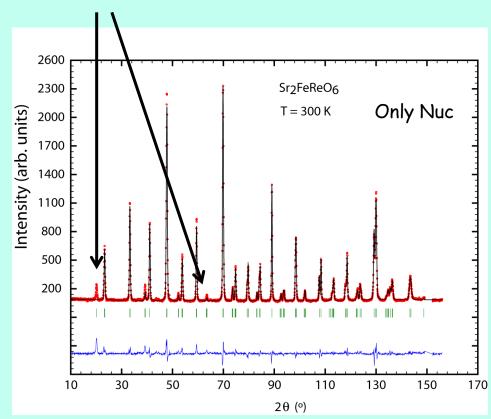
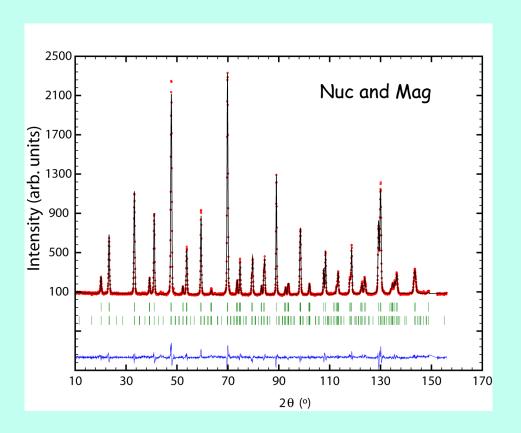
#### Magnetic Scattering, how can we detect it's existence? (1)



E.g.: Refining a neutron diffraction data set using a model for the crystallographic structure which is known from x-ray data

Strong intensity not accounted for by the model on allowed Bragg peaks or in additional Bragg peaks.





Missing intensity stronger at lower scattering angles ---> magnetic scattering

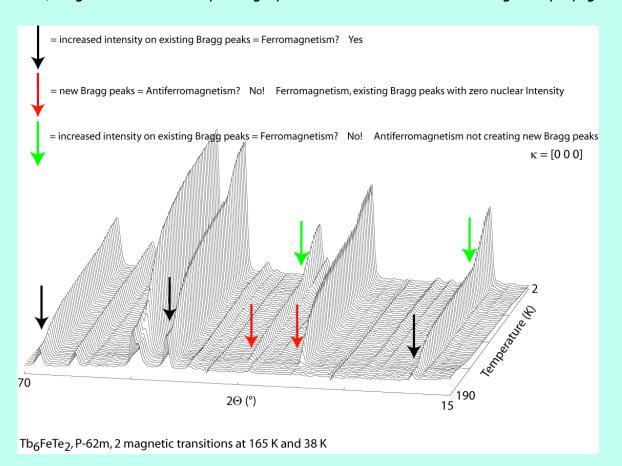
For room temperature data only if transition metals (Fe, Co) are present in the compound.

#### Magnetic Scattering, how can we detect it's existence? (2)

Measuring the temperature dependence (thermodiffractogram) of the neutron spectrum and look for changes.

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- I) Ferromagnetism, Ferrimagnetism: Intensity is added to existing Bragg peaks
- II) Antiferromagnetism:
- 1) Magnetic unit cell = Crystallographic unit cell, magnetic propagation vector  $\kappa$  = [0 0 0], new Bragg peaks can be created
- 2) a) Magnetic unit cell  $\neq$  Crystallographic unit cell, commensurate magnetic propagation vector, e.g.  $\kappa = [\frac{1}{2} \ \frac{1}{2} \ 0]$ , new Bragg peaks
  - b) Magnetic unit cell  $\neq$  Crystallographic unit cell, incommensurate magnetic propagation vector, e.g.  $\kappa$  = [0  $\frac{1}{2}$  0.21], new Bragg peaks



Exclude structural transitions as reasons for changes:

Peak splittings point to structural transitions through distortions.

Appearence of new peaks is accompanied by an intensity reduction of old peaks.

#### But:

Magnetostructural transitions where the magnetic phase transition triggers a structural transition or vice versa.

When does antiferromagnetism **not** create new reflections with  $\kappa = [0\ 0\ 0]$ ?

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e.g.: P2/m, Pmmm, P-62m, P432

When does antiferromagnetism Create new reflections with  $\kappa = [0\ 0\ 0]$ ?

e.g.: P2<sub>1</sub>/m, Pnma

#### In general:

Symmetry elements which create systematic extinctions within the crystallographic unit cell:

Lattice centering: A, B, C, F, I, R (e.g.:  $I \longrightarrow +\frac{1}{2} \frac{1}{2} \frac{1}{2}$ )

Screw axes:  $2_1$ ,  $3_1$ ,  $4_1$ ,  $4_2$ ,  $6_1$ ,  $6_2$ ,  $6_3$  (Rotation + Translation)

Glide planes: a, b, c, n, d (Reflection + Translation)

In structures where the crystallographic unit cell contains screw axes or glide planes a magnetic propagation vector  $\kappa$  = [0 0 0] creates new magnetic Bragg reflections for antiferromagnetic structures.

Structures having lattice centering like I4/m, R-3 create no new reflections for  $\kappa$  = [0 0 0], spins are parallel.

If the centering condition is lost, propagation vectors like  $\kappa$  = [1 0 0] or  $\kappa$  = [0 0 1.5] result and new reflections are formed, spins are anti-parallel.



#### Aim: Efficient use of neutron diffraction data to solve magnetic structures

Ideal diffraction data for the determination of magnetic structures:

- 1) High resolution neutron data of the compound in the paramagnetic phase: Initial structure
- 2) High intensity neutron thermodiffractogram from  $T > T_C$  or  $T_N$  down to lowest possible T:

Appearence of magnetic scattering and its T-dependence Existence of different transitions? Different magnetic propagation vectors?

