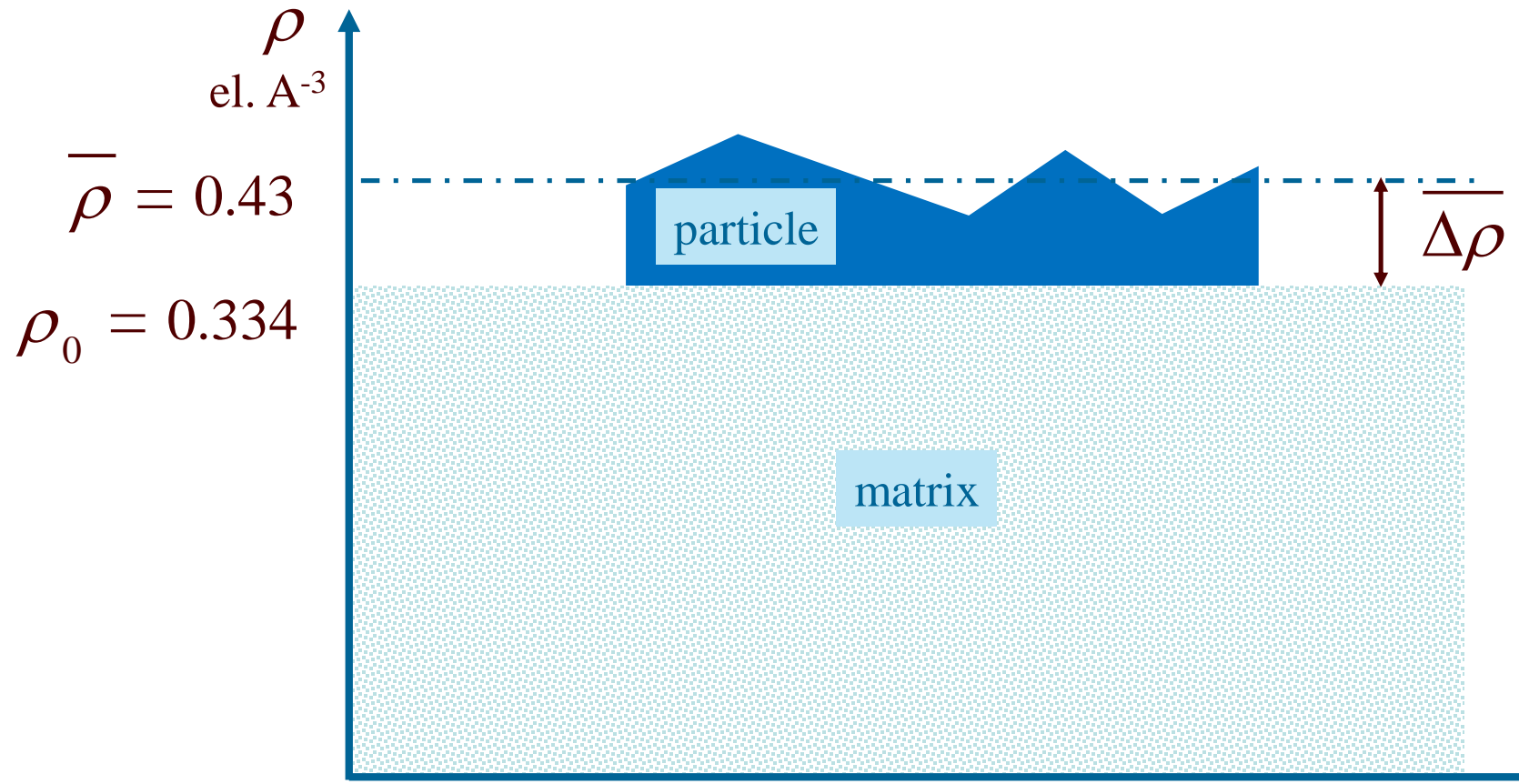
$$\gamma_e(\vec{r}) = \frac{1}{V} \int \rho_e(\vec{r}') \rho_e(\vec{r} + \vec{r}') d^3\mathbf{r}'$$

$$I(\vec{q}) = r_0^2 \int_V \gamma_e(\vec{r}) e^{-i\vec{q} \cdot \vec{r}} d^3\mathbf{r}$$

The scattered intensity is the Fourier Transform of the electronic density auto-correlation function

Particles in a homogeneous matrix (or buffer)

- A particle is described by the associated electron density distribution $\rho_p(\mathbf{r})$.
- In a matrix, what contributes to scattering is the **contrast** of electron density between the particle and the matrix $\Delta\rho(\mathbf{r}) = \rho_p(\mathbf{r}) - \rho_0$ that may be **very small** for biological samples.



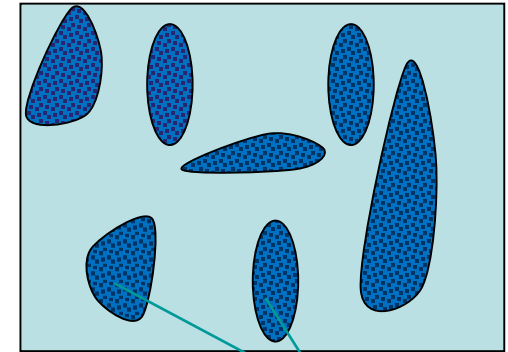
Particles inserted in a "homogeneous matrix"

- Scattering amplitude

$$f(\vec{q}) = -r_0 \int_{V_1} \Delta\rho(\vec{r}) e^{-i\vec{q}\cdot\vec{r}} d^3\mathbf{r}, \vec{q} \neq 0$$

Particles volume

Electronic Density Contrast



Particles

- $\Delta\rho(\vec{r})$ is the contrast of electronic density and describes the scattering object
- $f(\vec{q})$ is the Scattering Amplitude of the ensemble of the particles

- Scattering intensity per unit volume

$$I(\vec{q}) = \frac{1}{V} f(\vec{q}) f^*(\vec{q})$$

Irradiated volume

- $I(\vec{q})$ is expressed in cm^{-1} and is directly related to the measured intensity

Particles in solution

Particles in solution have random orientation, both in time (thermal motion) and in space (no long range correlations). The sample as a whole is therefore **isotropic**. As a result, the scattering intensity only depends on the **modulus** of \vec{q} , $q = 4\pi \sin(\theta) / \lambda$.

Scattering from a single particle in solution, averaged over time:

$$I_1(q) = \frac{1}{V} \overline{\langle f_1(\vec{q}) f_1^*(\vec{q}) \rangle}_{\Omega}$$

Modulus

Vector

The form factor $P(q)$ is the normalized signature in q -space of a particle in solution.

$$P(q) = \frac{\overline{\langle f_1(\vec{q}) f_1^*(\vec{q}) \rangle}_{\Omega}}{r_0^2 V_{\text{obj}}^2 \langle \Delta \rho \rangle^2}$$

$$P(0) = 1$$

Average Electronic Density contrast

Particle volume

Monodispersity and ideality

- **Monodispersity**

- Yes ← Identical particles
- No ← Size and Shape polydispersity

- **Ideality**

- Yes ← No correlations between particles positions
(No short-range or long-range interactions)
- No ← Correlations between particles positions
(Existence of short-range or long-range interactions)

Ideal and monodisperse solutions

- **Ideal**

$$I(q) = \sum_{i=1,N} I_i(q) = \frac{1}{V} \sum_{i=1,N} \overline{\langle f_i(\vec{q}) f_i^*(\vec{q}) \rangle}_{\Omega}$$

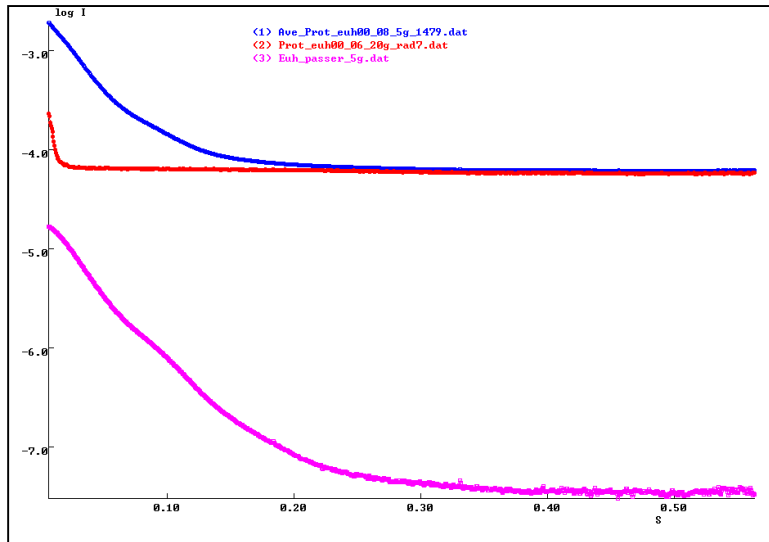
- **Monodisperse**

$$I_i(q) = I_1(q) \quad \forall i$$

- **Ideal and monodisperse**

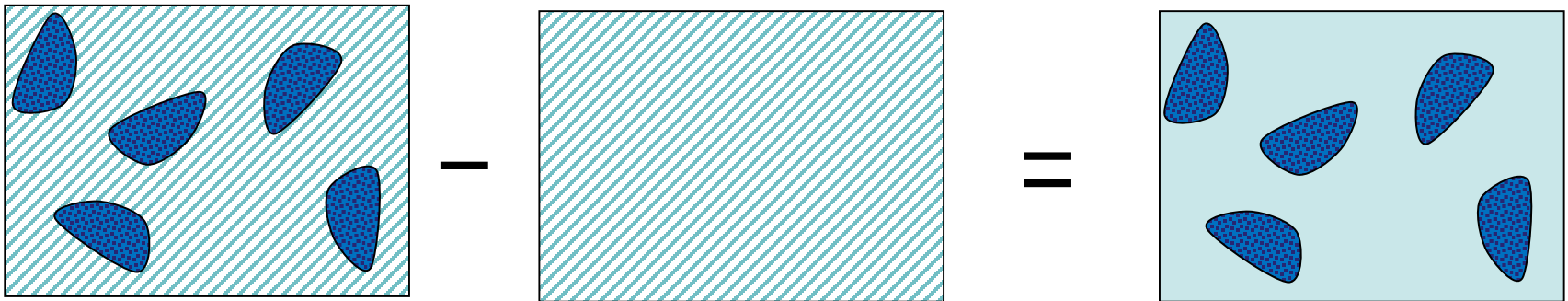
$$I(q) = N I_1(q) = \frac{N}{V} \overline{\langle f_1(\vec{q}) f_1^*(\vec{q}) \rangle}_{\Omega}$$

Solvent scattering and contrast



To obtain scattering solely from the contrasting particles, intrinsic solvent scattering must be measured **very accurately** and subtracted, which also permits to subtract contribution from parasitic background (slits, sample holder etc) which should be reduced to a minimum.

$$I_{\text{solution}}(q) - I_{\text{buffer}}(q) = I_{\text{particles}}(q)$$



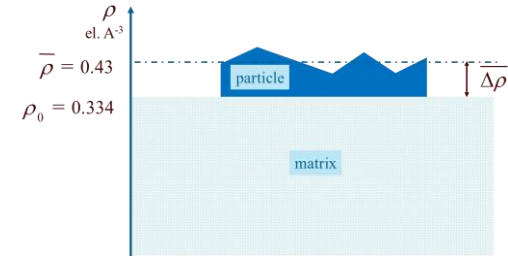


Do not confuse !



contrast effect

$$\Delta\rho(\vec{r}) = \rho(\vec{r}) - \rho_0$$

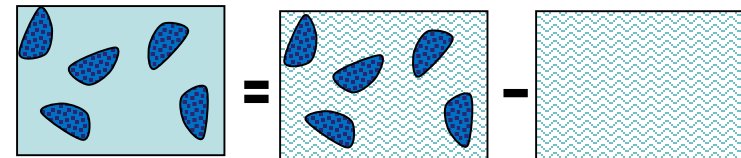


$$f(\vec{q}) = -r_0 \int_V \Delta\rho(\vec{r}) e^{-i\vec{q}\cdot\vec{r}} d^3\mathbf{r}, \vec{q} \neq 0$$



$$I(q) = \frac{1}{V} \left\langle f(\vec{q}) f^*(\vec{q}) \right\rangle_\Omega$$

buffer subtraction



$$I_{particles,exp}(q) = I_{solution,exp}(q) - I_{buffer,exp}(q)$$

$$i_1(q)$$

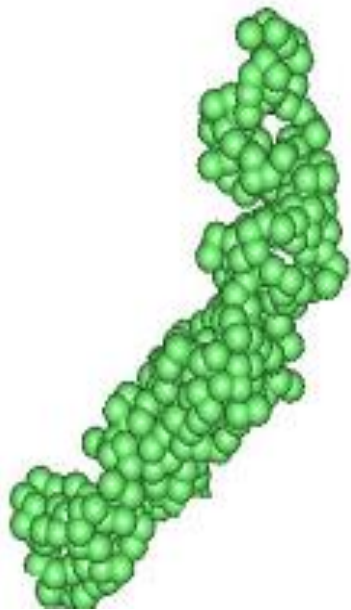
Ideality

$$I(q)$$

Monodispersity



One must check that both assumptions are valid for the sample under study.



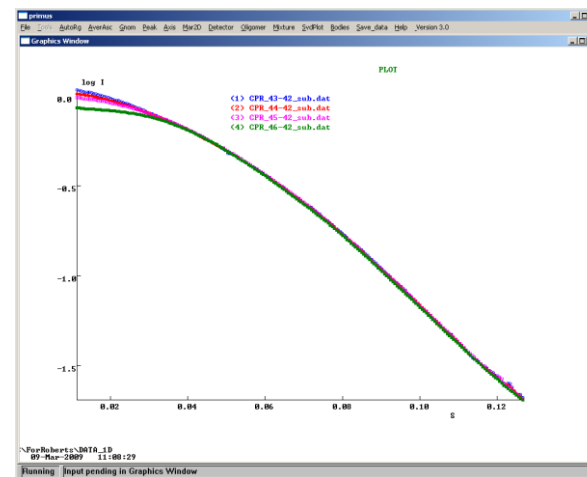
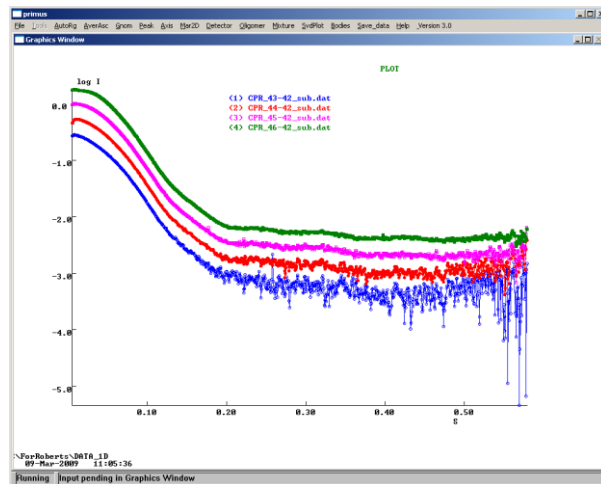
molecule

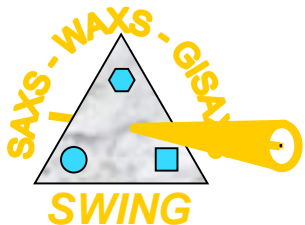


experimental

Checking the validity of both assumptions for the sample under study is crucial for non erroneous data interpretation

- Size Monodispersity must be checked **independently**
 - Purification protocol :SEC, DLS, AUC, MALS, etc.
- Ideality : reached by working in buffers with screened interactions or at high dilution
 - In practice : measurements at decreasing concentrations and checks whether the scattering pattern is independent of concentration.

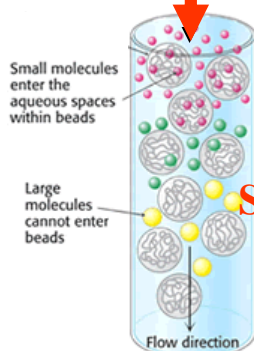




SE-HPLC / Solution Sampler



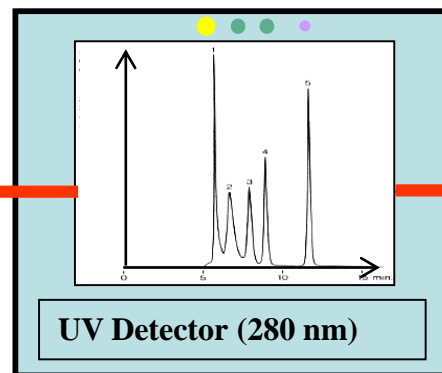
Flow rate 250 $\mu\text{l}/\text{min}$



Size Exclusion

- Monodisperse solution
- Aggregation is eliminated
- Oligomeric conformations can be distinguished
- Equilibrium states can be transiently separated
- Perfect background subtraction
- Automatic concentration series

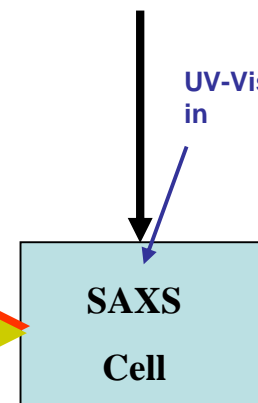
Incident X-ray



Flow rate 5-40 $\mu\text{l}/\text{min}$

Pure sample

- Small volumes ($\sim 10 \mu\text{l}$)
- No dilution
- High rate (a few minutes)



UV-Vis
out

BIOPHYSICAL INFORMATION

Biophysical information

- Guinier Analysis
- Kratky plot : why is it so interesting ?
- « Real-space SAXS » : Distance correlation function $P(r)$

Biophysical information

- Guinier Analysis
- Kratky plot : why is it so interesting ?
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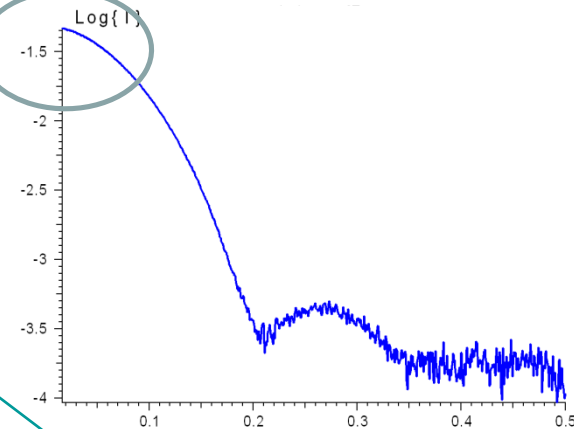
Asymptotic behaviour at small angles : Guinier law

Close to $q=0$, the scattering intensity of a particle can be described by a Gaussian curve.

The validity domain actually depends on the shape of the particle and is around $q < 1.3 / R_g$ for a globular shape.

$$I(q) = I(0) \exp\left(\frac{-q^2 R_g^2}{3}\right)$$

Extrapolated intensity at origin



Radius of gyration



*Prof. André Guinier
1911-2000
Orsay, France*

Guinier law, in Log scale :

$$\ln[I(q)] = \ln[I(0)] - \frac{q^2 R_g^2}{3}$$

The Guinier law is equivalent of a linear variation of $\ln(I(q))$ vs q^2 (Guinier plot). Linear regression on the experimental Guinier plot directly provides R_g and $I(0)$.

Mass retrieval from Guinier analysis

$$I(q) = I(0) \exp\left(\frac{-q^2 R_g^2}{3}\right)$$

Absolute Unit : cm^{-1}

Classical electron radius

$$I(0) = \frac{c \cdot M \cdot r_0^2}{N_A} \cdot [v_p (\rho_{prot} - \rho_{buf})]^2$$

Mass concentration

Protein specific volume

Electronic density contrast

$$R_g^2 = \frac{\int_V r^2 (\rho_{prot}(\vec{r}) - \rho_{buf}) d\vec{r}}{\int_V (\rho_{prot}(\vec{r}) - \rho_{buf}) d\vec{r}}$$

R_g depends on the volume
AND on the shape of the particle

$I(0)$ gives an independent estimation
of the molar mass of the protein
(only if the mass concentration
and specific volume are precisely

Typically : known ...)

$$M (\text{kDa}) = (1200 \sim 1600) * I_0 (\text{cm}^{-1}) / C (\text{mg/ml})$$

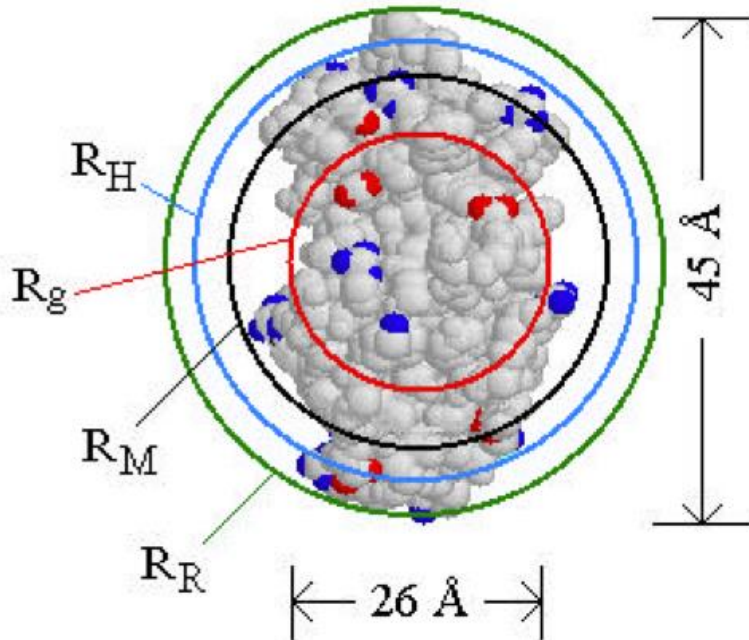
For globular proteins : $R_g (\text{\AA}) \approx 6.5 * M^{\frac{1}{3}}$, M in kDa
For unfolded proteins : $R_g (\text{\AA}) \approx 8.05 * M^{0.522}$

Bernado et al. (2009), Biophys. J., 97 (10), 2839-2845.

Radius of gyration

$$R_{g_{\text{exp}}}^2 = \frac{\int_V r^2 \Delta \rho(\vec{r}) d\vec{r}}{\int_V \Delta \rho(\vec{r}) d\vec{r}}$$

Lysozyme



Useful definitions of R_g

$$R_g^2 = \frac{1}{N} \sum_i \|\vec{r}_i - \vec{r}_{COM}\|^2 \quad \text{by atoms}$$

$$R_g^2 = \frac{\int_V r(r) r^2 dr}{\int_V r(r) dr} \quad \text{by electron density}$$

$$R_g^2 = \frac{1}{2N(N-1)} \sum_i \sum_j \|\vec{r}_i - \vec{r}_j\|^2 \quad \text{by atom pairs}$$

$$R_g^2 = \frac{\int_0^\infty r^2 p(r) dr}{\int_0^\infty p(r) dr} \quad \text{by pair distribution}$$

graphic: www.silver-colloids.com/Papers/hydrodynamic-radius.pdf

R_g radius of gyration

R_H hydrodynamic radius (not always $> R_g$!)

R_R maximum hard sphere radius

R_M radius of mass-equivalent sphere

* center of mass of the *electron* density

Sphere $R_g = \sqrt{\frac{3}{5}} R$

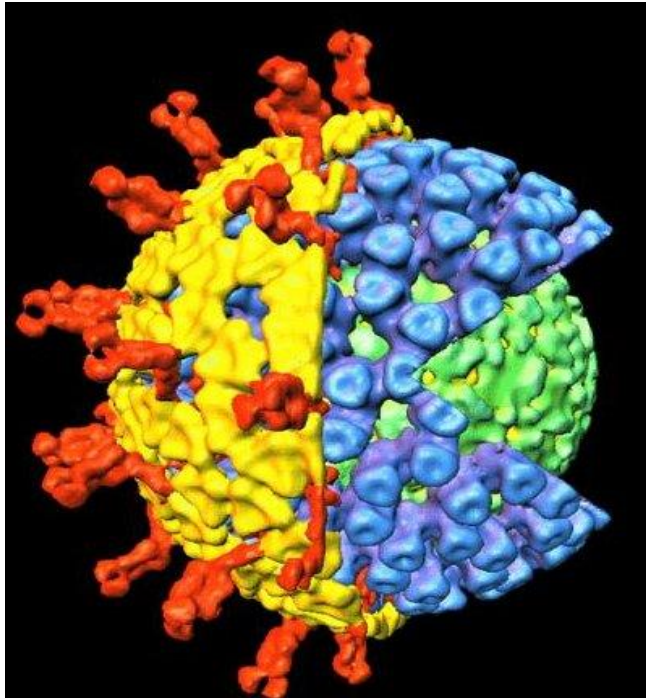
Thin rod $R_g = \sqrt{\frac{1}{12}} L$

Thin disk $R_g = \sqrt{\frac{1}{2}} R_{\text{disk}}$

Courtesy: Richard Gillilan, Cornell U., USA

Basic law of reciprocity in scattering

Rotavirus VLP : diameter = 750 Å, 44 MDa

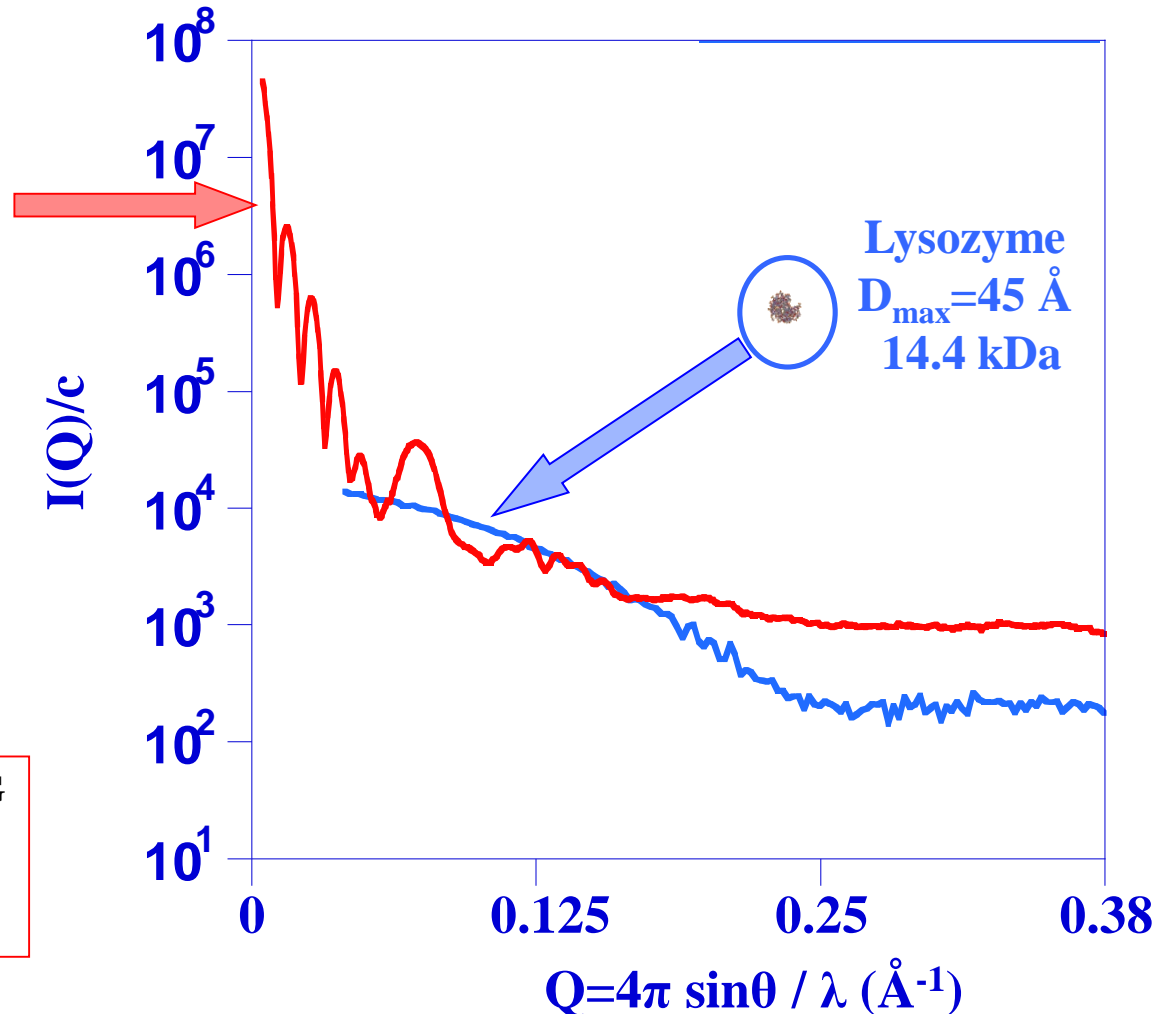


Typical range on beamline SWING

$$2 \cdot 10^{-3} < q < 7 \cdot 10^{-1} \text{ Å}^{-1}$$

$$750 \text{ Å} > D_{\text{max}}$$

$$3100 \text{ Å} > d_{\text{Bragg}} > 9 \text{ Å}$$

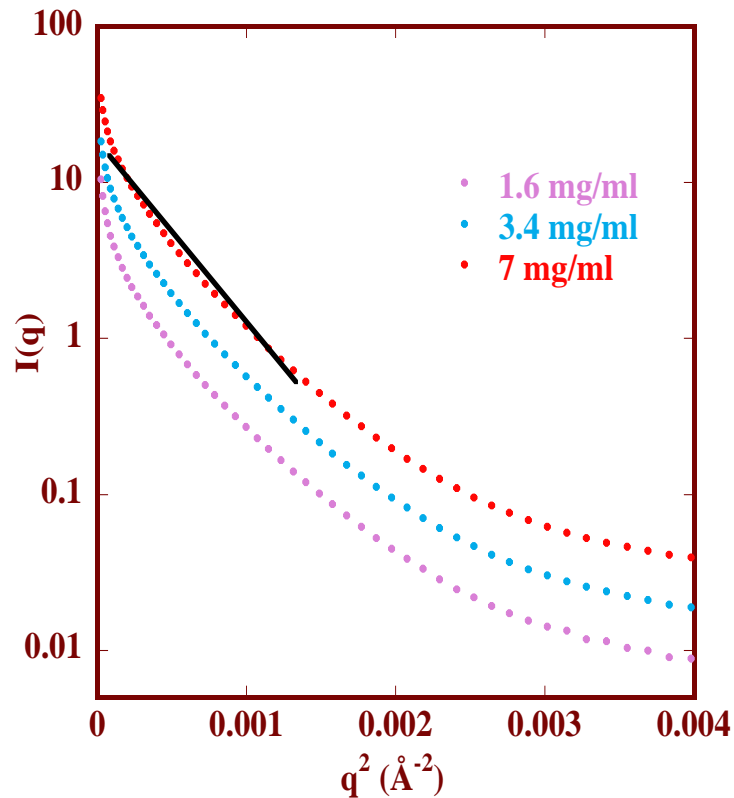


Long distance correlations \longleftrightarrow modulations at small Q

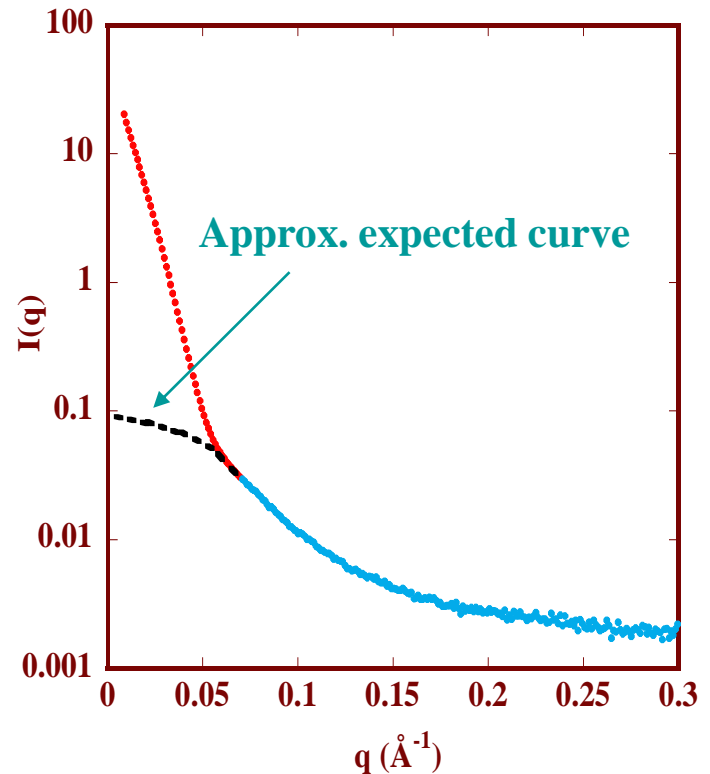
Evaluation of the solution properties

Irreversible aggregation

→ Useless data: the whole curve is affected



$I(0)$: > 150 fold the expected value for the given MM



Swing – Domaine 1-242 de RRP44 – 07/08

(Courtesy D. Durand, IBBMC, Orsay)

Evaluation of the solution properties

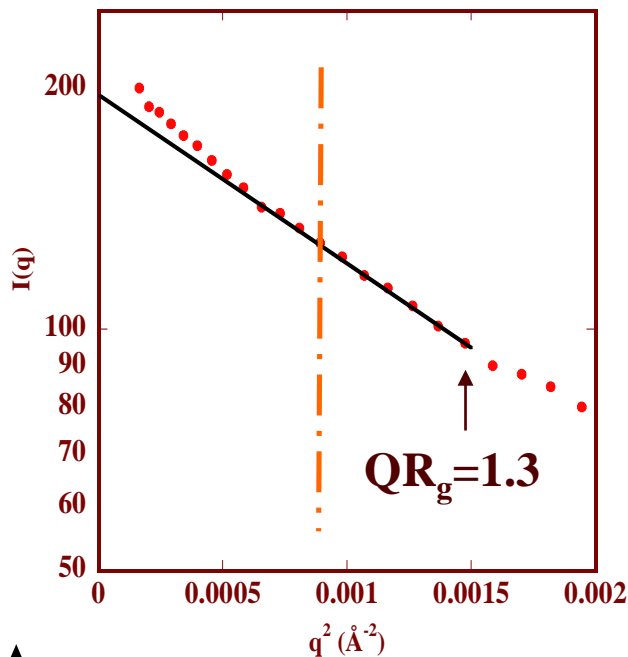
Weak aggregation



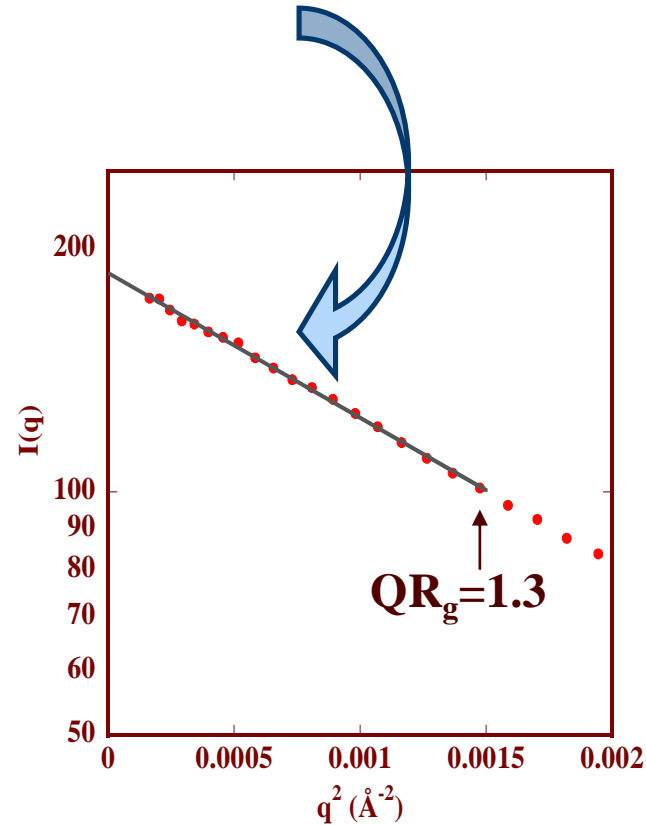
possible improvement

centrifugation, buffer change

Nanostar –PR65 protein



$R_g \sim 38 \text{ \AA}$ – too high!!



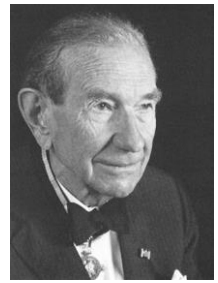
$R_g \sim 36 \text{ \AA}$

(Courtesy D. Durand, IBBMC, Orsay)

Biophysical information

- Guinier Analysis
- Kratky plot : why is it so interesting ?
- « Real-space SAXS » : Distance correlation function $P(r)$

Kratky Plot

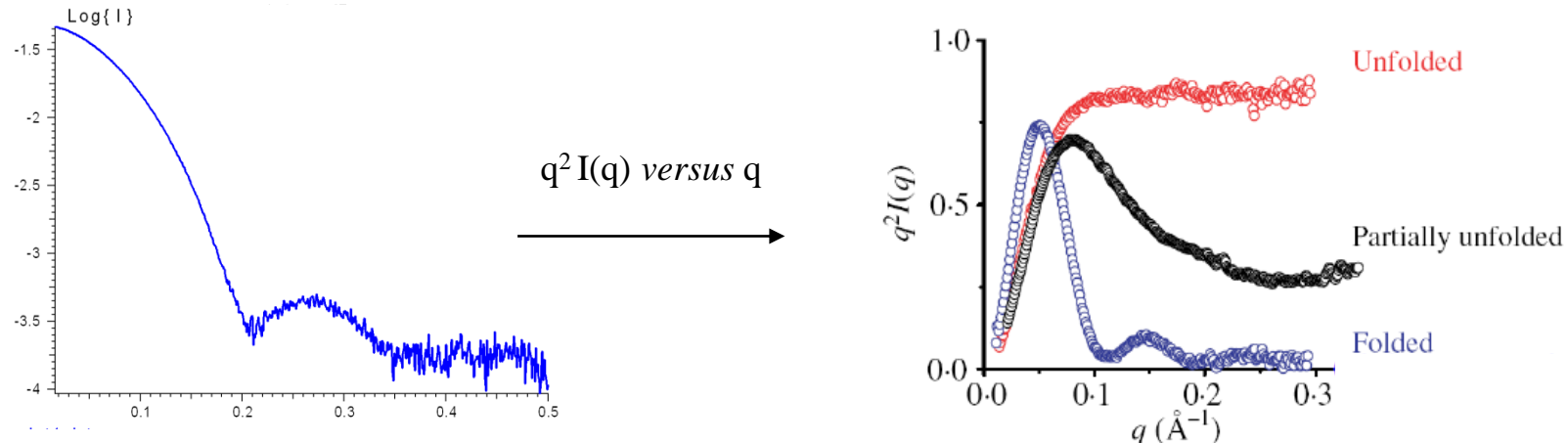


Prof. Otto Kratky
1902-1995
Graz, Austria

SAXS provides a sensitive means to *evaluate the degree of compactness* of a protein:

- To determine whether a protein is globular, extended or unfolded
- To monitor the folding or unfolding transition of a protein

This is most conveniently represented using the so-called Kratky plot:



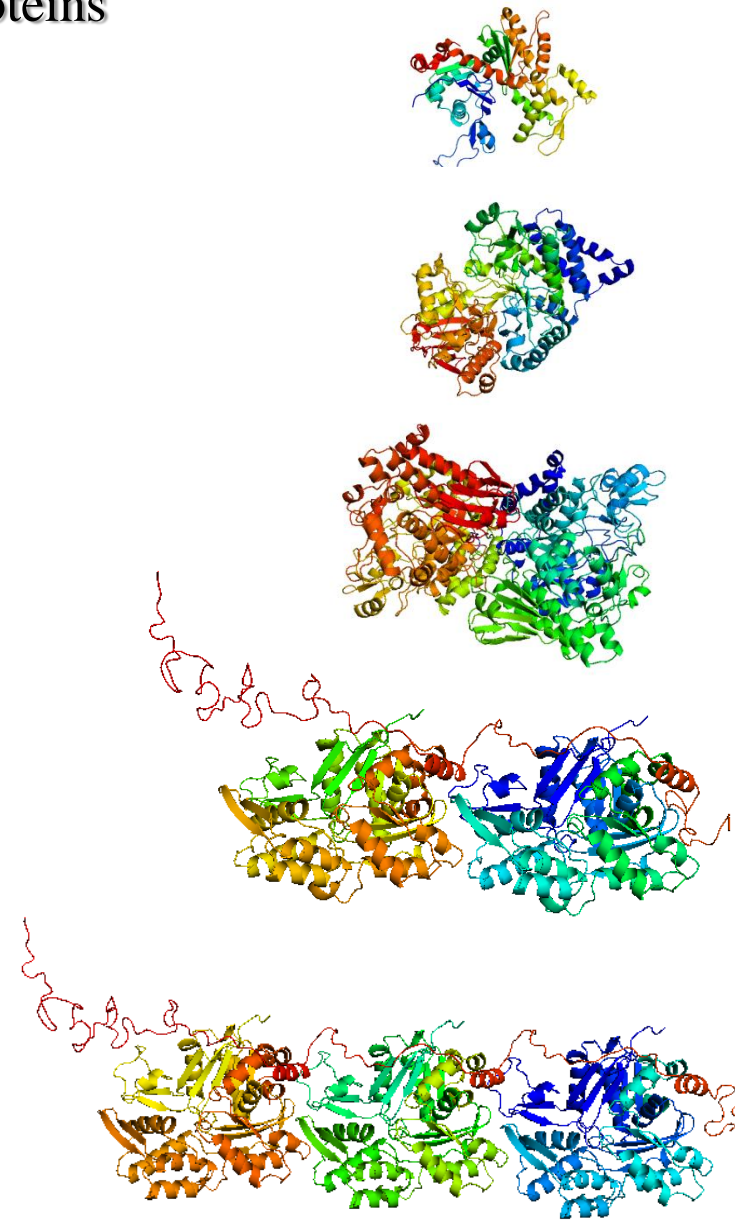
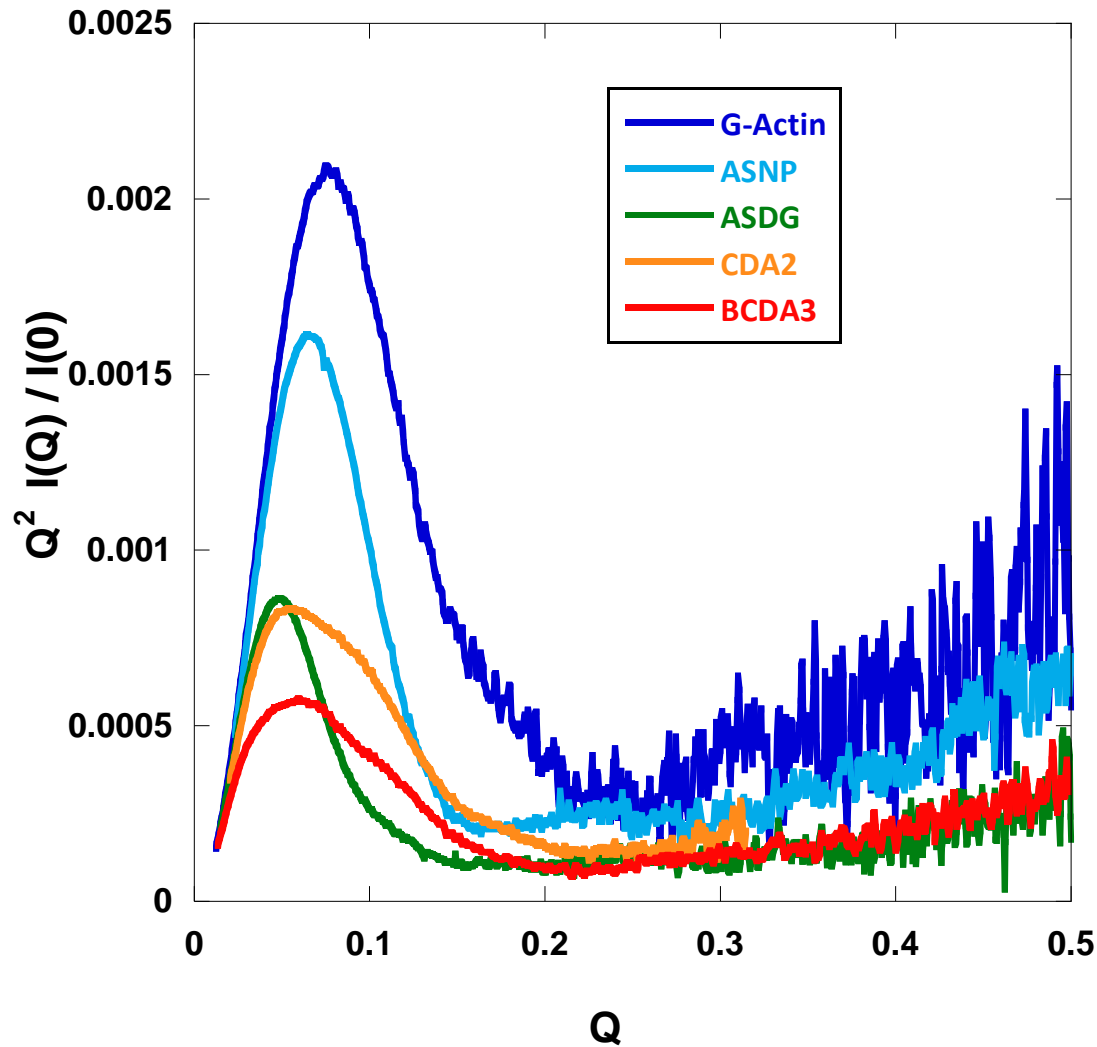
Putnam, D., et al. (2007) *Quart. Rev. Biophys.* 40, 191-285.