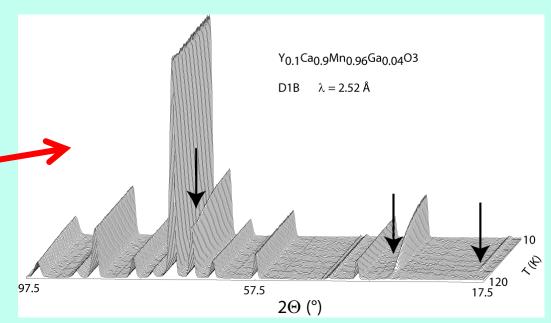
3) High resolution neutron data (and synchrotron data) at lowest T: Final structure, Existence of a structural phase transition?

Example 1:  $Y_{0.1}Ca_{0.9}Mn_{0.96}Ga_{0.04}O_3$ 

Temperature dependence of the neutron diffraction data:

#### Thermodiffractogram

Conclusion: 1 magnetic transition at about 95 K





1) Determination of the structure in the paramagnetic phase

Paramagnetic.pcr

Paramagnetic.prf

2) Determination of the magnetic propagation vector

Structure10K.pcr Structure10K.prf

Difference pattern

Ksearch1.sat Ksearch.kub

$$\kappa = (\frac{1}{3}, \frac{1}{3}, 0)$$

$$\kappa = (0, 0, 0)$$



#### Finding a magnetic structure by trial and error:

 $\kappa = [0\ 0\ 0] \rightarrow \text{magnetic unit cell} = \text{nuclear unit cell}$ 

Refine the magnetic contribution using a second phase in the FULLPROF pcr file

Give all the magnetic atoms separately and imagine the magnetic coupling between the magnetic moments in the cell: Mn on  $(0,0,\frac{1}{2})$ ,  $(\frac{1}{2},0,0)$ ,  $(0,\frac{1}{2},\frac{1}{2})$ ,  $(\frac{1}{2},\frac{1}{2},0)$ 

	!! ! Data for PHASE number: 2 ==> Current R_Bragg for Pattern# 1: 4.50	Couplir
	!LAYCA magnetic	
	4 0 0 0.0 0.0 1.0 1 0 1 0 0 0.000 0 3 0	Site
	! P -1	1
Site	!S1Î S12 S13 T1	2
1	!Atom Typ Mag Vek X Y Z Biso Occ Rx Ry Rz ! Ix	3
2	0.00 0.00 0.00 0.00 241.00 0.00 311.00 291.00 321.00 0.00 0.000 0.000 0.000 0.000 0.000 0.000 0.0000 0.00000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.00000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.000000	4
3	0.00	
4	0.000 0.000 0.000 0.000 0.000 0.000 0.000 0.000 0.0000 0.000	
	0.00 0.00 0.00 0.00 241.00 0.00 311.00 291.00 321.00 0.000 0.000 0.000 0.000 0.000 0.000 0.0000 0.00 0.00 0.00 0.00 0.00 0.00 0.00 1> Profile Parameters for Pattern # 1	

ings within a 4-fold site:

Here we assume e.g.	Site		Coup	oling		
a ferromagnetic coupling in all three	1	+	+	+	+	
directions:	2	+	+	-	-	
$F_x F_y F_z$	3	+	-	-	+	
Trial_and_Error1.pcr	4	+	-	+	-	
Trial_and_Error.prf		F	Α	G	С	
Trial_and_Error1.new						
\\\\\EullDnof Suit	ta\winnla					

Resulting fit is very bad!!



#### Trial and error 2 (or 41 or 99 or ...)

! Data for PHASE number: 2 ==> Current R_Bragg for Pattern# 1: 4.50										
LAYCA magnetic										
!Nat Dis Mom Pr1 Pr2 Pr3 Jbt Irf Isy Str Furth										
4 0 0 0.0 0.0 1.0 1 0 1 0 0 0.000 0 3 0										
P -1 <space for="" generation<="" group="" hkl="" symbol="" td=""></space>										
!Nsym Cen Laue MagMat										
1 1 1 1										
S11 S12 S13										
!M11 M12 M13 M21 M22 M23 M31 M32 M33 Ph										
$\begin{array}{cccccccccccccccccccccccccccccccccccc$										
!Atom Typ Mag Vek X Y Z Biso Occ Rx Ry Rz										
! IX IY IZ beta11 beta22 beta33 MagPh										
MN1 MMN3 1 0 0.00000 0.00000 0.50000 0.20484 0.96000 2.000 2.000 2.00										
0.00 0.00 0.00 241.00 0.00 311.00 291.00 321.0										
0.000 0.000 0.000 0.000 0.000 0.000 0.0000										
0.00 0.00 0.00 0.00 0.00 0.00 0.00										
MN1 MMN3 1 0 0.50000 0.00000 0.00000 0.20484 0.96000 2.000 -2.000 2.00										
0.00 0.00 0.00 241.00 0.00 311.00 -291.00 321.0										
$egin{array}{cccccccccccccccccccccccccccccccccccc$										
MN1 MMN3 1 0 0.00000 0.50000 0.50000 0.20484 0.96000 -2.000 -2.000 2.00										
0.00 0.00 0.00 241.00 0.00 -311.00 -291.00 321.0										
0.000 0.000 0.000 0.000 0.000 0.000 0.0000										
0.00 0.00 0.00 0.00 0.00 0.00 0.00										
MN1 MMN3 1 0 0.50000 0.50000 0.00000 0.20484 0.96000 -2.000 2.000 2.00										
0.00 0.00 0.00 241.00 0.00 -311.00 291.00 321.0										
0.000 0.000 0.000 0.000 0.000 0.000 0.000										
0.00 0.00 0.00 0.00 0.00 0.00 0.00										
The Date of Date of the Date o										

 $A_x$   $G_y$   $F_z$ 

More than 100 possibilities to put F, A, C, G in the three directions. They can appear twice  $\rightarrow$  Trial and error ??? or three times.

#### Couplings within a 4-fold site:

Site	C	Coupling								
1	+	+	+	+						
2	+	+	-	-						
3	+	-	-	+						
4	+	-	+	-						
	F	Α	G	С						



Refine Trial\_and\_error2.pcr → Trial\_and\_error2.prf

C:\FullProf\_Suite\winplotr.exe

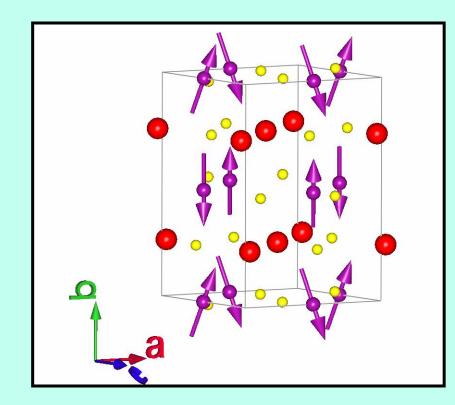
Results looks really good!

Difference plot shows all magnetic peaks to be refined

Trial\_and\_error2.new

 $R_{mag}$  is low

 $A_xG_yF_z$  seems to be the solution, let's publish!





Short reminder on the way a magnetic moment is defined:

$$m_j = \sum_{\kappa} S_{\kappa j} e^{-2\pi i \kappa t}$$

If only one propagation vector  $\kappa$ :

$$m_i = S_{\kappa i} e^{-2\pi i \kappa t}$$

$$S_{\kappa j} = \sum_{n} C_{\kappa j n} S_{\kappa n}$$

n Coefficients (variables) to be determined for n Basis vectors for spin on atom j

 $\mathbf{S}_{\kappa j}$  are the Fourier components of the magnetic atom j in the zeroth unit cell, t is the translation (e.g. t = (1 0 0) to neighbouring unit cell in x-direction)

Linear combination of n Basis vectors gives the magnetic moment value and its direction

Examples of Basis vectors:

(100) (010) (001) corresponds to the x, y and z directions of an orthogonal lattice

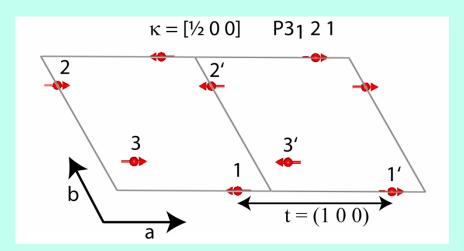
But can be as well e.g. in a hexagonal unit cell: (110) or (120)

They can be real or complex

#### Short reminder of the meaning of the term: "magnetic propagation vector"



The magnetic propagation vector describes how the magnetic structure within the nuclear unit cell is periodically repeated in space. It is equivalent to the translations used to describe the periodicity of crystal structures.



With the use of the propagation vector one can describe the relation between the magnetic moments (size and direction) on equivalent atom sites in different unit cells.

If e.g. there is only one Basis vector  $S_{\kappa}$  = (1 0 0) and  $\kappa$  =  $[\frac{1}{2}$  0 0] then the magnetic moment in the zeroth unit cell (t = (0 0 0)) is:

$$m_j = S_{\kappa j} e^{-2\pi i \kappa t} = C_{\kappa j} S_{\kappa} (\cos 2\pi \kappa t + i \sin 2\pi \kappa t) = C_{\kappa j} (100) \cos 2\pi (\frac{1}{2}00) (000) = C_{\kappa j} (100)$$

In the neighbouring unit cell in x-direction one has with  $t = (1 \ 0 \ 0)$ :

$$m_j = S_{\kappa j} e^{-2\pi i \kappa t} = C_{\kappa j} S_{\kappa} (\cos 2\pi \kappa t + i \sin 2\pi \kappa t) = C_{\kappa j} (100) \cos 2\pi (\frac{1}{2}00) (100) = C_{\kappa j} (-100)$$

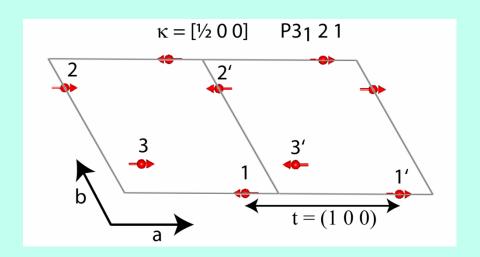
 $\rightarrow$  Antiferromagnetic alignment of magnetic moments in neighbouring unit cells along the x-direction

We know now that a magnetic moment on an atom site is described by a linear combination of Basis vectors (multiplied by their coefficients) and that the periodicity within the crystal lattice is described by the propagation vector.



If one has a crystallographic unit cell where the magnetic atom occupies a site with multiplicity 1 (e.g. Wykoff site  $1d(\frac{1}{2}0\frac{1}{2})$  in Pmmm) this is all the information we need.

If, however, the magnetic site(s) has(ve) higher multiplicities we have to determine the relation the different sites which are created through the symmetry elements of the spacegroup have to each other.



One has to determine the relation between the moment size and direction of the magnetic moment on site 1 to those on the symmetry related sites 2 and 3 within the crystallographic unit cell

We have seen that it is possible to do this using a trial and error method, however, in most of the cases this will be very tedious or nearly impossible and furthermore lead to often non-physical results. This is the reason why the method of Magnetic symmetry analysis has been developed by E.F. Bertaut in the sixties. Only, however, in the last 20 years the method has been applied routinely due to the availability of easy to use programs like BASIREPS or SARAH.



#### Magnetic symmetry analysis (a simple glimps):

The symmetry of the crystallographic structure is described by its spacegroup and the corresponding symmetry elements.