Folded particle : *bell-shaped curve* (asymptotic behaviour $I(q) \sim q^{-4}$)

Random polymer chain : *plateau* at large q-values (asymptotic behaviour in $I(q) \sim q^{-2}$)

Extended polymer chain : *increase* at large q-values (asymptotic behaviour in $I(q) \sim q^{-1.x}$)

Kratky Plots of folded proteins



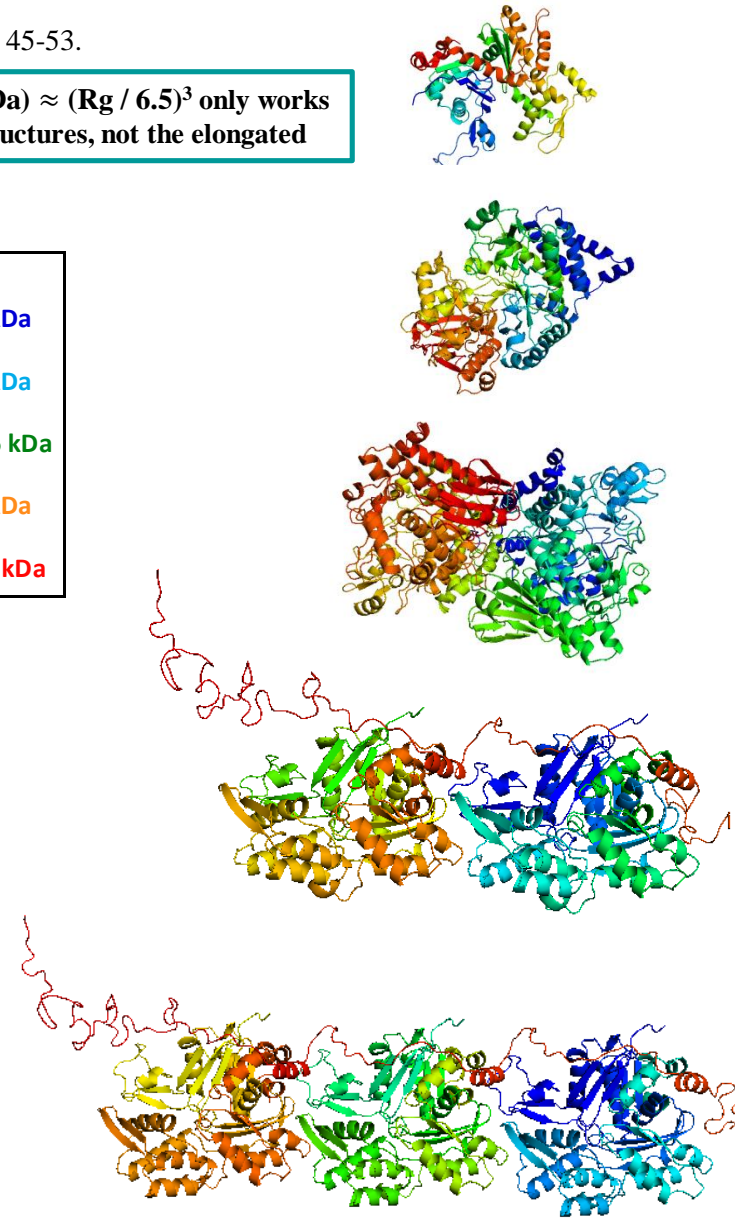
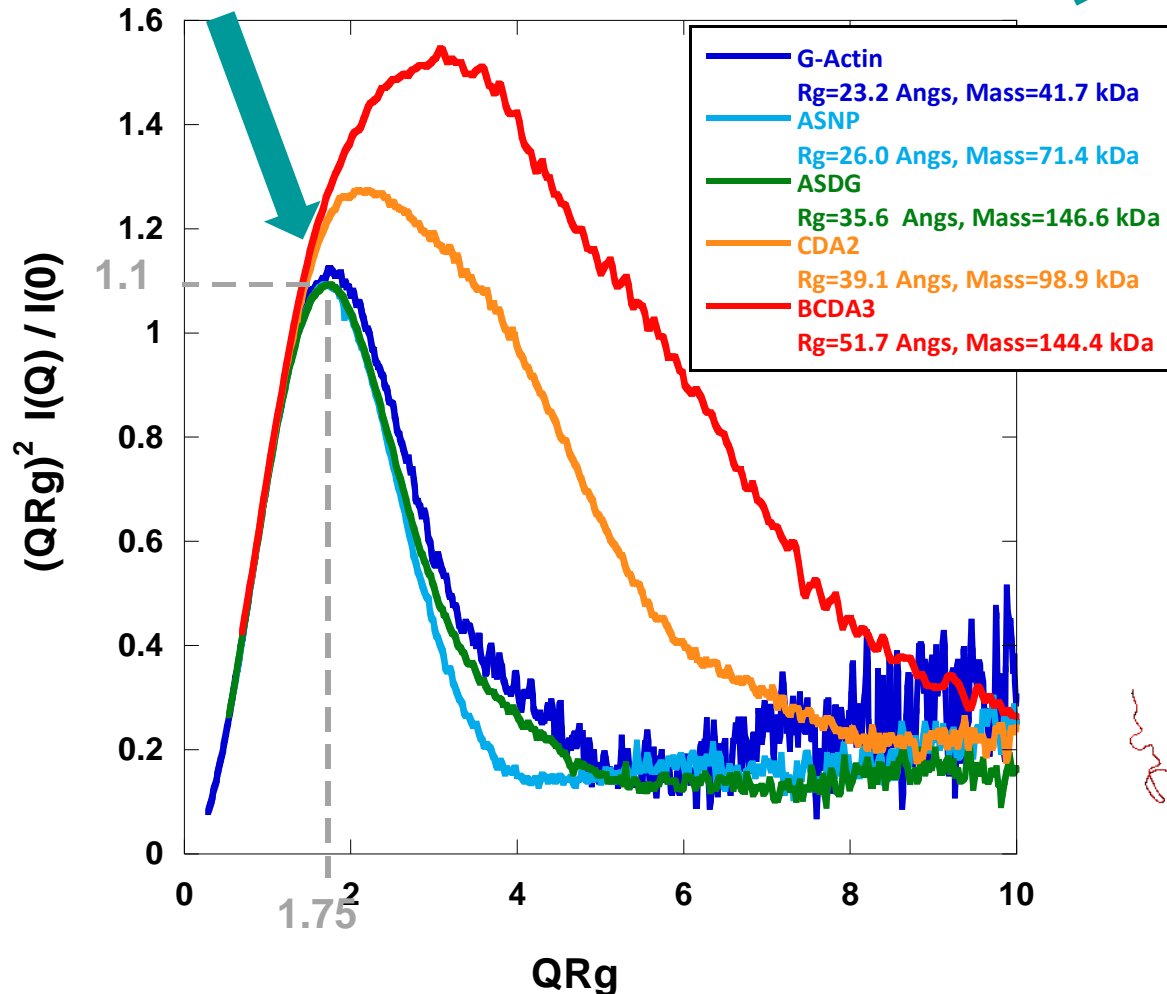
Folded proteins display a bell shape. Can we go further?

Dimensionless Kratky Plots of folded proteins

Introduced for biology in Durand et al. (2010), J. Struct. Biol. **169**, 45-53.

For globular structures, DLKPs fold into the same maximum

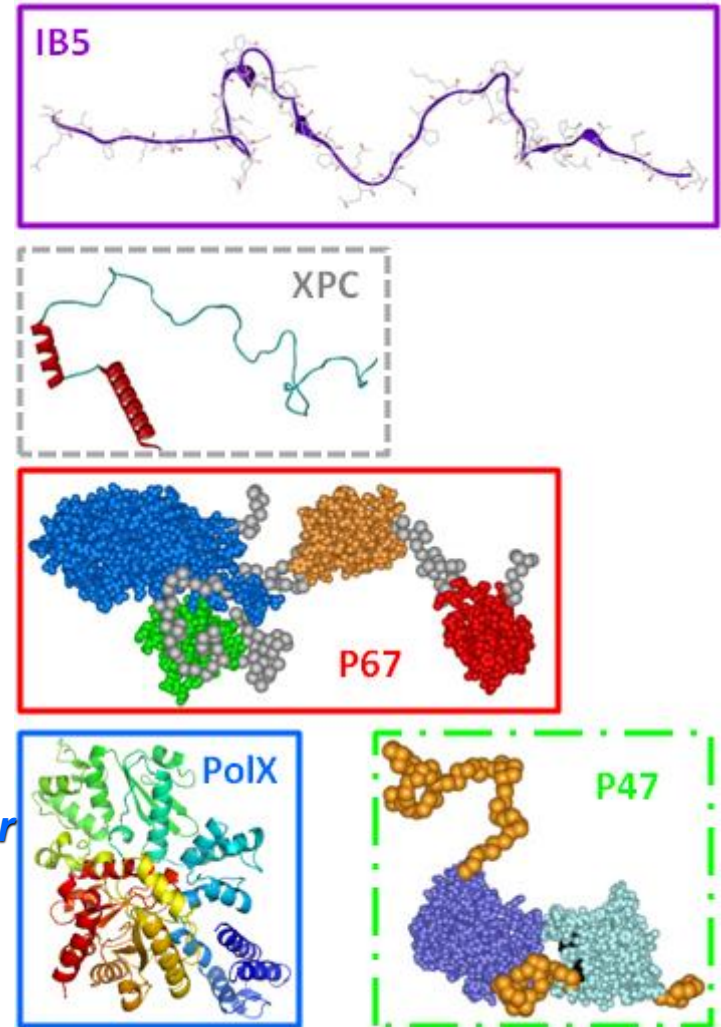
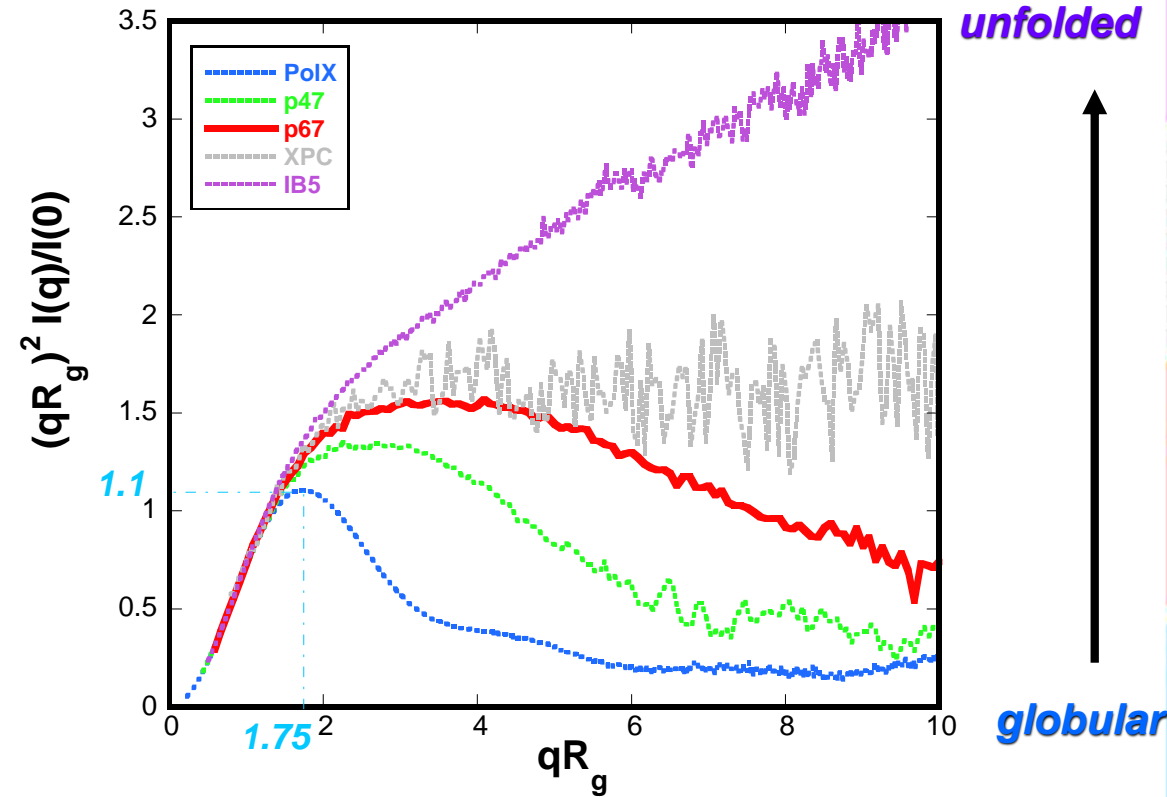
The relation $M_{\text{Rg}}(\text{kDa}) \approx (\text{Rg} / 6.5)^3$ only works for the globular structures, not the elongated



The maximum value on the dimensionless bell shape tells if the protein is globular.

Dimensionless Kratky Plots of (partially) unfolded proteins

Receveur-Bréchet V. and Durand D (2012), Curr. Protein Pept. Sci., 13:55-75.



The bell shape vanishes as folded domains disappear and flexibility increases.

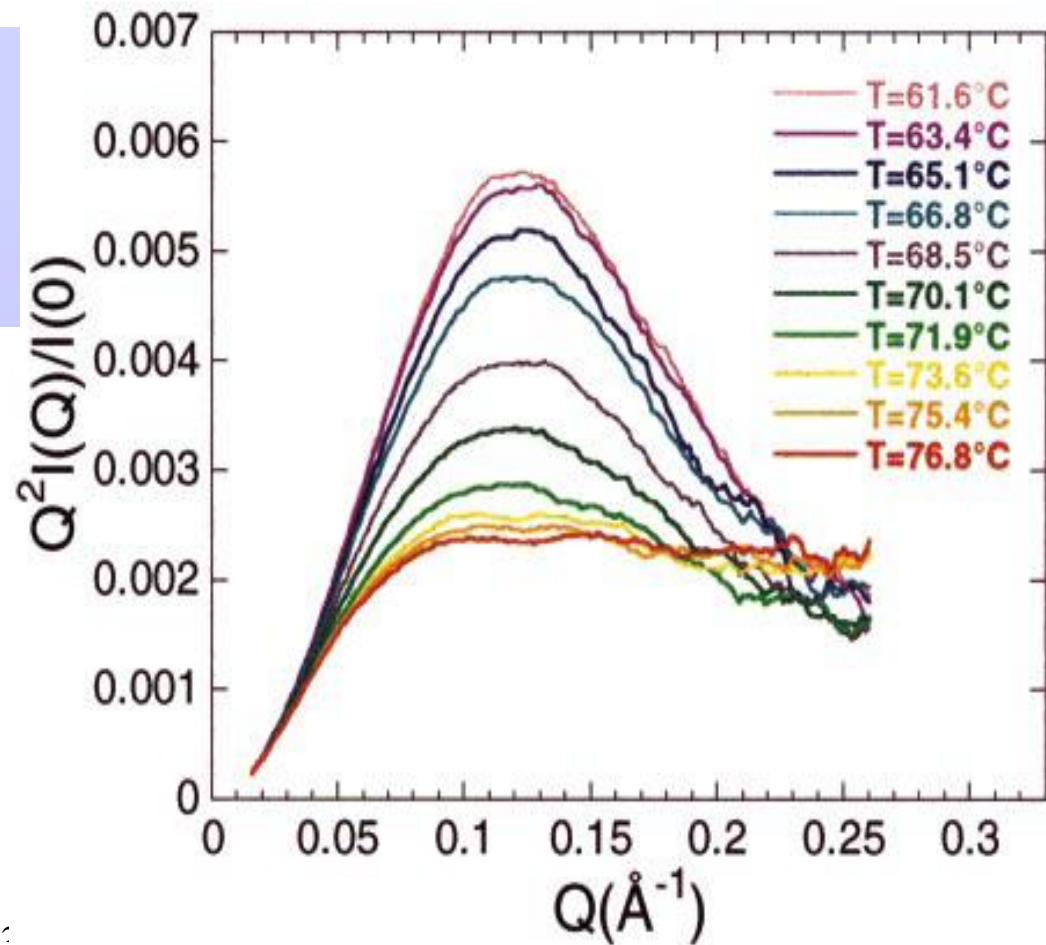
The curve increases at large Q as the structure extends.

Kratky Plot : NCS heat unfolding



In practice, thin Gaussian chains do not exist.

In spite of the plateau at T=76°C, NCS is not a Gaussian chain when unfolded, but a thick chain with persistence length

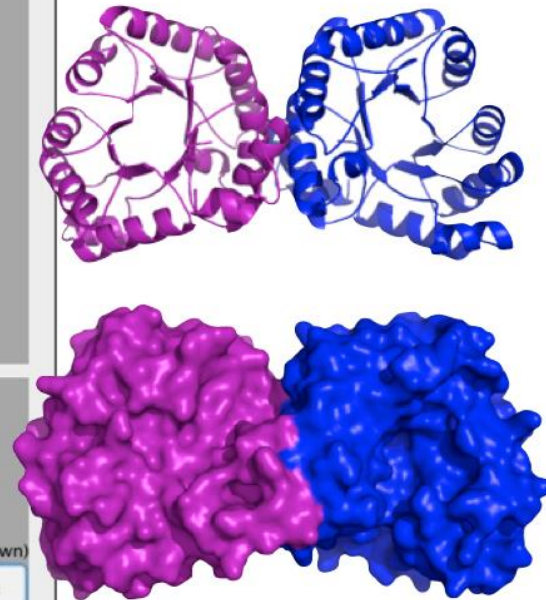
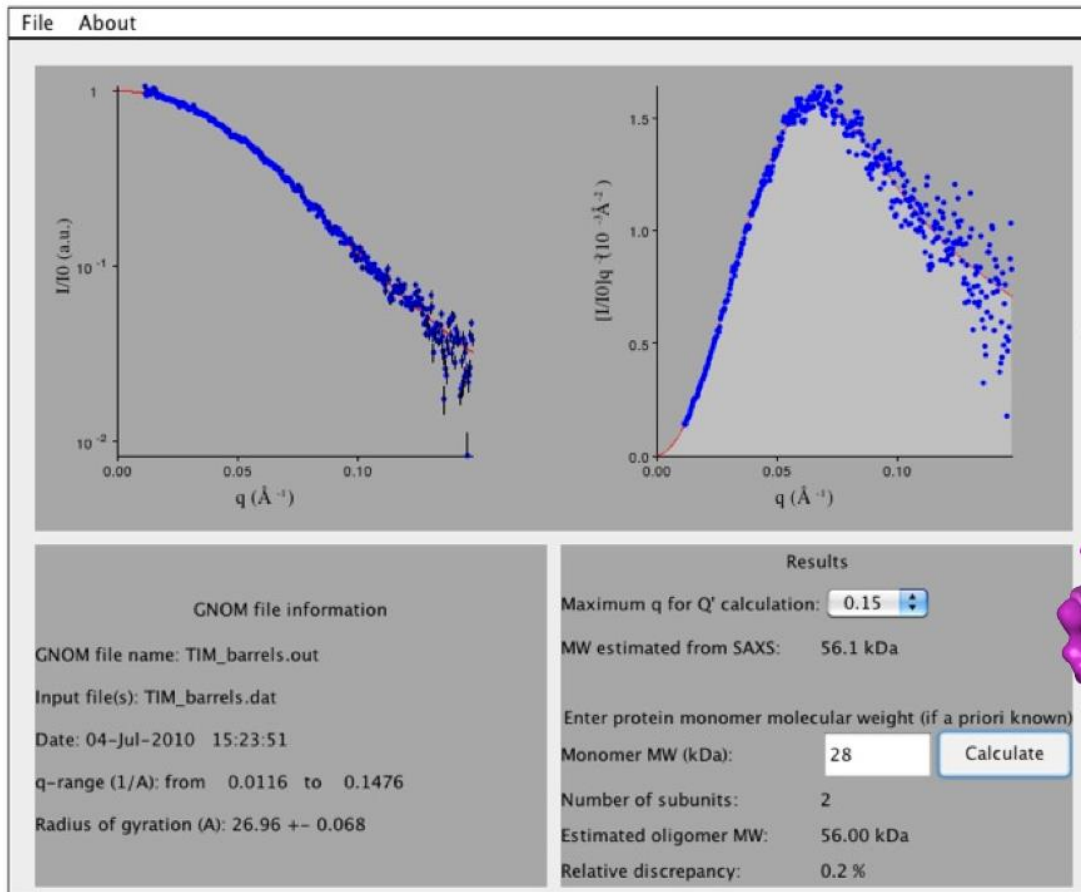


Pérez et al., *J. Mol. Biol.*(2001), 308, 721-743

Molecular Weight estimation based on Porod invariant

<http://www.ifsc.usp.br/~saxs/saxsmow.html>

- does not require knowledge of concentration
- relies on Porod Volume theory + **structural database**
- does **not** work for proteins with unfolded domains



Other methods for MW estimation based on similar though different grounds were developed

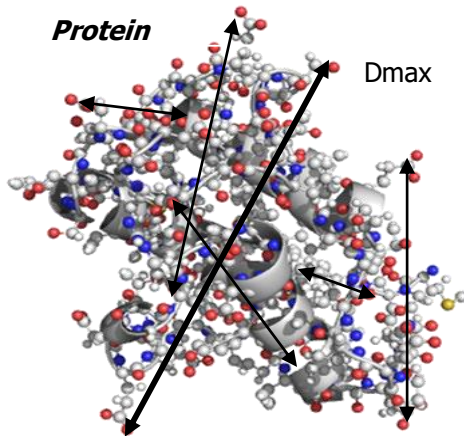
Rambo R. And Tainer J. (2013), Nature, **496**, 477-481.