As the periodicity of the magnetic unit cell is not necessarily the same as the periodicity of the nuclear unit cell its symmetry can be lower.

The symmetry elements of the original crystallographic spacegroup which leave the magnetic propagation vector invariant form a subgroup called the « little group ».

The operators (matrixes) of these symmetry elements (permutation matrixes) multiplied by the operators which describe how the magnetic moment components are transformed as axial vectors by the symmetry elements of the little group form the magnetic representation.

This means that the magnetic representation describes the effect of the symmetry operators of the little group on the position of the atoms and on the magnetic moments.

This magnetic representation can be simplified by applying an appropriate rotation bringing the representation in the block-diagonal form where one is only left with square blocks of matrixes on the diagonal and zeros everywhere else.

Each square block forms a smaller representation and if they cannot be reduced further into a block-diagonal form; they are called irreducibel.

The different irreducible representations contain the information on the possible magnetic structures which fulfill the symmetry restrictions of the crystal structure. This limits dramatically the number of magnetic structures which have to be taken into account for the refinement.

The basis vectors of these irreducible representations can be calculated using existing programs (BASIREPS, SARAH) and will serve for the refinement of the magnetic structures.



#### Landau theorie of second order phase transitions:

The free energy of a system which sees a magnetic ordering with wave vector  $\kappa$  must be invariant under the action of the symmetry elements of the little group of  $\kappa$ .

Or simpler: If a magnetic ordering is of second order only a single irreducible representation should be involved.

This reduces strongly the possible magnetic structures: Group theory has determined a limited number of possible irreducible representations (IREP), Landau theory tells us now that only one IREP should be involved as a magnetic structure is formed.

So we are happy!! But:

- I) This is not always valid (higher order terms in the free energy)
- II) What happens if we have the Presence of several magnetic sublattices, e.g. A and B:
- 1) If the intrasite interactions  $I_{AA}$  and  $I_{BB}$  are large compared to the intersite interaction  $I_{AB}$  then one can expect two different transitions and different IREPS for each site.
- 2) If the intersite  $I_{AB}$  is dominant then both sites A and B should order at the same temperature and follow the same IREP
- 3) If only one intrasite interaction  $I_{AA}$  is large and the intersite interaction  $I_{AB}$  as well strong then the second site B gets polarized and adopts a magnetic structure similar to site A. At lower temperatures a second magnetic transition will appear for the ordering of site B. Due to the strong  $I_{AB}$  its symmetry should still follow the same IREP as A.
- III) The long range magnetic order induces a strong magnetostructural effect: first order transition.

E.g.: Strong single-ion anisotropy of the magnetic ion. Especially for rare earth ions which have a strong unquenched orbital angular moment. Interaction with the crystal field leads to easy axis or easy plane anisotropy which can determine the moment orientation in contradiction to the group theory. Does not necessarely lead to a visible structural transition.

Clemens Ritter

FullProf School 2023



## Intermediate summary:

Magnetic structures are defined by the moment size and direction of a first magnetic site, its relation to the other magnetic atoms on the symmetry related sites within the nuclear unit cell and finally by the magnetic propagation vector.

Thanks to magnetic symmetry analysis and the Landau theory one can normally strongly reduce the number of parameters to be determined for the description of a magnetic structure from  $3 \times n$ , where n is the number of magnetic atoms in the unit cell.

The solutions will be physically reasonable.

The complicated mathematics and the tedious calculations involved in the magnetic symmetry analysis are nowadays done by easy to use programs.

In the second part of the talk we will learn how to use the program BASIREPS which calculates IREPS and their Basis vectors (BV) using very limited input information.





BASIREPS: Included in the Fullprof\_Suite of programs



#### BASIREPS:

calculates through magnetic symmetry analysis the allowed irreducible representations and their basis vectors

needs as input the crystallographic spacegroup, the position of the magnetic atoms and the magnetic propagation vector

creates a file called \*.fp which contains the results of the calculations in a form which can be copied directly to the \*.pcr file for the refinement and a file called \*.bsr which contains details of the calculations.



#### 2 Examples:

1) 
$$Y_{0.1}Ca_{0.9}Mn_{0.96}Ga_{0.04}O_3$$

Mn on 0 0 
$$\frac{1}{2}$$

(Wykoff site 4b)

$$\kappa = [0\ 0\ 0]$$
  $T_N = 100\ K$ 

2) TbFe<sub>3</sub>(
$$BO_3$$
)<sub>4</sub>

(Wykoff site 3a)

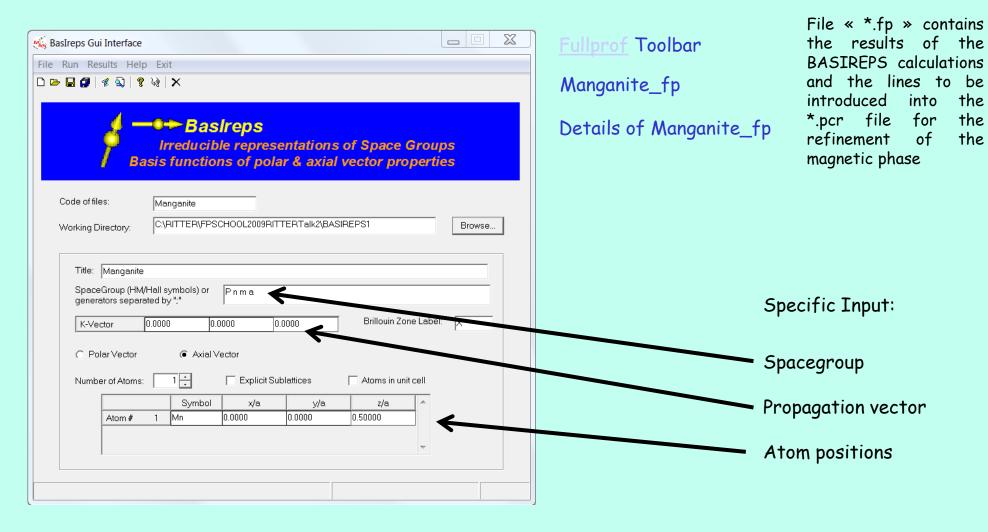
Fe2 on 0.79 0.45 0.34 (Wykoff site 6c) Tb on 0.66 0.66 0

(Wykoff site 3a)

$$\kappa = [0\ 0\ \frac{1}{2}]$$
  $T_N = 40\ K$ 



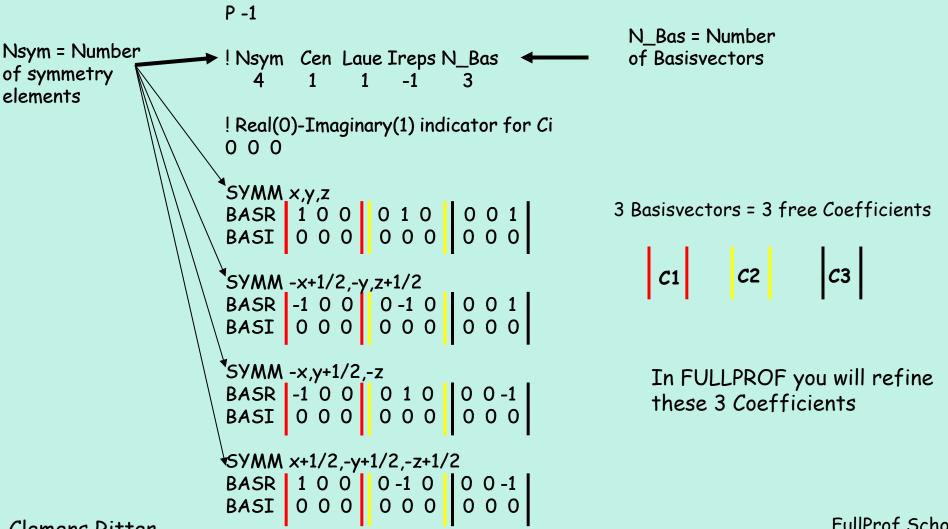
#### 1) $Y_{0.1}Ca_{0.9}Mn_{0.96}Ga_{0.04}O_3$





#### One of the possible 4 solutions:

This has to be copied to the FULLPROF pcr file in order to try a refinement



Clemens Ritter

Data for PHASE number: 2 ==> Current R\_Bragg for Pattern# 1: Model of Basireps Input pcr file Jbt=1: Phase is treated as pure magnetic phase Irf=-1: Satellite reflections generated from space group symbol Isy=-2: The basis functions of the irreducible representation are read in the form: SYMM ... BASR ... BASI ... Nvk = Number of propagation vectors Magnetic phase !Nat Dis Mom Pr1 Pr2 Pr3 Jbt Irf Isy Str Furth ATZ Nvk Npr More 0 0.0 0.0 1.0 1 -1 -2 0 0 17501.770 **1** 5 0 <--Space group symbol for hkl generation P -1 ! Nsym Cen Laue Ireps N\_Bas Laue Class: mmm = 3 -1 ! Real(0)-Imaginary(1) indicator for Ci 0 0 0 Ireps: Number of irreducible representations, often equal to the number of magnetically independent atoms, here therefore = 1 SYMM x,y,z BASR 1 0 0 0 1 0 0 0 1 BASI 0 0 0 0 0 0 0 0 SYMM - x + 1/2, -y, z + 1/2BASR -1 0 0 0 -1 0 0 0 1 BASI 0 0 0 0 0 0 0 0 SYMM -x,y+1/2,-zBASR -1 0 0 0 1 0 0 0 -1 BASI 0 0 0 0 0 0 0 0 SYMM x+1/2,-y+1/2,-z+1/2BASR 1 0 0 0 -1 0 0 0 -1 BASI 0 0 0 0 0 0 0 0

```
Typ = Magnetic formfactor MMN3 = Mn<sup>3+</sup>
         Mag = Irreducible representation used for this Atom
                   Vek = Magnetic propagation vector used for this Atom
!Atom Typ Mag Vek X
                                          Biso
                                                  Occ
                                                           C1
                                                                  C2
                                                                        C3
                           C4
                                  C5
                                         C6
                                               C7
                                                     C8
                                                             C9
                                                                   MagPh
               1 0.00000 0.00000 0.50000 0.22051 0.96000
MN1 MMN3 1
                                                         -0.373
                                                                 0.290 -4.112
                           0.00
                                  0.00
                                         241.00
                                                   0.00
                                                          11.00
                                                                 21.00 31.00
                   0.00
                                0.000 0.000 0.000 0.000 0.000 0.0000
                          0.000
                           0.00
                                  0.00
                                        0.00
                                               0.00
                                                    0.00
                                                            0.00
                                                                     0.00
!-----> Profile Parameters for Pattern # 1
          Shape1
                    Bov
                          Str1
                                  Str2
                                         Str3 Strain-Model
! Scale
           0.28936 0.00000 0.00000 0.00000
 2.8989
                                               0.00000
 311.00000 121.000
                     0.000
                             0.000
                                     0.000
                                             0.000
                                У
                                      GauSiz LorSiz Size-Model
                 W
                         X
 0.122186 -0.184594 0.205900 0.000000 0.000000 0.000000 0
         141.000
                              0.000
                                      0.000
  131,000
                  151.000
                                               0.000
                                                       0.000
                                 alpha
                                             beta
                                                                #Cell Info
               b
                                                       gamma
     а
                         C
 5.300002 7.451008 5.261796 90.000000 90.000000 90.000000
 91.00000 101.00000 111.00000
                                 0.00000
                                            0.00000
                                                      0.00000
          Pref2
                  Asy1
! Pref1
                          Asy2
                                   Asy3
                                           Asy4
 0.00000 0.00000 0.03350 0.01794 0.00000 0.00000
  0.00
           0.00
               271.00
                         281.00
                                    0.00
                                            0.00
! Propagation vectors:
 0.0000000 0.0000000
                        0.0000000
                                       Propagation Vector 1
  0.000000
             0.000000
                        0.000000
```