Biophysical information

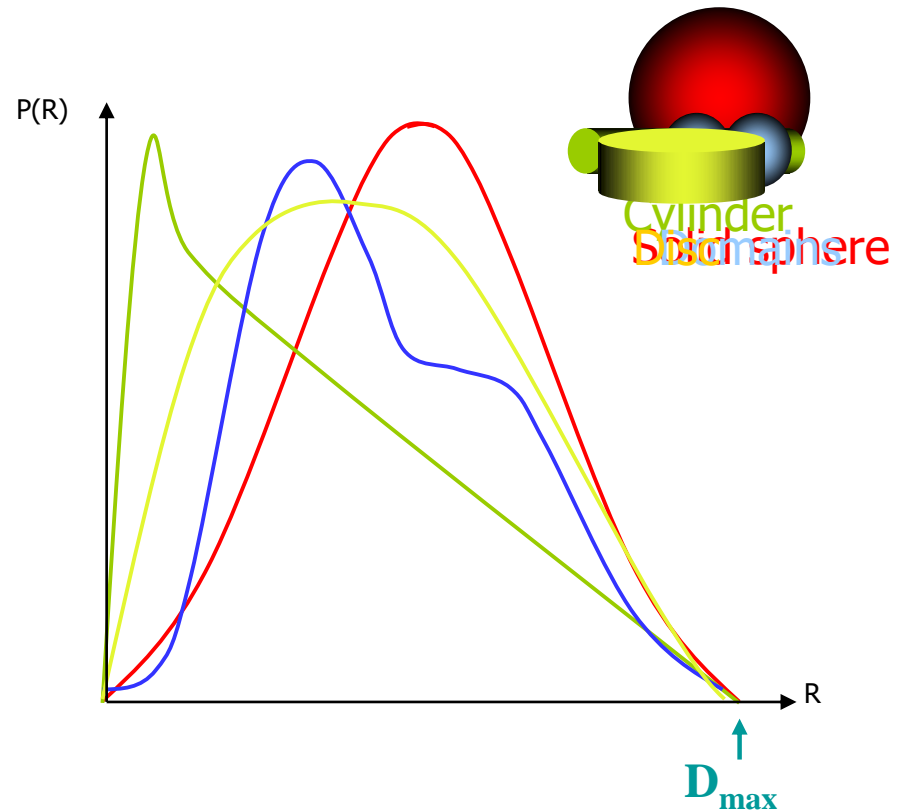
- Guinier Analysis
- Kratky plot : why is it so interesting ?
- « Real-space SAXS » : Pair distribution function $P(r)$

Pair Distribution Function $p(r)$

The distance distribution function $p(r)$ is proportional to the average number of atoms at a given distance, r , from any given atom within the macromolecule.



$p(r)$ vanishes at $r = D_{max}$

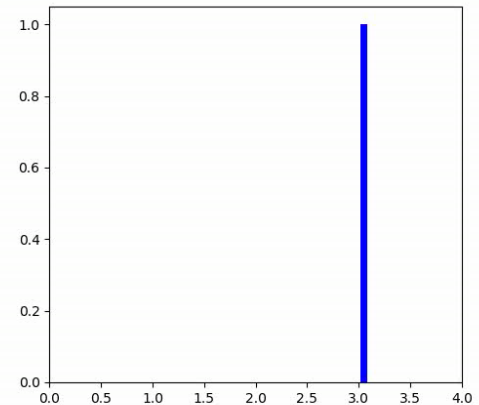
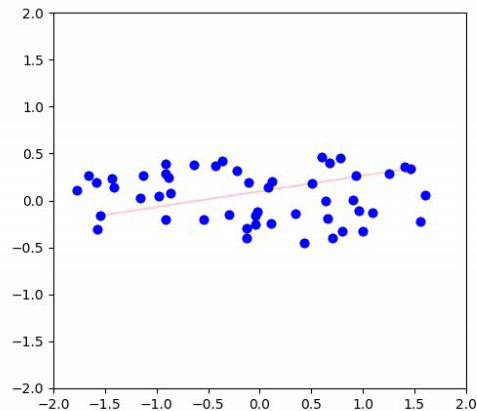
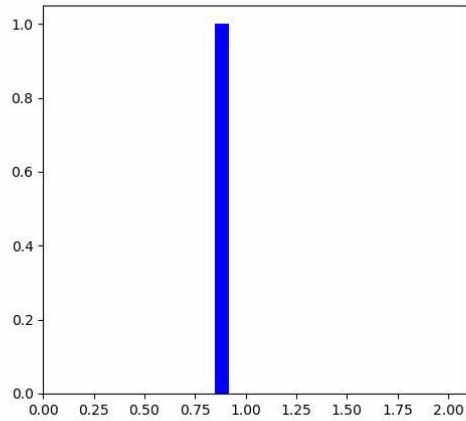
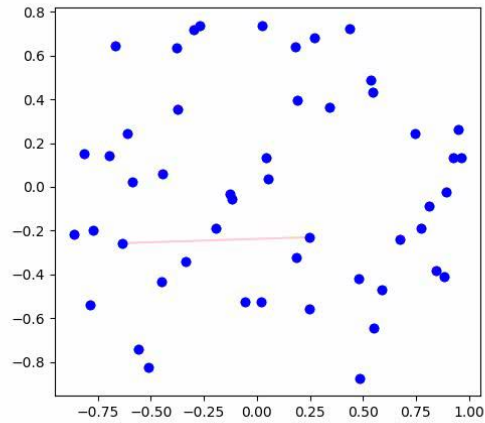


The distance distribution function characterises the shape of the particle in **real space**

Pair Distribution Function $p(r)$

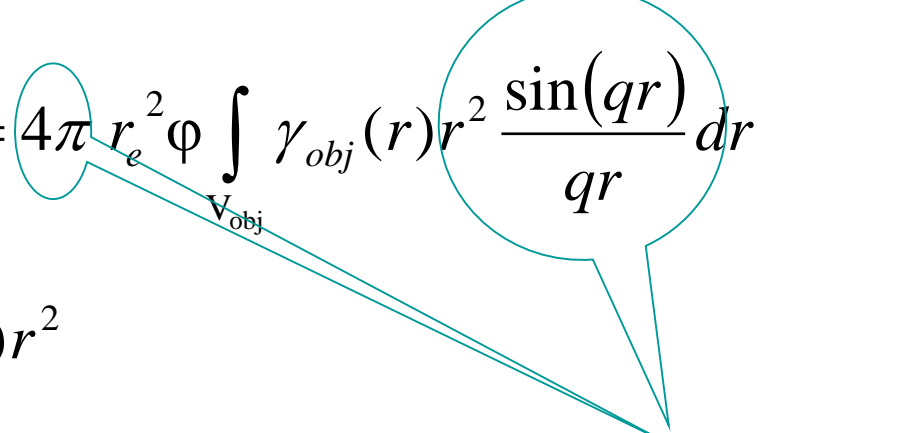


Marc-André Delsuc:
thanks for the python script



Relation between $p(r)$ and $I(Q)$

Intensity is the Fourier Transform of self-correlation function $\gamma_{obj}(r)$:

$$I(q) = 4\pi r_e^2 \varphi \int \gamma_{obj}(r) r^2 \frac{\sin(qr)}{qr} dr$$


And :

$$p(r) = \gamma_{obj}(r) r^2$$

Then :

$$I(q) = 4\pi r_e^2 \varphi \int_0^D p(r) \frac{\sin(qr)}{qr} dr$$

And :

$$p(r) = \frac{r^2}{2\pi^2 \varphi r_e^2} \int_0^\infty q^2 I(q) \frac{\sin(qr)}{qr} dq$$

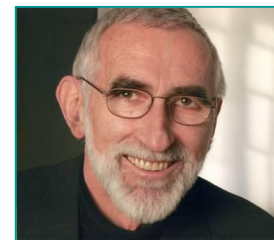
Fourier Transform for isotropic samples

$p(r)$ could be directly derived from $I(q)$. Both curves contain the same information.

However, direct calculation of $p(r)$ from $I(q)$ is made difficult and risky by $[q_{min}, q_{max}]$ truncation and data noise effects.

Back-calculation of the Distance Distribution Function

Glatter, O. *J. Appl. Cryst.* (1977) **10**, 415-421.



Prof. Otto Glatter
Guinier Prize 2012
Graz, Austria

Main hypothesis : the particle has a « finite » size, characterised by D_{\max} .

- D_{\max} is proposed by the user
- $p(r)$ is expressed over $[0, D_{\max}]$ by a linear combination of orthogonal functions

$$p_{\text{theoret}}(r) = \sum_1^M c_n \varphi_n(r)$$

- $I(q)$ is calculated by Fourier Transform of $p_{\text{theoret}}(r)$

$$I(q) = 4\pi r_e^2 \varphi \int_0^{D_{\max}} p_{\text{theoret}}(r) \frac{\sin(q \cdot r)}{q \cdot r} dr$$



Dr. Dmitri Svergun
Guinier Prize 2018
Hamburg, Germany

Svergun (1988) : program "GNOM"

$M \sim 30 - 100 \Rightarrow$ ill-posed LSQ \Rightarrow regularisation method

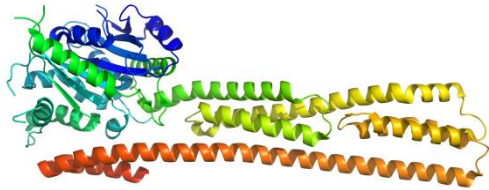
+ "**Perceptual criteria**" : smoothness, stability, absence of systematic deviations

- Each criterium has a predefined weight
- The solution is given a score calculated by comparison with « ideal values »

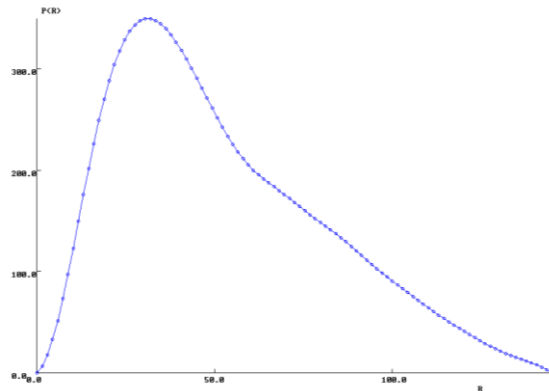
Distance Distribution Function

Experimental examples

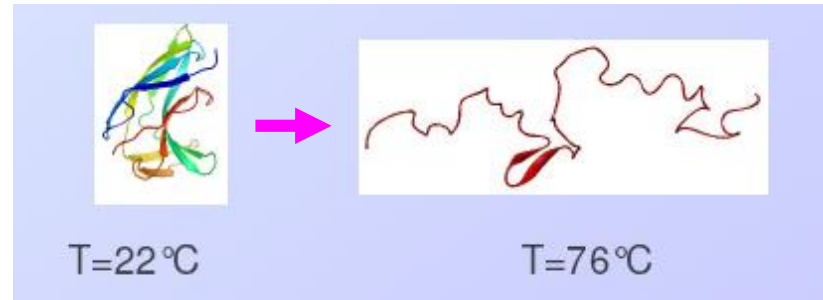
GBP1



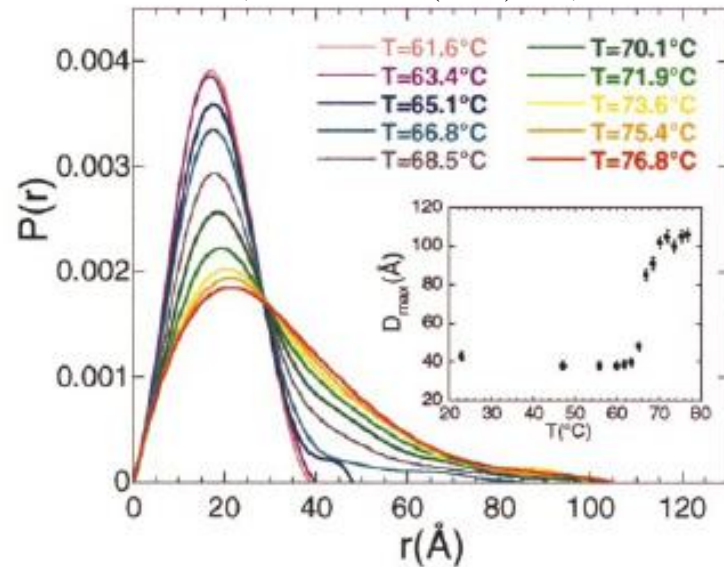
Real space: $R_g = 42.34$, $\langle R_g \rangle = 0.2775E+06$



Heat denaturation of Neocarzinostatin



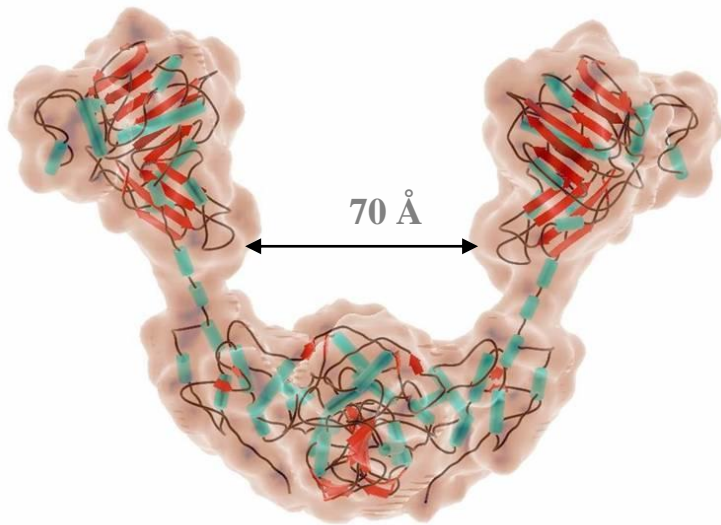
Pérez et al., J. Mol. Biol. (2001) 308, 721-743



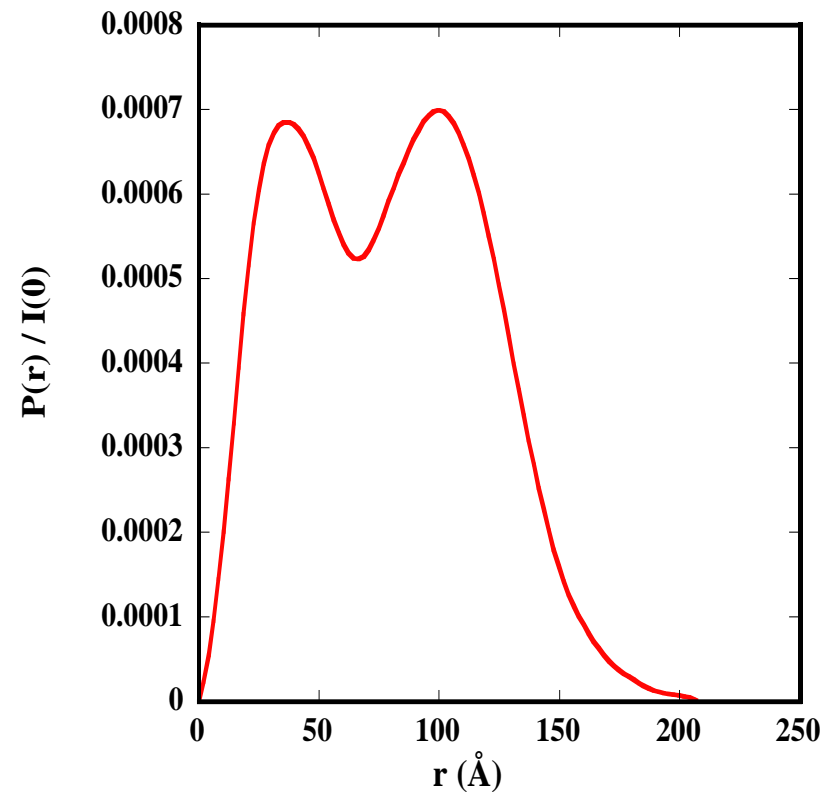
Distance Distribution Function

Experimental examples

Topoisomerase VI



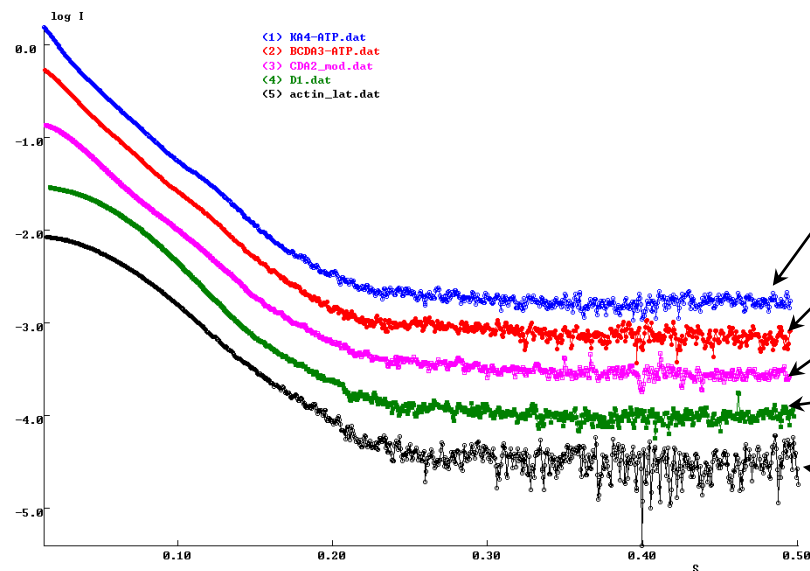
Bimodal distribution



M. Graille et al., Structure (2008), 16, 360-370.

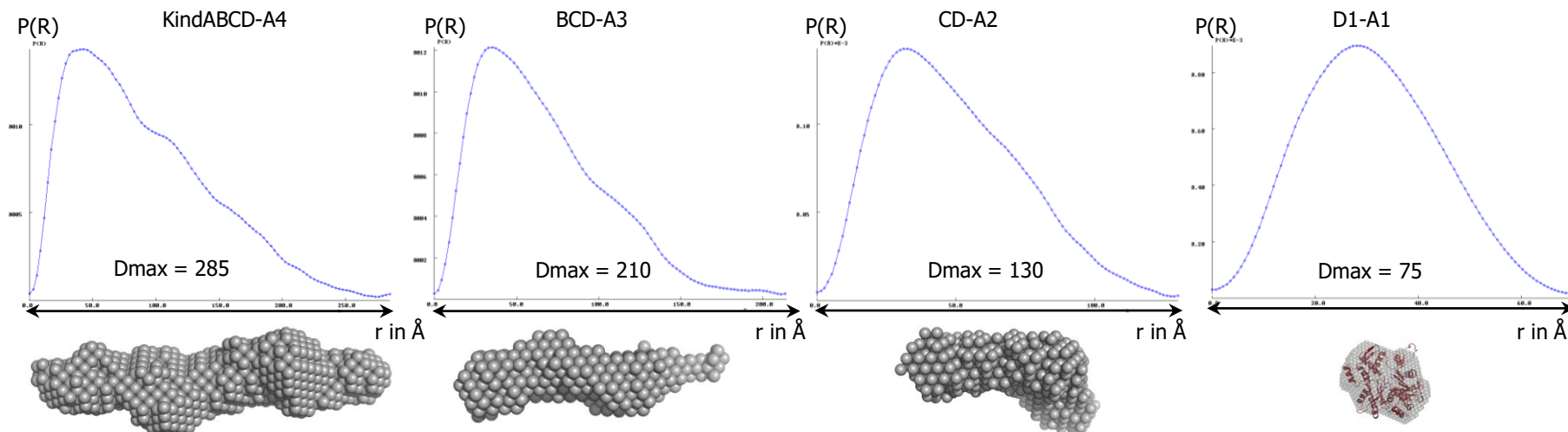
Distance Distribution Function

Scattering curves obtained on different complexes Spire-Actin and Actin alone



Complexes	Radius of gyration	Maximum diameter
	75.5 Å	285 Å
	55.5 Å	210 Å
	38.9 Å	130 Å
	25 Å	75 Å
	23.1 Å	70 Å

Histogram of intramolecular distances and ab initio molecular envelopes determined using DAMMIF



Distance Distribution Function

The radius of gyration and the intensity at the origin can be derived from $p(r)$ using the following expressions :

$$R_g^2 = \frac{\int_0^{D_{\max}} r^2 p(r) dr}{2 \int_0^{D_{\max}} p(r) dr}$$

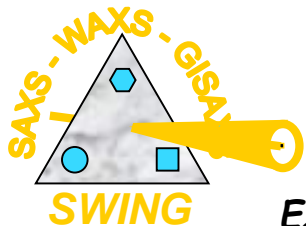
and

$$I(0) = 4\pi r_e^2 \phi \int_0^D p(r) dr$$

This alternative estimate of R_g makes use of the whole scattering curve, and is less sensitive to interactions or to the presence of a small fraction of oligomers.

Comparison of estimates from Guinier analysis and from $P(r)$ is a useful cross-check.