Trying out the 4 different solutions proposed by BASIREPS

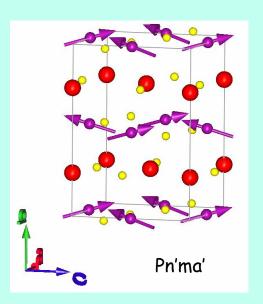


Copy the corresponding lines of the manganite.fr file created by BASIREPS into the \*.pcr file

Refine and decide from the R-factors and from the visual inspection of the fit whether the solution is correct

Best Fit with couplings:

$$\frac{x \ y \ z}{+ \ + \ +}$$
 $+ \ + \ - \ + \ +$ 
 $- \ + \ +$ 



Manganite\_pcr
Model1

WinplotR

As this model does not work return to the fp file and copy a different IREP into the pcr file

Manganite\_fp

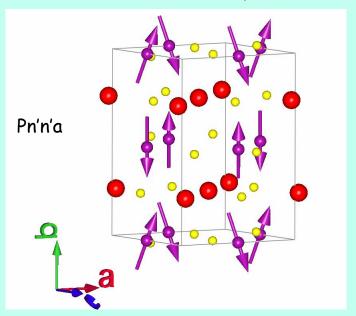
Manganite\_pcr Model2

Manganite\_pcr Model3

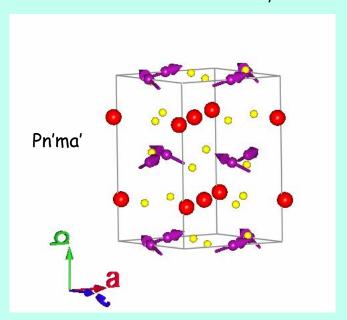
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Compare with the intuitively (trial and error) found solution:  $A_xG_yF_z$ Pseudosymmetry and not sufficient resolution are causing this situation Only BASIREPS gives you the correct spin direction,

Trial and error:  $A_xG_yF_z$ 



Symmetry analysis:  $A_x F_y G_z$ 

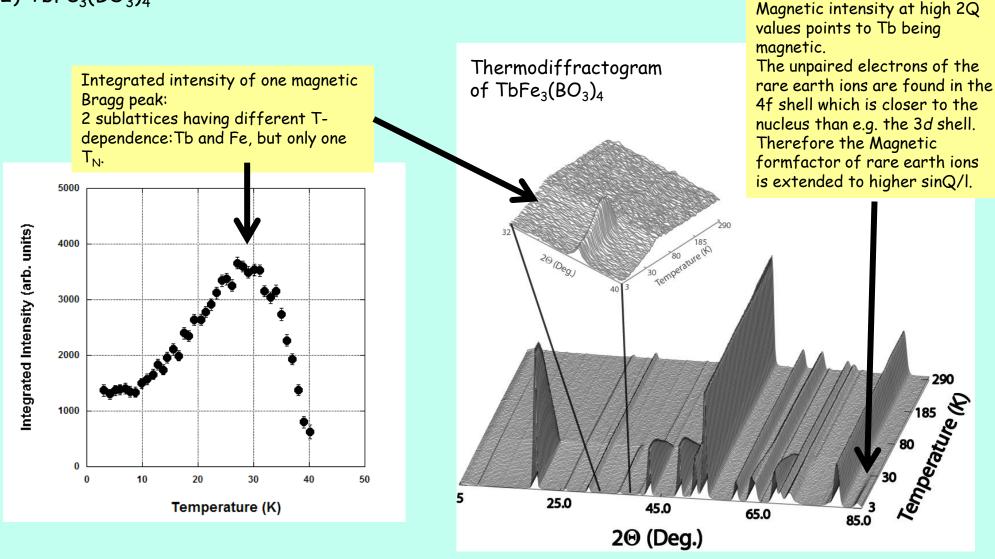


Pseudosymmetry and not sufficient resolution are causing this situation:

b = 
$$7.4507$$
 c =  $5.2616$ 



2) TbFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>



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 $\mathsf{TbFe}_3(\mathsf{BO}_3)_4$ 

1) Fit at 50 K > T<sub>N</sub>

Nuclear Structure pcr

Nuclear Structure prf

2) Compare data at 50 K and 2 K and fit the peak positions of the magnetic peaks

WinplotR

3) Create a \*.sat file for K-search, apply the zeroshift of the 50 K refinement to the determined peak positions

TbFe3(BO3)4\_sat

4) Run K-search from the green version of <u>WINPLOTR</u>

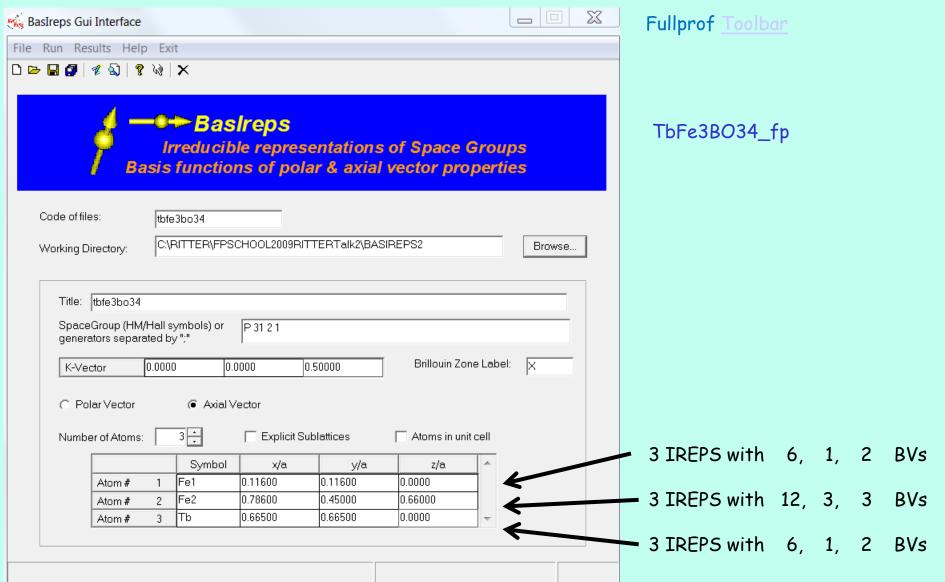
5) Look at the result file TbFe3(BO3)4\_kub

(Notice that there are a lot of solutions of type " $(\frac{1}{8}, \frac{1}{4}, \frac{1}{4})$ " or " $(\frac{1}{8}, \frac{1}{4}, 0)$ " or " $(\frac{1}{4}, \frac{1}{4}, \frac{1}{4})$ ",  $\rightarrow$  Ignore these in 99% of cases!!)

$$\rightarrow \kappa = [0\ 0\ \frac{1}{2}]$$

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#### $\mathsf{TbFe}_3(\mathsf{BO}_3)_4$



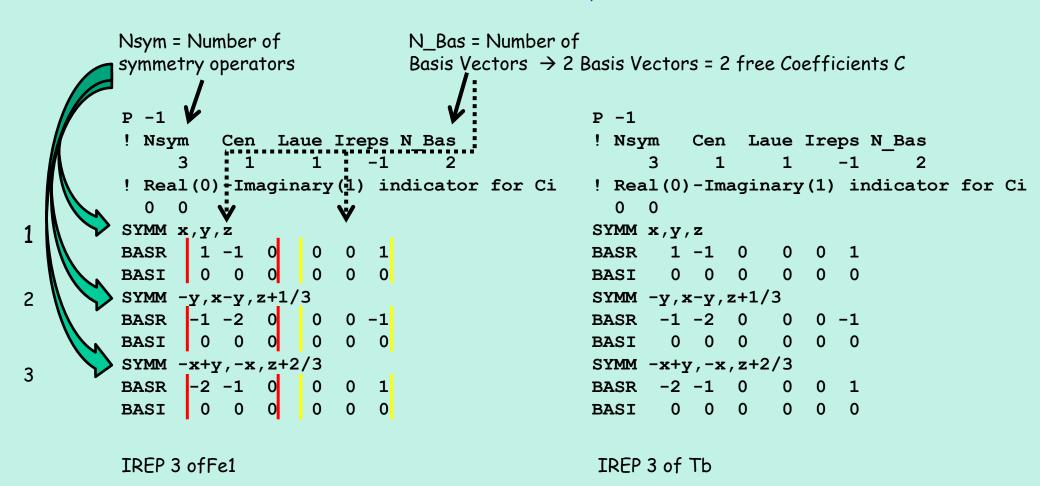
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 $\mathsf{TbFe}_3(\mathsf{BO}_3)_4$ 

IREP3: TbFe3BO34\_fp



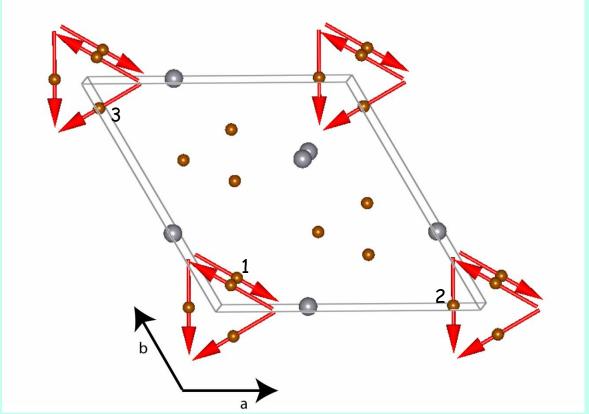
Both IREPS are identical! This is normal as both atoms are sitting on the same Wykoff site 3a (x x 0)

 $\mathsf{TbFe}_3(\mathsf{BO}_3)_4$ 



```
P -1
! Nsym
        Cen Laue Ireps N Bas
! Real(0)-Imaginary(1) indicator for Ci
    0
  0
SYMM x,y,z
BASR
      1 -1 0
BASI
      0 0
SYMM -y, x-y, z+1/3
BASR -1 -2 0
BASI
                   0 0
SYMM -x+y, -x, z+2/3
BASR -2 -1 0
                  0 1
                0 0
      0 0 0
BASI
```

Magnetic structure with only BV1



IREP3 of Fe1



```
\mathsf{TbFe}_3(\mathsf{BO}_3)_4
```

IREP 3 of Fe2 has 6 symmetry operators and 3 BVs  $\rightarrow$  3  $C_i$ 

```
P -1
! Nsym
        Cen Laue Ireps N Bas
                    -1
          1
                          3
! Real(0)-Imaginary(1) indicator for Ci
 0 0 0
SYMM x,y,z
      1 0 0
BASR
             0 0 0
      0 0 0
BASI
SYMM -y, x-y, z+1/3
      0 -1 0
               1 1 0
BASR
      0 0 0
               0 0 0
                        0 0 0
BASI
SYMM -x+y, -x, z+2/3
BASR -1 -1 0
               1 0 0
      0 0 0 0 0
BASI
SYMM y,x,-z
      0 -1 0 -1 0
                   0
BASR
      0 0 0
               0 0 0
                        0 0 0
BASI
SYMM x-y, -y, -z+2/3
BASR -1 0 0
               1 1 0
               0 0 0
                        0 0 0
BASI
      0 0 0
SYMM -x, -x+y, -z+1/3
BASR -1 -1 0
               0 1 0
                        0 0 -1
               0 0 0
      0 0 0
                        0 0 0
BASI
```

```
P -1
! Nsym
        Cen Laue Ireps N Bas
          1
     3
                1
                     -1
                            2
! Real(0)-Imaginary(1) indicator for Ci
  0 0
SYMM x,y,z
BASR
      1 -1 0
      0 0 0 0 0 0
BASI
SYMM -y, x-y, z+1/3
BASR -1 -2 0
                0 0 -1
BASI
      0 0 0
                0 0 0
SYMM -x+y, -x, z+2/3
BASR -2 -1 0
                0 0 1
BASI
      0 0 0 0 0 0
```