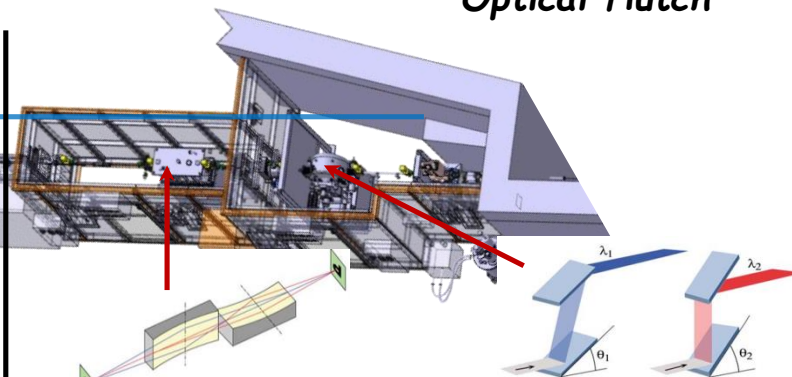
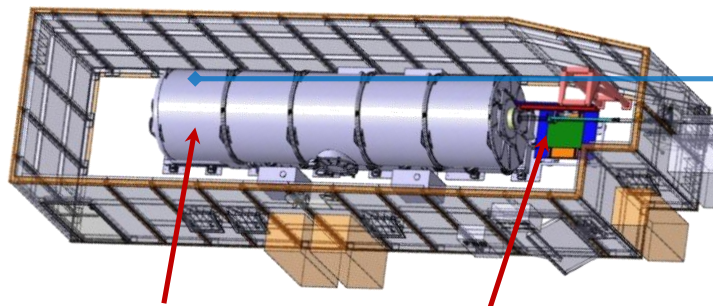
A FEW EXPERIMENTAL CONSIDERATIONS



Set-up for BioSAXS at beamline SWING

Experimental Hutch

Optical Hutch

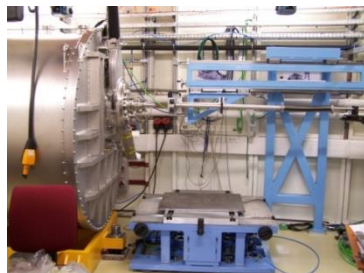


Tunnel under vacuum with X-rays detector

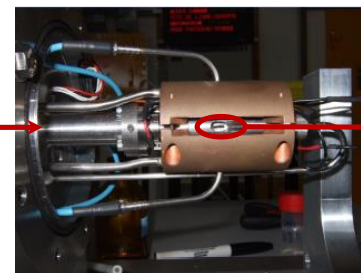
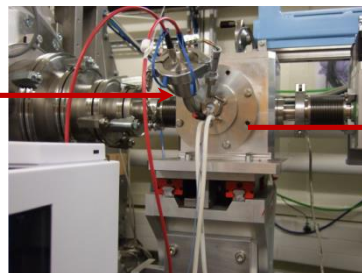
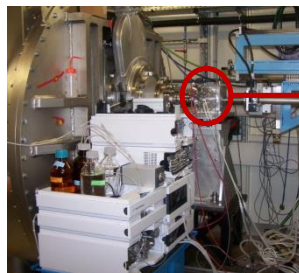
X,Y and Z motorized table

Double crystals mirror

Monochromator



Sample environment dedicated to the biology :



SEC-HPLC device

SEC-SAXS

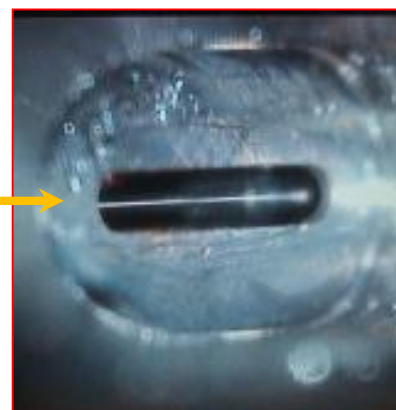
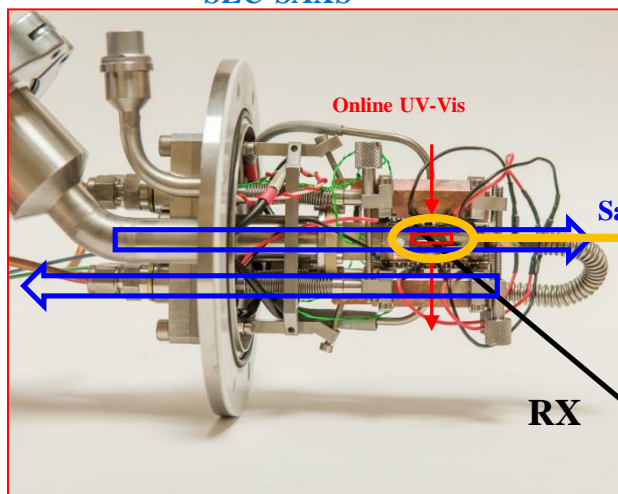
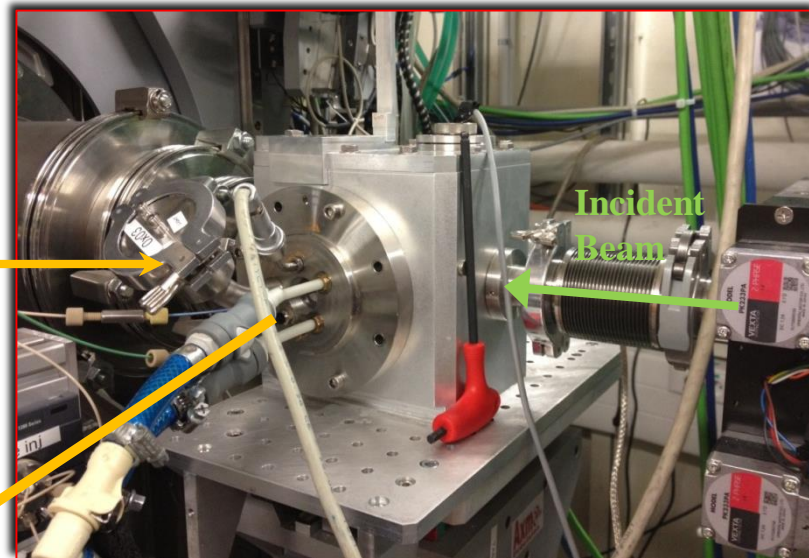
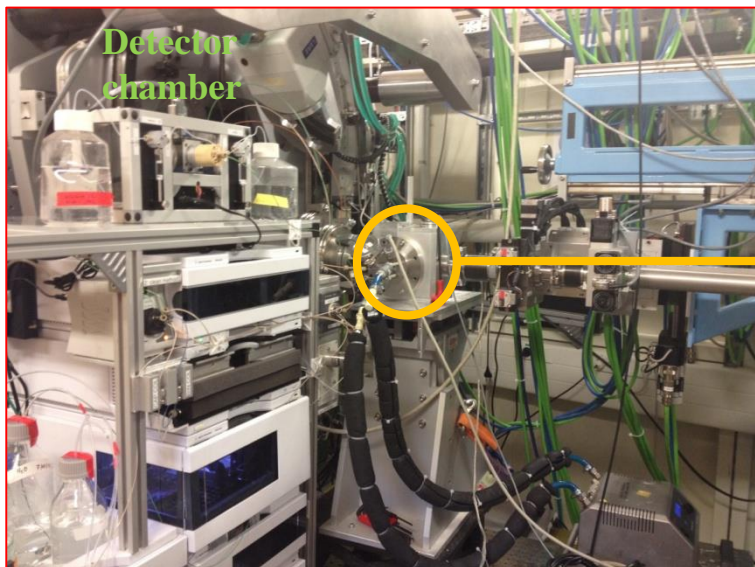
SAXS cell vacuum chamber

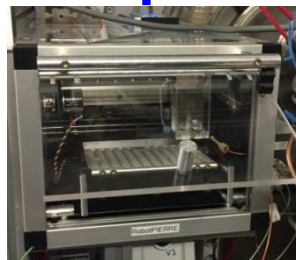
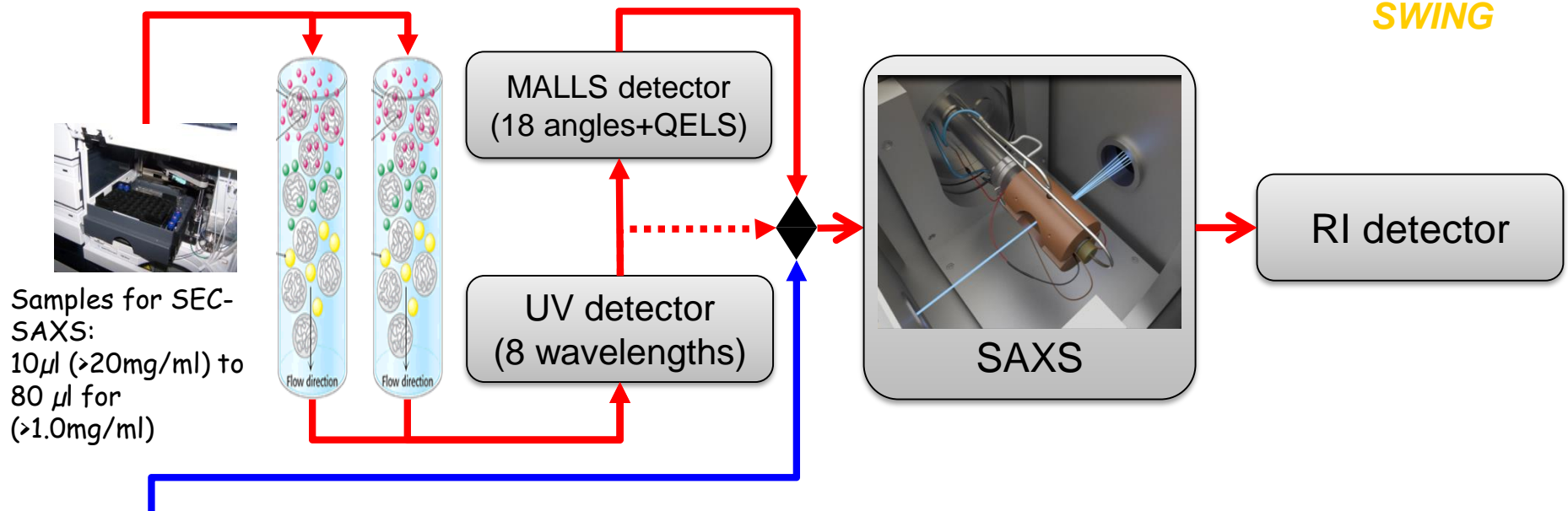
Details of the SAXS cell

Quartz capillary

Set-up for BioSAXS at Beamline SWING

G. David and J. Pérez (2009), J. Appl. Cryst



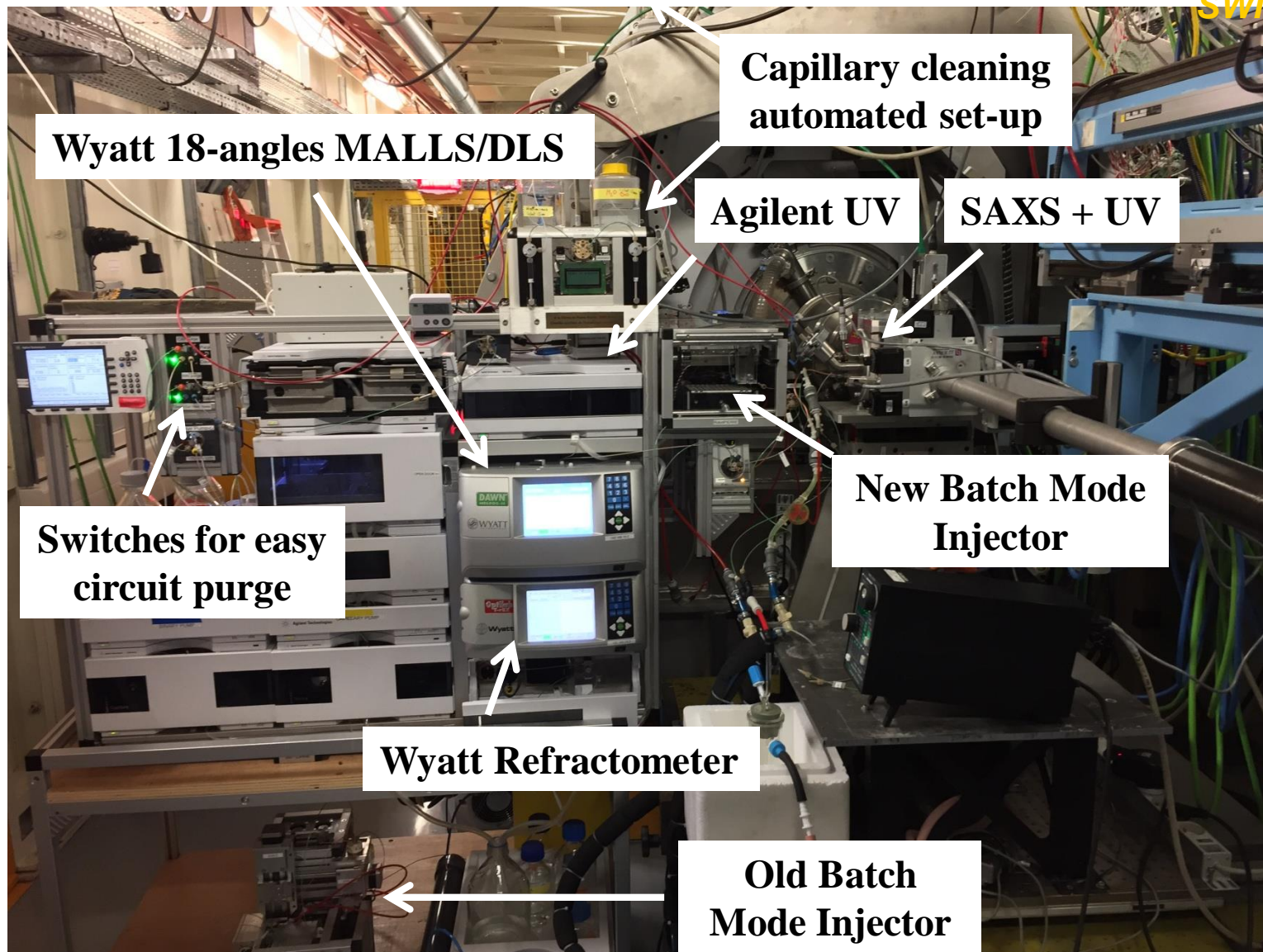


Samples for BATCH-SAXS:
10 μ l (>1 mg/ml) to 50 μ l for (>0.1 mg/ml)
Less than 3 minutes per sample.

Automated process for injection and detection (passerelle GUI):

- Automated injection for both Batch and SEC-SAXS
- Users just need to define a list of: Sample names, Rack Positions, Volumes and Column circuits
- Multiple injections for both Batch and SEC-SAXS can be programmed

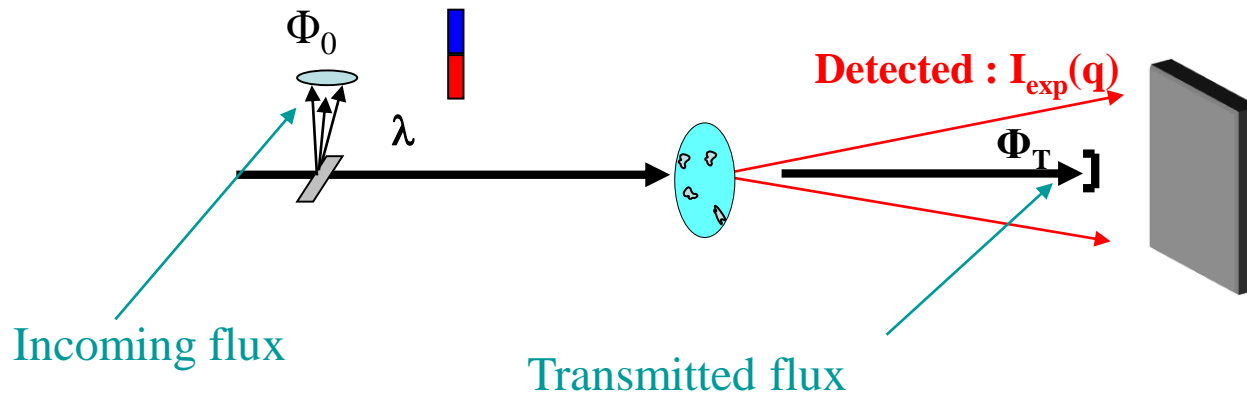
⇒ 46 proteins recorded using only SEC-SAXS in 21 hours.



Transmission and buffer measurements are crucial

- Transmission

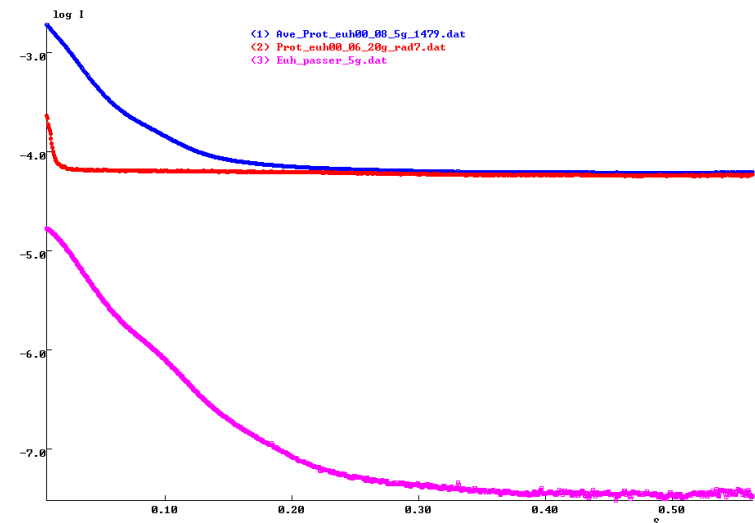
- The experimental scattering intensity must be normalised by transmitted intensity.
- Transmission intensity must be measured with high accuracy ($\sim 0.1\%$).



- Buffer

- Buffer and protein samples must be measured in the same cell for correct subtraction of parasitic background arising from slits and holder walls.
- The buffer in the buffer sample must be identical to that of the protein sample (dialysis, SEC, ...).

$$I_{\text{particles}}(q) = I_{\text{sample}}(q) - I_{\text{buffer}}(q)$$



Calibration of the set-up using water scattering

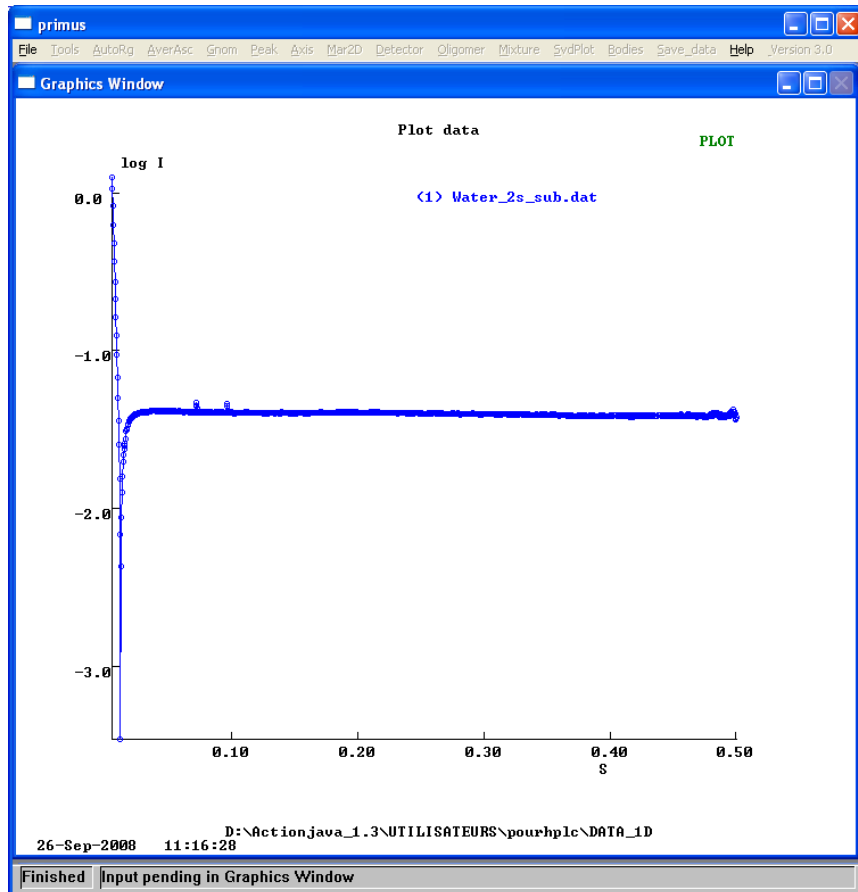


Liquid scattering (theory): $I(q) = \text{constant at small } q = r_0^2 Z^2 \rho_A^2 \cdot kT \kappa_T$

$$I_{\text{H}_2\text{O}, \text{theory}} = 0.0163 \text{ cm}^{-1}$$

Molecular density

Isothermic compressibility



Water is used as primary reference to get the absolute intensity scale

- Capillary diameter = 1.6 mm
- Average of 2 frames of 2s
- Empty capillary subtracted
- Normalized by solid angle
- Normalized by transmitted intensity

Example:

$$I_{\text{H}_2\text{O}, \text{exp}} = 0.042 \text{ Exp. Units}$$

$$I_{\text{H}_2\text{O}, \text{exp}} = K_{\text{exp}} * I_{\text{H}_2\text{O}, \text{theory}}$$

$$\rightarrow \text{Here : } K_{\text{exp}} = 2.56 \text{ Exp.Units / cm}^{-1}$$

➔ For any sample in that capillary : $I_{\text{theory}}(\text{cm}^{-1}) = I_{\text{exp}} / K_{\text{exp}} = I_{\text{exp}} / 2.56$

Particles in solution

Relation between the the number of measured photons ΔN_{ph} on a given pixel of the detector, making a solid angle $\Delta\Omega$, and the Scattering Intensity per unit volume :

Differential cross-section

Number of detected scattering photons
in a given pixel

$$I(q) = \frac{1}{V} \frac{d\sigma}{d\Omega}(q) = \frac{\Delta N_{\text{ph}}}{N_0} \frac{1}{T \cdot e} \frac{D^2}{PxSize^2}$$

Distance sample-pixel

Irradiated volume

Sample thickness

Sample transmission

Number of incident photons

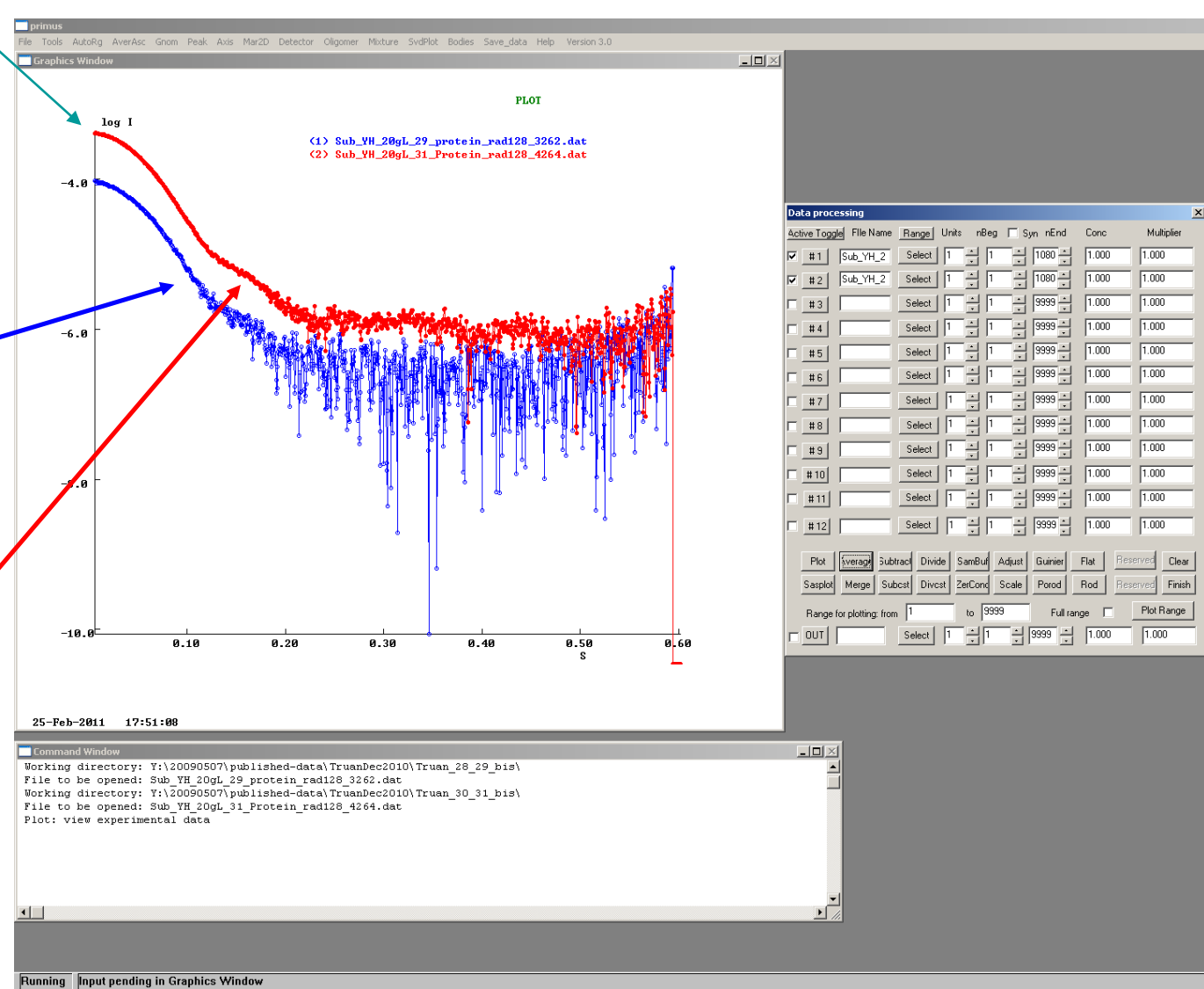
Scattering Intensity per unit Volume

Merge low c (@ low q) and high c (@ high q) curves

Here, slight repulsive interactions alter the concentrated curve at small angles

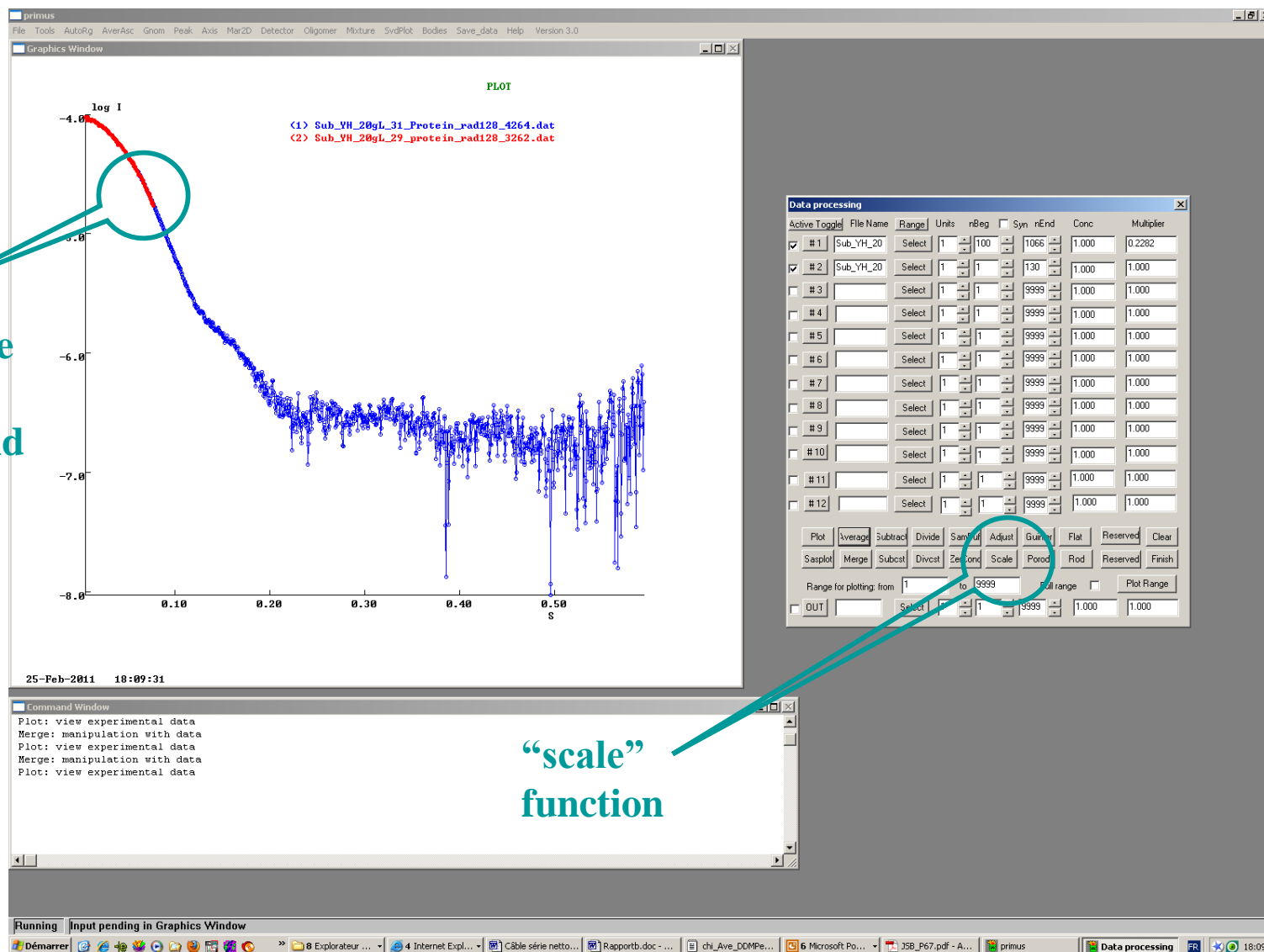
small angle data using the lowest concentration curve or an extrapolation to zero concentration from a series of dilute solutions (correction of interparticle effects)

larger angle data using the most concentrated solution



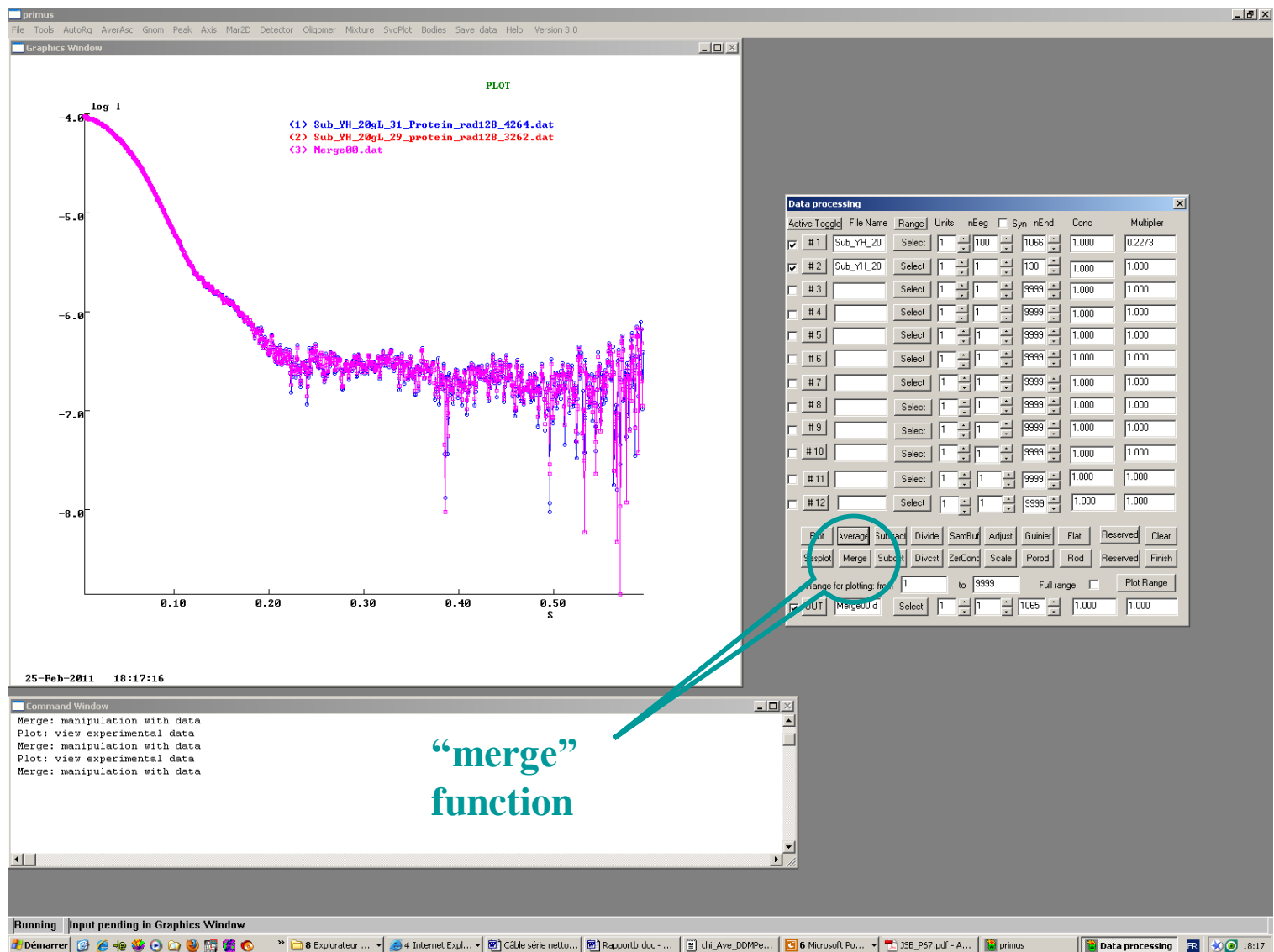
PRIMUS: merging data

The common range should be as restricted as possible to avoid adding noise



“scale”
function

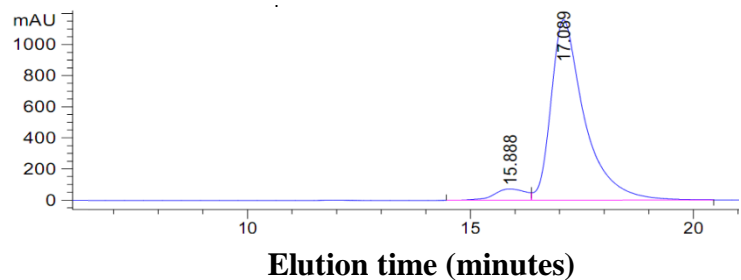
PRIMUS: final merged curve



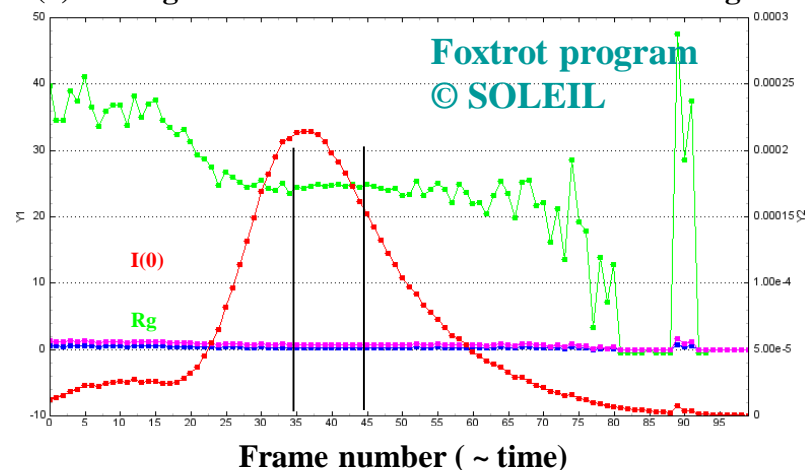
2nd case : the solution is a slow equilibrium or an unwanted mixture

- Use on-line HPLC data collection (typ 50 μ l)

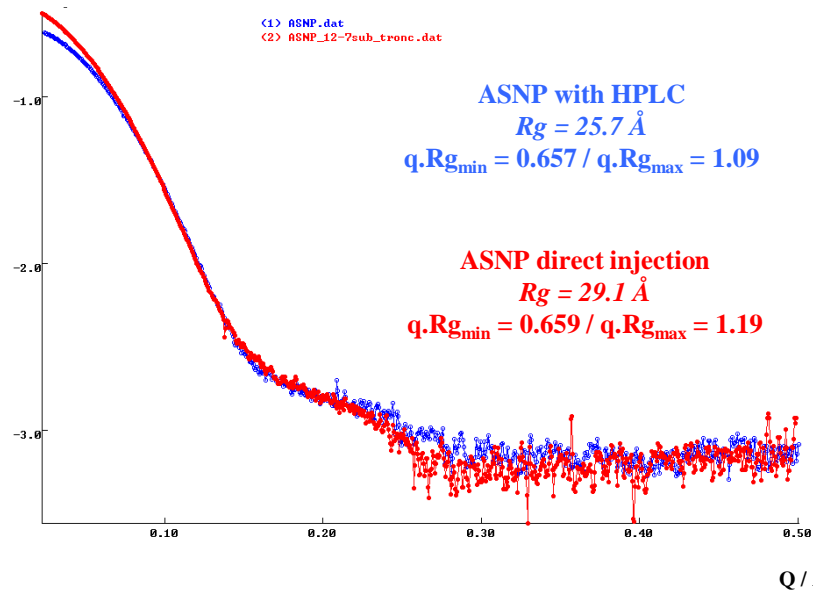
ASNP elution profile, monitored by UV absorption at 280 nm



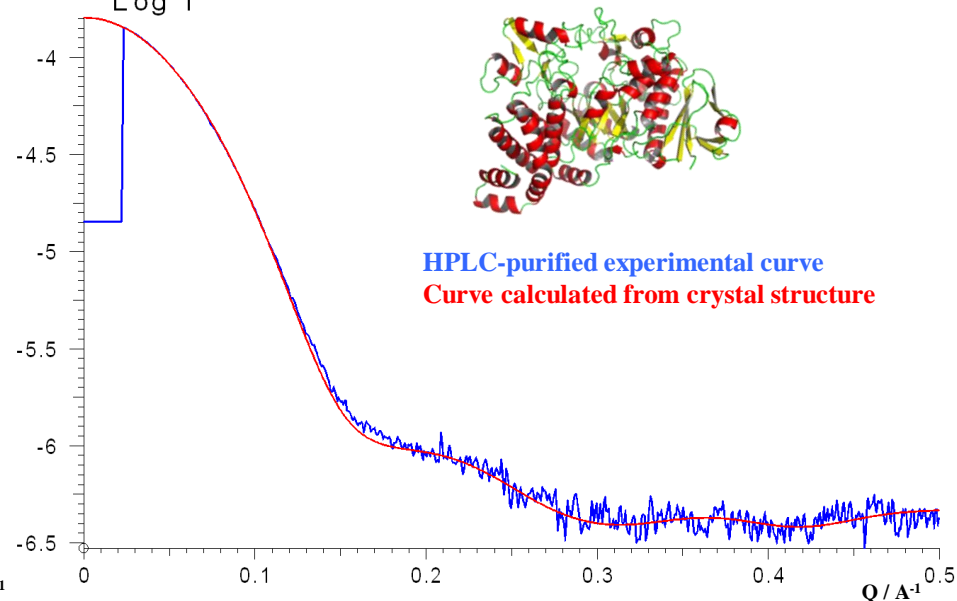
$I(0)$ and R_g determined for each SAXS frame during elution



Comparison between HPLC-purified and Direct injection curves

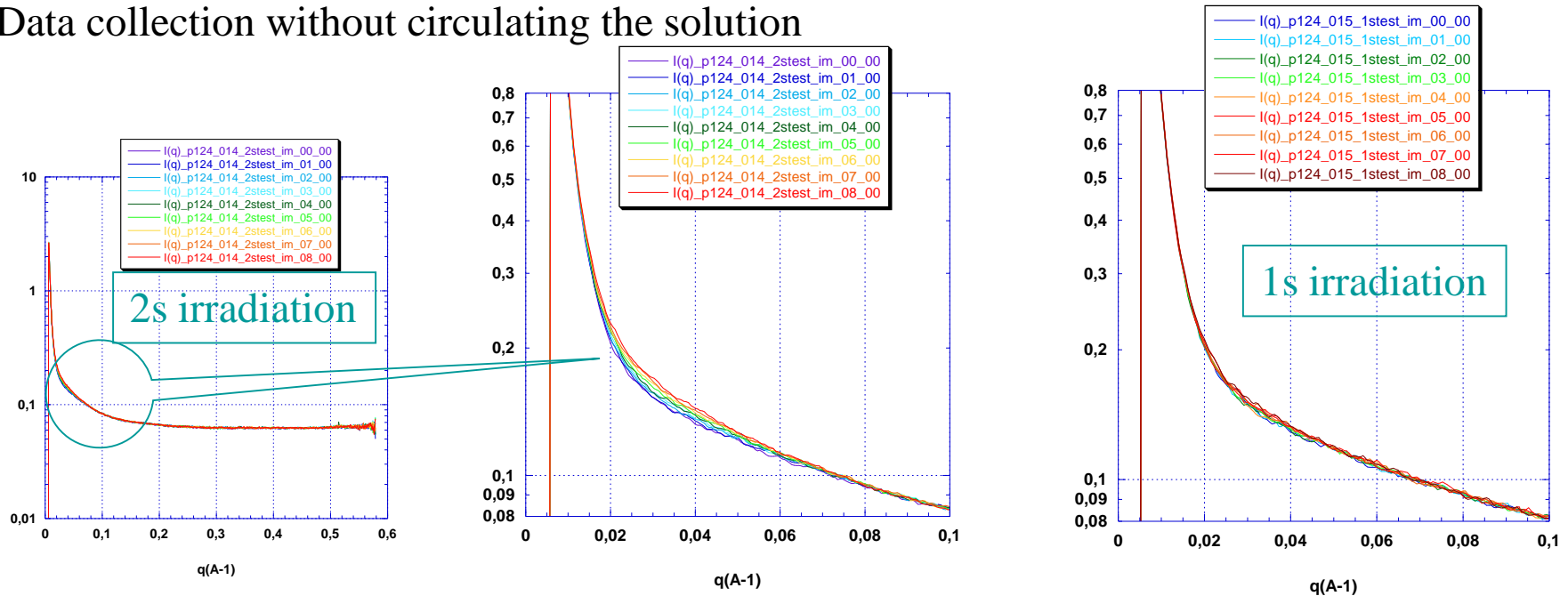


Fitting the HPLC-purified experimental curve with the crystal structure

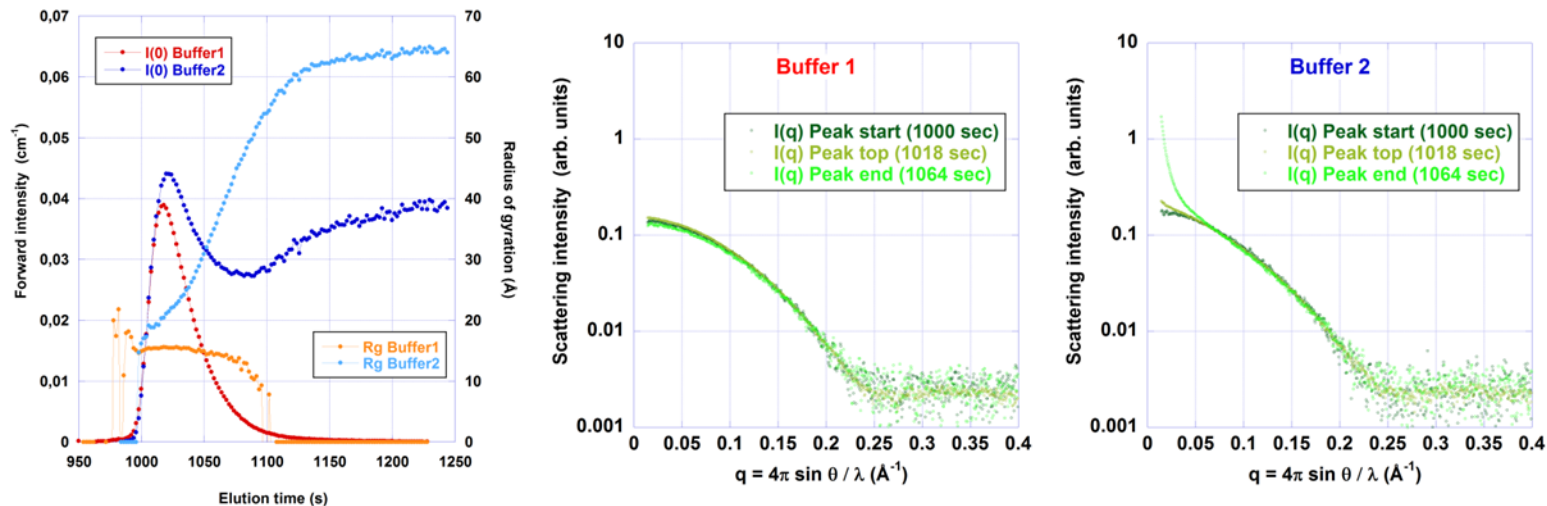


Unwanted Radiation effects

Data collection without circulating the solution



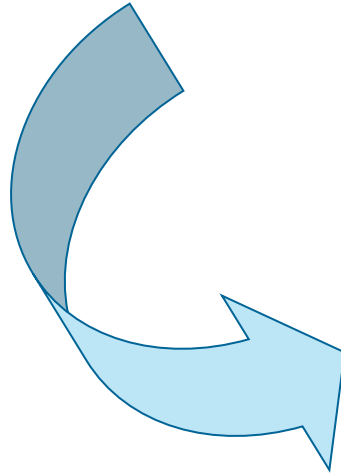
Even with a circulating solution (SEC-SAXS here)



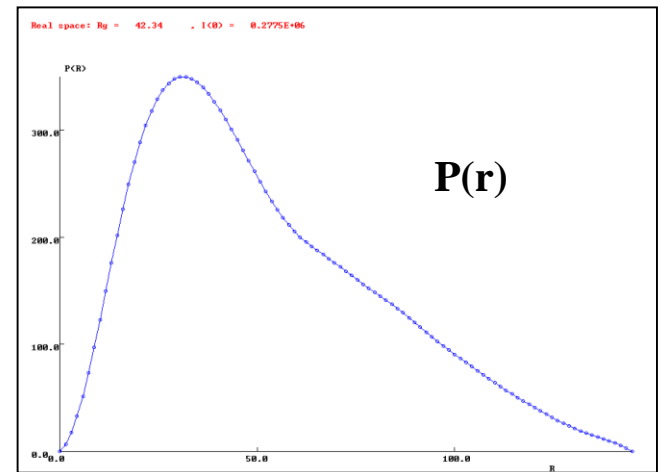
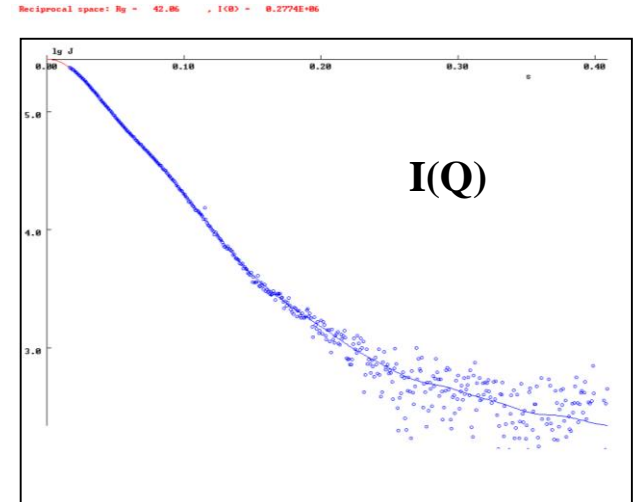
Pérez J., Vachette P. (2017)
In: Biological Small Angle
Scattering: Techniques,
Strategies and Tips.
Advances in Experimental
Medicine and Biology, vol
1009. Springer, Singapore

At this stage

We have gone from

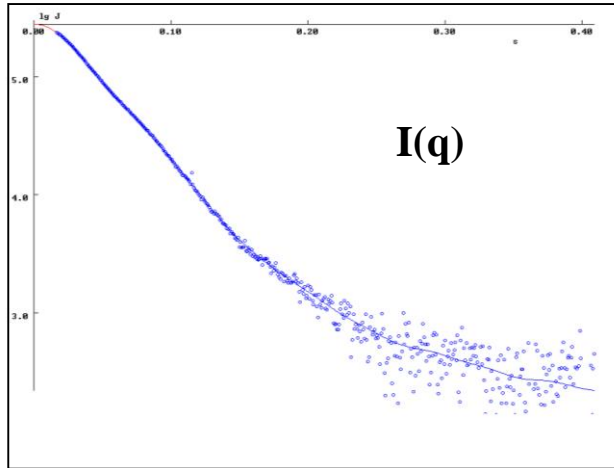


to

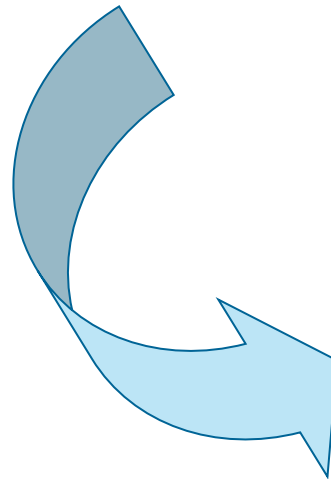
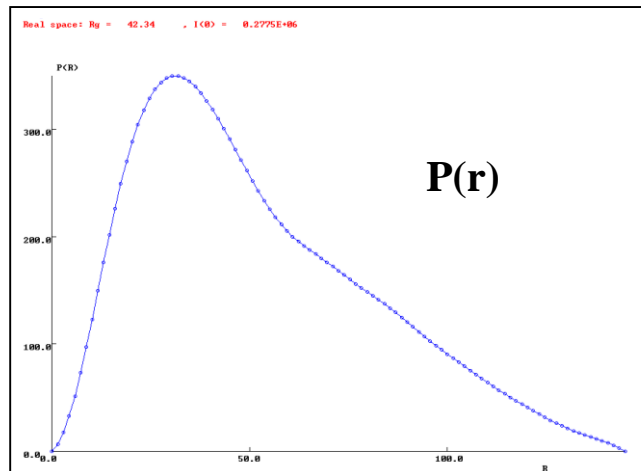


Now, we have to go from

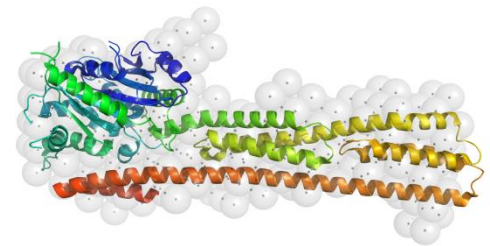
Reciprocal space: $R_g = 42.86$, $I(0) = 0.2774E+06$



Real space: $R_g = 42.34$, $I(0) = 0.2775E+06$

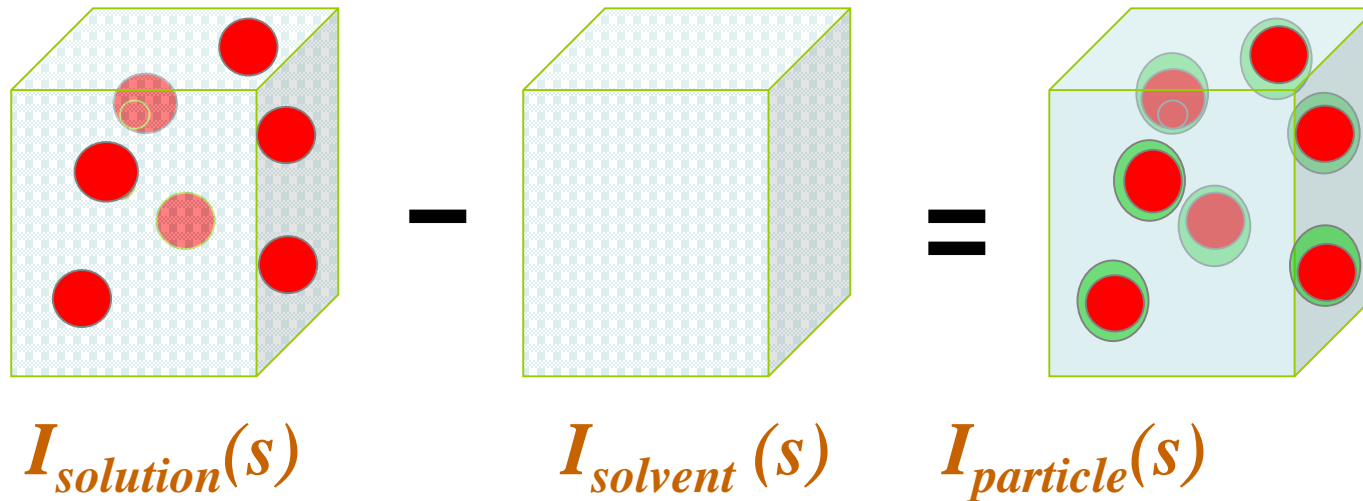


to



MODELLING

Solvent scattering and contrast



To obtain scattering from the particles, solvent scattering must be subtracted to yield the effective density distribution $\Delta\rho(\mathbf{r}) = \rho(\mathbf{r}) - \rho_0$, where ρ_0 is the scattering density of the solvent.

Further, the bound solvent density may differ from that of the bulk.

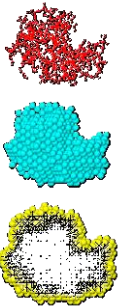
CRY SOL: from atomic coordinates to a SAXS curve

$$I_{th}(q) = \left\langle \left| A_a(\vec{q}) - \rho_s A_s(\vec{q}) + \delta\rho_b A_b(\vec{q}) \right|^2 \right\rangle_{\Omega}$$

$A_a(q)$ = molecular scattering amplitude in vacuum

$A_s(q)$ = scattering amplitude from excluded volume

$A_b(q)$ = scattering amplitude from the hydration shell, layer of arbitrary thickness 3\AA



In CRY SOL program, in order to gain computing time, $I(q)$ is developed in a series of Bessel functions and spherical harmonics :

$$I_{calc}(q) = \sum_{l=0}^L \sum_{m=-1}^l \left| A_{lm}(q) - \rho_0 C_{lm}(q) + \delta\rho B_{lm}(q) \right|^2$$

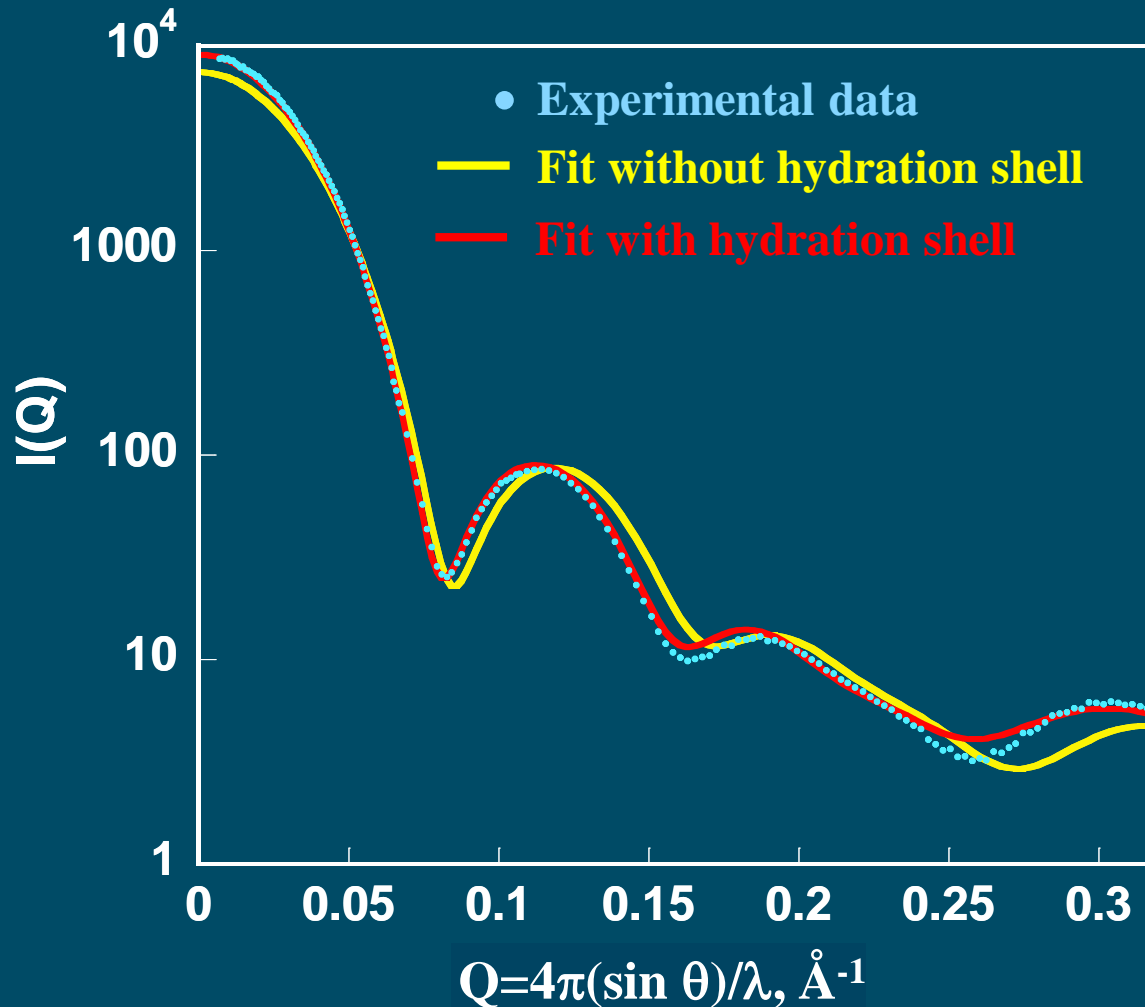
The experimental scattering curves are then fitted using only 3 parameters in order to minimize the discrepancy χ :

- the general scale of $I_{calc}(q)$
- the total excluded volume V , which is equivalent to modifying the average contrast ρ_0
- the contrast of the border layer $\delta\rho$

$$\chi^2 = \frac{1}{N-1} \sum_{i=1}^N \left[\frac{I_{exp}(q_i) - scale * I_{calc}(q_i)}{\sigma_{exp}(q_i)} \right]^2$$

Effect of the hydration shell

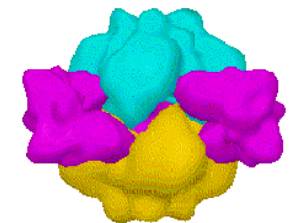
T state of *E. coli* allosteric ATCase



y
z



A



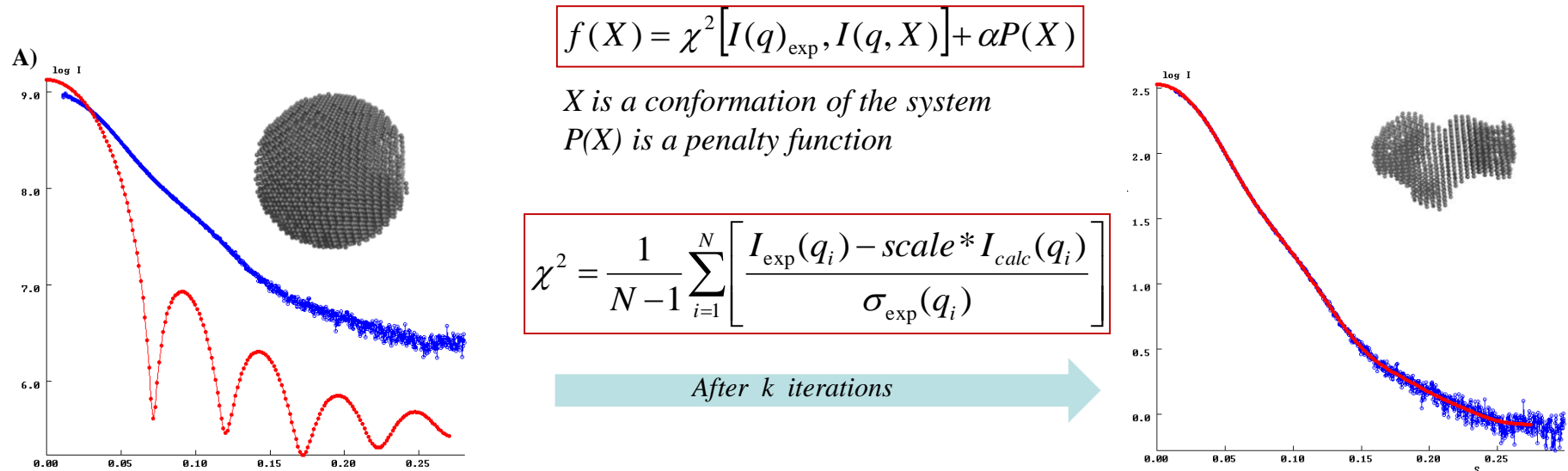
Ab initio shape modelling : nothing is known excepted the curve !

Principle of the method : any structure volume of homogeneous electronic density can be approximated at any resolution by a set of spheres of small enough diameter

Starting model = sphere with a radius $R = D_{\max}/2$ with N scattered beads ($r_0 \ll R$)

The number of the “dummy atom” $N \approx (R/r_0)^3$

Each sphere is associated to a position j and an index X_j corresponding to the type of the phase ($X_j = 0$ for the solvent and $X_j = 1$ for the molecule)



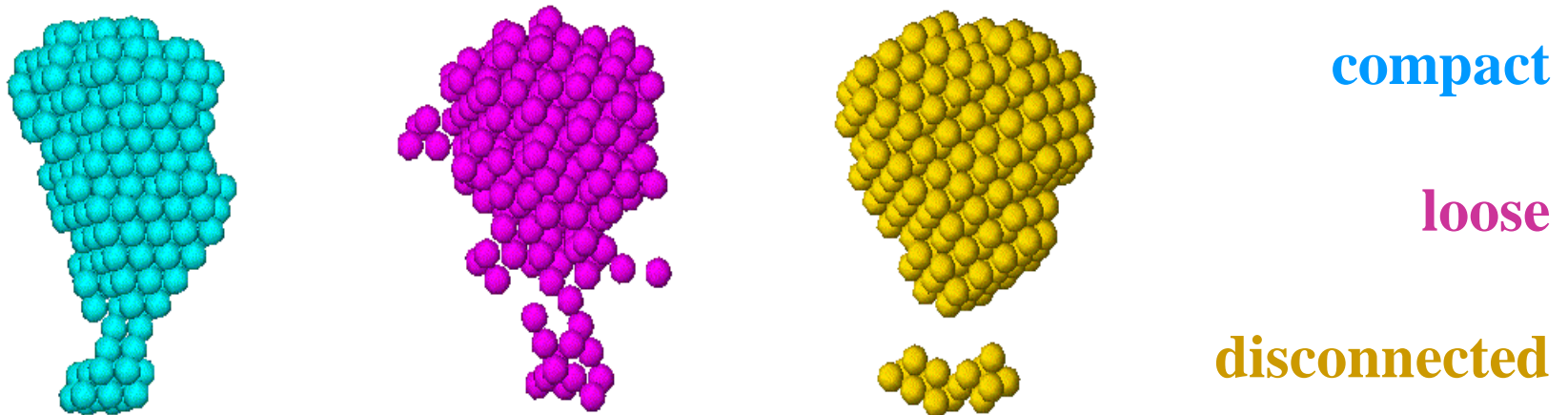
D. I. Svergun, M. Kozin, M. Petoukhov, V. Volkov (1999). *Biophys J.* 2879-2886.

3D shape reconstructions from SAXS data with DAMMIN

- Obtaining 3D shapes from SAXS data is a ill-defined problem that can be ****partially**** solved by introducing additional information to ****reduce**** ambiguity of interpretation
- Using simulated annealing, finds a compact dummy atoms configuration X that fits the scattering data by minimizing

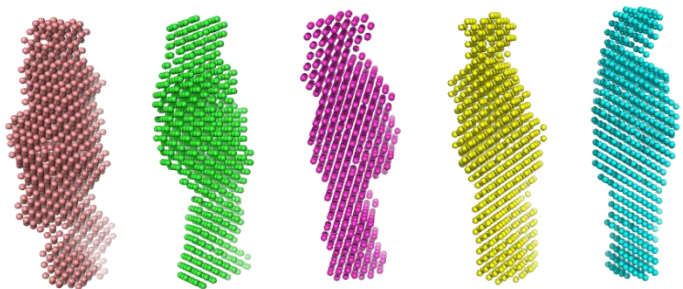
$$f(X) = \chi^2[I_{\text{exp}}(s), I(s, X)] + \alpha P(X)$$

where χ is the discrepancy between the experimental and calculated curves, $P(X)$ is the penalty to ensure compactness and connectivity, $\alpha > 0$ its weight.



3D shape reconstructions from SAXS data with DAMMIN

- A series of runs (10-50) are performed to compare the different shapes obtained from the same data.
- After the run, an optimal superposition of models is realized with the program suite DAMSEL and DAMSUP.
- The algorithm defines a criteria of similarity, called « Normalized Spatial Discrepancy » or NSD, which measures the agreement between any pair of models.
- Similar shapes results in $NSD < 1$, very similar shapes $NSD \approx 0.5$



Shp1 Shp2 Shp3 Shp4 Shp5

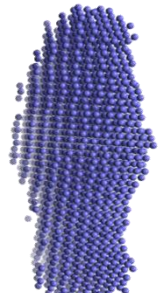
File	Aver	1	2	3	4	5
1	0.52	0.00	0.51	0.52	0.50	0.52
2	0.52	0.51	0.00	0.52	0.49	0.52
3	0.53	0.52	0.52	0.00	0.53	0.52
4	0.53	0.50	0.49	0.53	0.00	0.54
5	0.53	0.52	0.52	0.52	0.54	0.00

Mean value of NSD : 0.535
Standard deviation of NSD : 0.008

Damsel.log



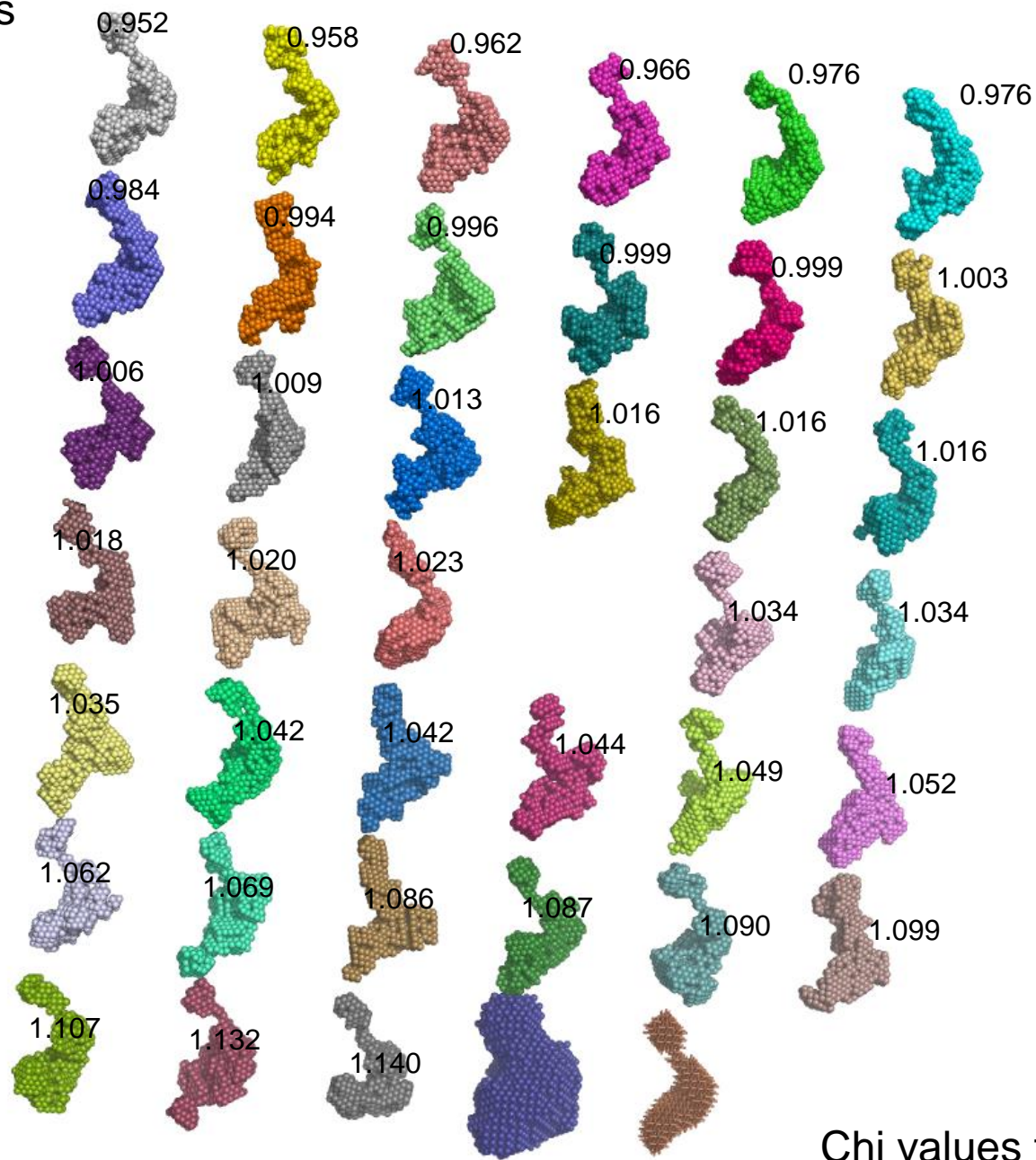
Damfilt
(intersection)



Damaver
(all superimposed)

- Models are conserved if its $NSD < \text{Mean of NSD} + 2 * \text{standart deviation}$
- The model with the lowest NSD is the shape which has the most similarities with other, and ****can**** be regarded as the most representative of envelopes in accordance with the SAXS data
- Be careful with damfilt.pdb because $I_{\text{damfilt}}(q) \neq I_{\text{exp}}(q)$

NSD values



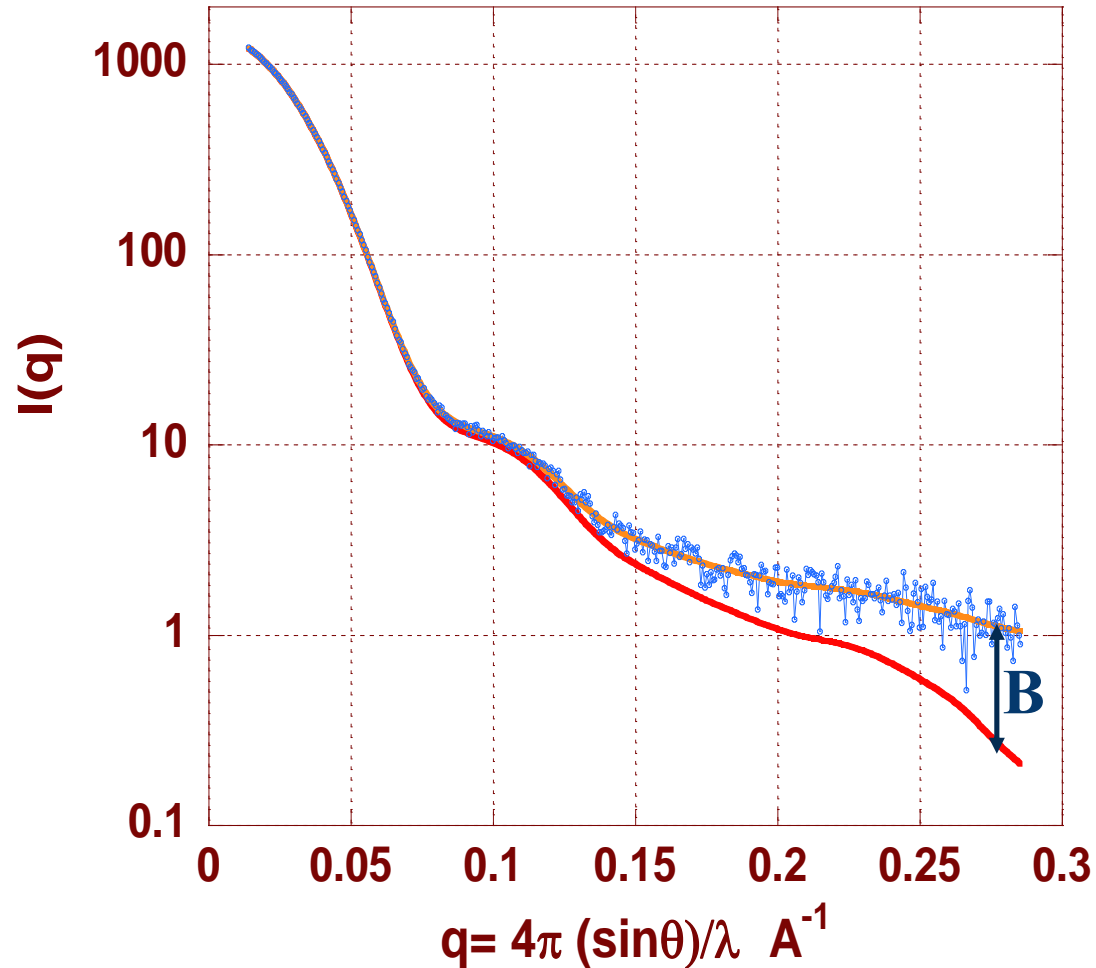
Chi values from 1.704 to 1.942

Be aware : “Porod law” is forced for ab initio shape determination

DAMMIN : shape determination
Model with uniform density



Fitting data with approximate q^{-4}
high angle trend by subtracting a
constant.



SASREF: Rigid body modeling against SAXS data

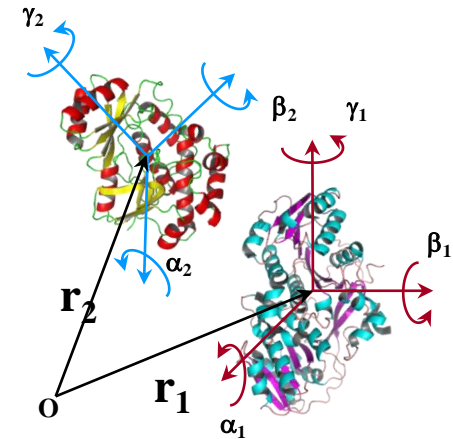
SASREF : when atomic structures of domains are known, but no their mutual organization

The objective is to find the relative orientation of each subunit with a correct agreement with the SAXS data of the complex

The scattering intensity $I(q)$ of the complex is equal to the sum squared of the amplitudes of each subunit

$$I(S) = \left\langle \left| \sum_{k=1}^K A^{(k)}(\vec{S}) \right|^2 \right\rangle_{\Omega}$$

$$A^{(k)}(\vec{S}) = \exp(i\vec{S} \cdot \vec{r}_k) \prod (\alpha_k \cdot \beta_k \cdot \gamma_k) [C^{(k)}(\vec{S})]$$



The amplitude are calculated with CRY SOL from the high resolution structure of each monomer

The algorithm of minimization is the same used with DAMMIN with a penalty function (interconnectivity of the subunits, the steric clashes) and possibility to give information about contacting residues from other experiences.

$$f(X) = \sum_i \chi_i^2 + \alpha_{dist} P_{dist}(X) + \beta_{cross} P_{cross}(X) + \gamma_{cont} P_{cont}(X)$$

DADIMODO : All-atom rigid body + linkers refinement vs. SAXS / NMR data

Collab : Christina Sizun & François Bontems (ICSN, Gif sur Yvette))

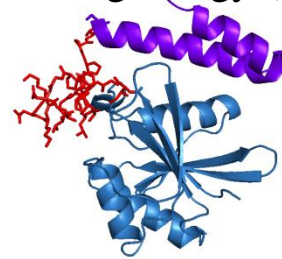
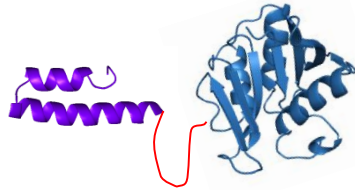
F. Mareuil, et al. (2007) *Eur Biophys J.*

Evrard et al. (2011), *J. Appl. Cryst.*

Modelling approach : complete atomic model

Full structure initiated with :

- Crystal or NMR domain structures
- Homology models



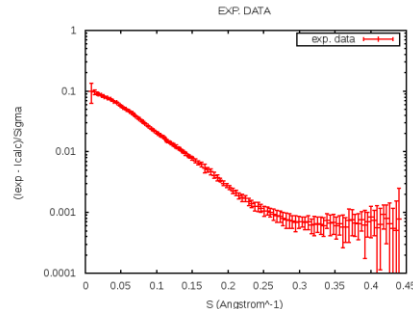
External information:

- Sequence
- Sub-parts moved as rigid-bodies (user-defined)
- A correct stereochemistry is maintained at all steps by minimizing energy (Amber 99 Force Field)

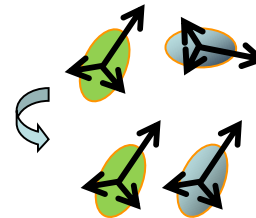
Experimental data:

- SAXS
- NMR
- RDC

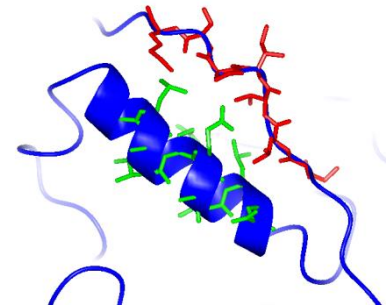
ADR (chem. shift map.)



SAXS score



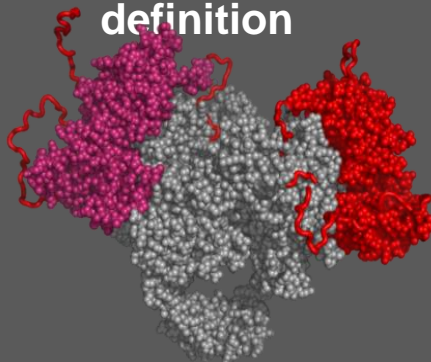
RDC score



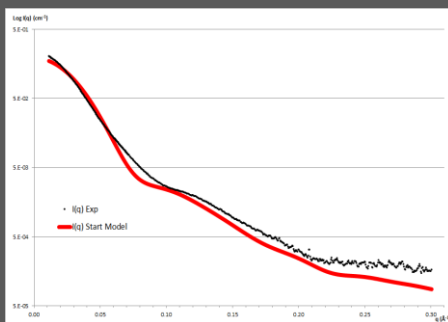
ADR score

Optimisation of the structure via a genetic algorithm

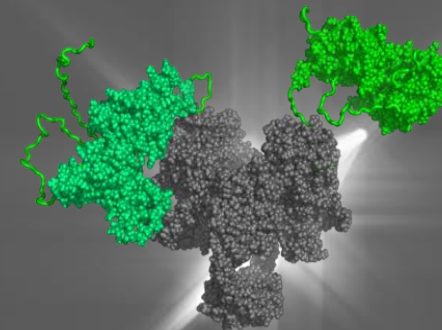
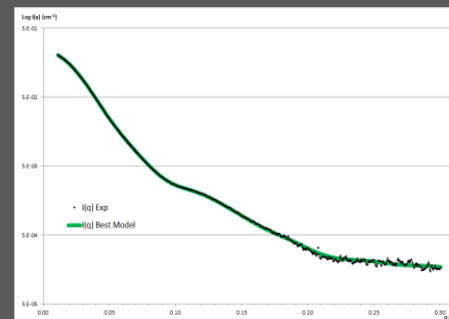
Full atom Pdb
file +
rigid domains
definition



Saxes curve



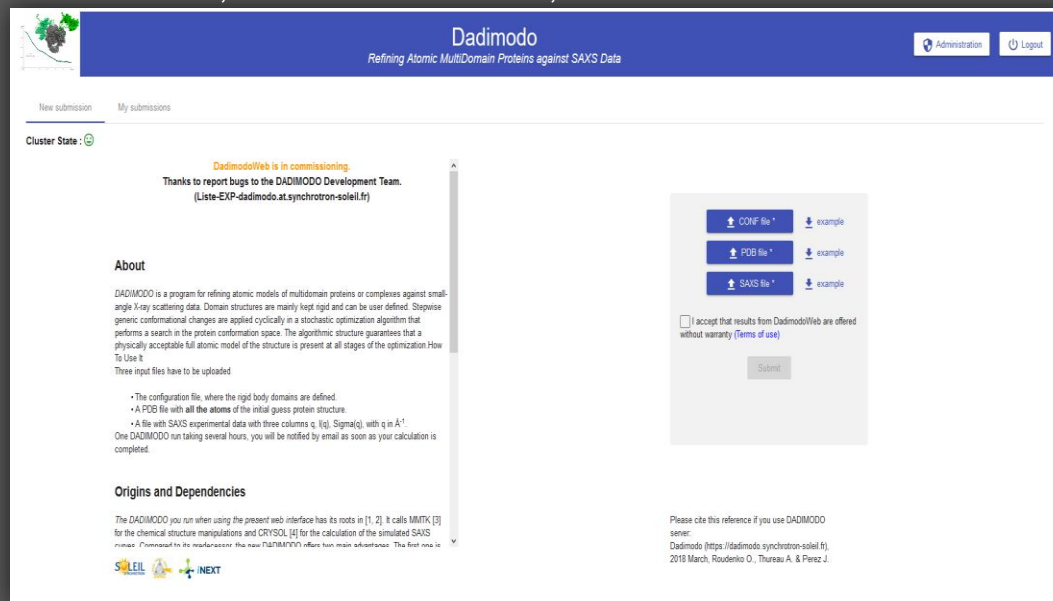
Petrella et al., 2019, Structure, *in press*

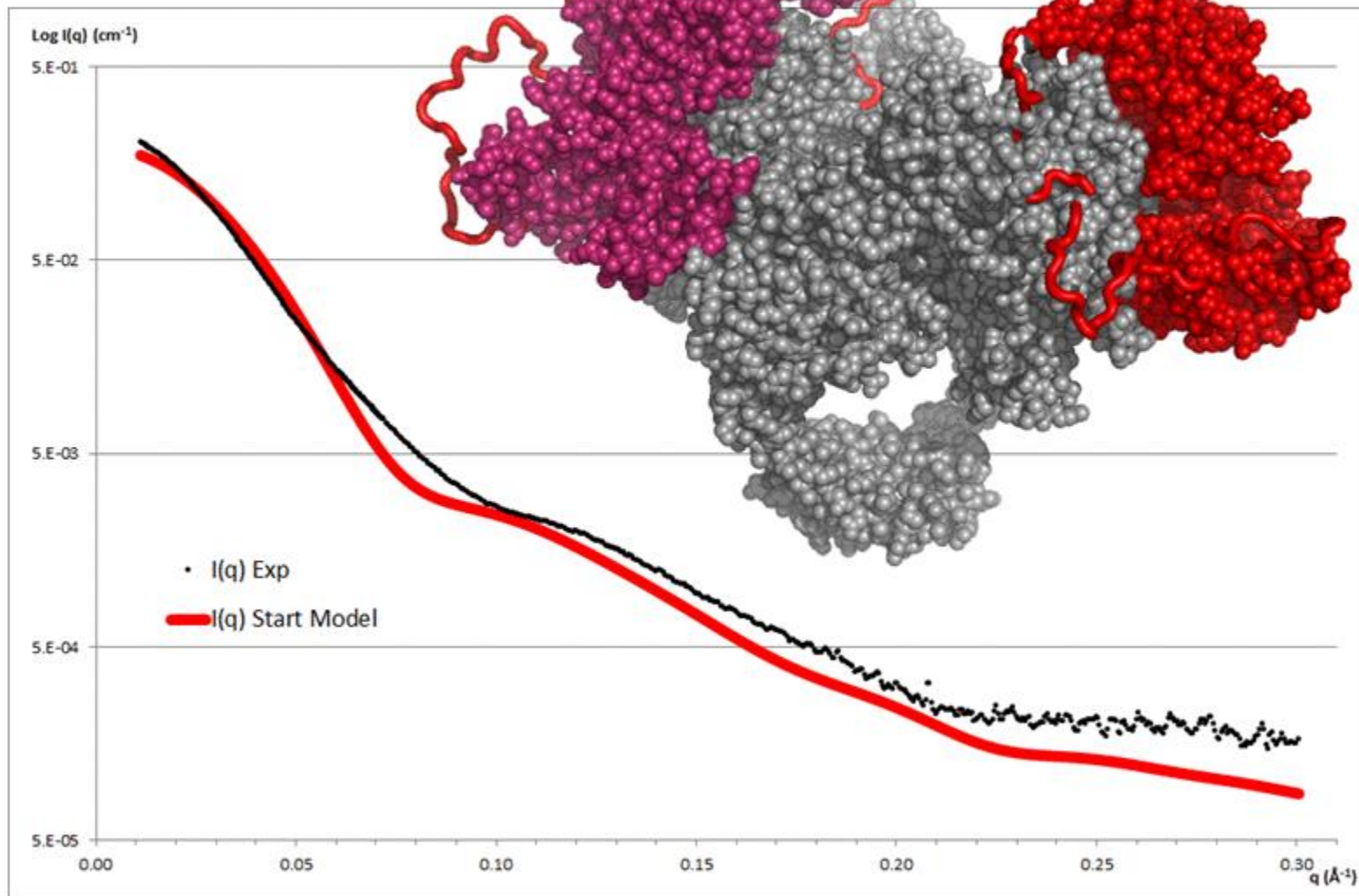


5 SAXS compatible models

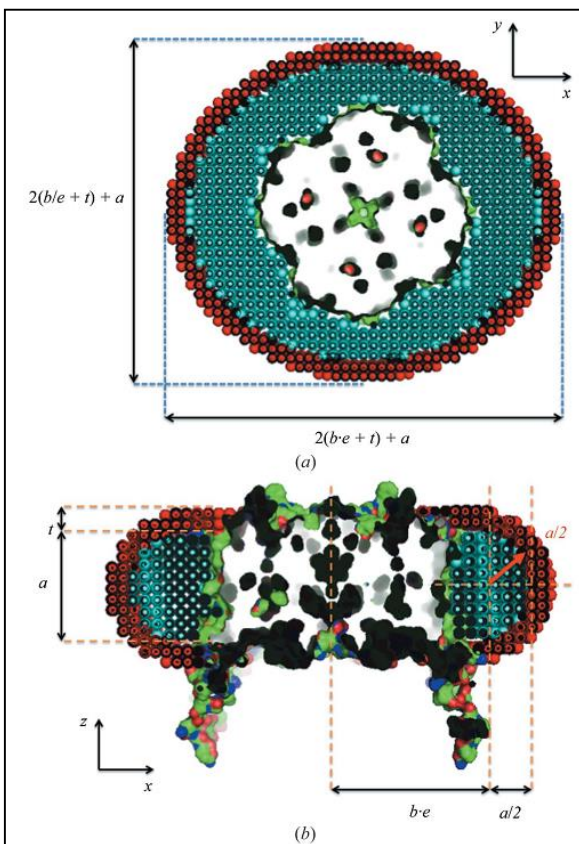
<https://dadimodo.synchrotron-soleil.fr>

Roudenko O., Thureau A. & Pérez J., March 2018

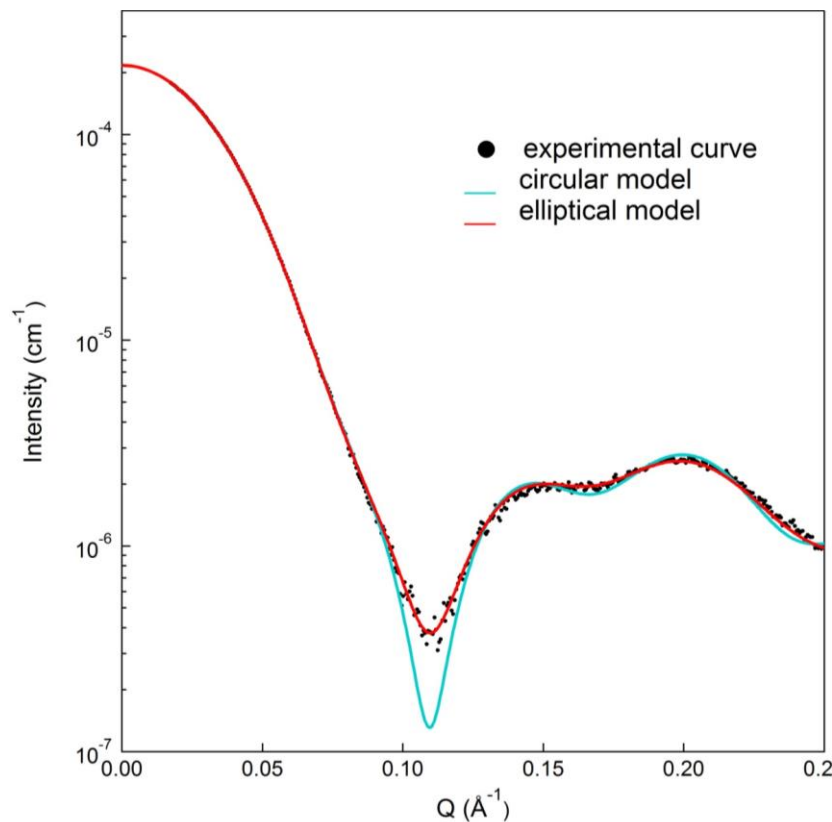




Pérez J. & Koutsioubas, A. (2014). Acta Cryst.D70



Algorithm of the *Memprot* program. The program essentially creates PDB files with the models made of the full-atom protein structure and the parameterized coarse-grained detergent corona, and *CRY SOL* is called to calculate the SAXS curves. An overall sorting on the χ value is performed to keep the best model.



Berthaud et al. (2012), *JACS*, 134 (24), 10080-10088

A FEW CONCLUSIONS

- A scattering pattern can be accurately calculated from atomic coordinates, thereby providing a link between high resolution and BioSAXS work.
- Using SAXS patterns, *ab initio* methods can propose possible shapes of a macromolecule
- Based on all-atom partial structural information, stochastic algorithms may provide possible models for missing regions
- Analysis and modeling require a monodisperse and ideal solution, which has to be checked independently.

SAXS is at his best when it is used to distinguish between several preconceived hypotheses.