The BVs of the Wykoff site 6c (Fe2) are different than those of site 3a (Fe1, Tb)

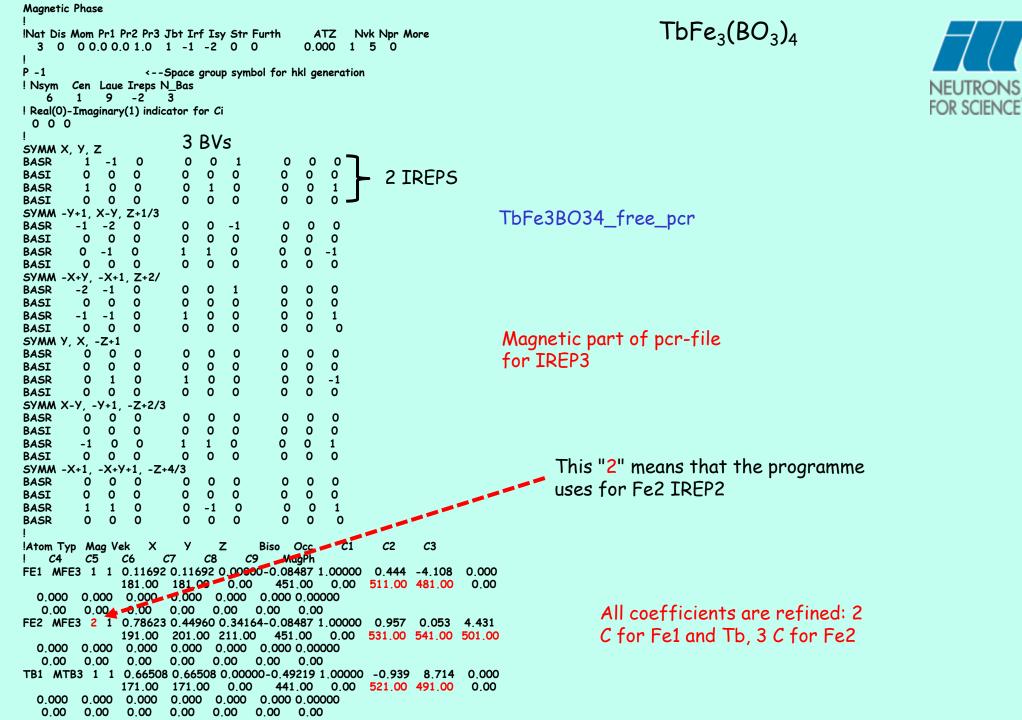
#### $\mathsf{TbFe}_3(\mathsf{BO}_3)_4$

BASI

#### Construction of the pcr file for IREP3

```
I)
                                                                                               P -1
     ! Nsym
     ! Real(0)-Imaginary(1) indicator for Ci
                                                  ! Real(0)-Imaginary(1) indicator for Ci
                                                                                               ! Real(0)-Imaginary(1) indicator for Ci
      0 0
                                                    0 0
                                                                                                 0
                                                                                                   0
    SYMM x,y,z
                                                  SYMM x,y,z
                                                                                               SYMM x,y,z
    BASR
                0
                    0
                                                  BASR
                                                          -1
                                                                                               BASR
                                                                                                                                 Add zeros to
    BASI
                                                  BASI
                                                  BASR
                                                                                                                                      get the same
                                                  BASI
                                                                                               BASI
                               Fe1
                                                                                                                                      number of BVs
                                                  BASR
                                                  BASI
          -2 -1
    BASR
                                                                                               BASI
    BASI
                                                  BASR
                                                                                               BASR
     P -1
                                                  BASI
             Cen Laue Ireps N Bas
                                                  BASR
                                                                                               SYMM -x+y, -x, z+2/3
                                                  BASI
                                                                                               BASR
     ! Real(0)-Imaginary(1) indicator for Ci
                                                  BASR
                                                  BASI
     SYMM x,y,z
                                                                                               BASI
                                                                                                                                    __ Add zeros under
    BASR
                                                  BASR
                                                                                                                                      the symmetry
                                                                                               BASR
                                                                                               BASI
                                                                                                                                      operators which
                                                  BASI
    BASI
                                                                                                                                      do not concern
                                                  BASR
                                                                                               BASI
                                                  BASI
                                                                                                                                      Fe1 and Tb
                                                  SYMM v,x,-z
    BASI
                                                  BASR
                                                                                               BASI
                                                  BASI
                                                                                               BASR
                                                                                               SYMM
                                                  BASI
                                                                                               BASR
    BASI
    BASI
    P -1
                                                     Group together
                                                                                                 Remove the identical
     ! Real(0)-Imaginary(1) indicator for Ci
                                                     under the same
      0 0
                                                                                                 lines as IREP3 of Fe1
    SYMM x,y,z
                                                     symmetry
                                                                                                 and Tb are the same
     BASI
```

operators





# $TbFe_3(BO_3)_4$ Refinement

Winplote

1) Refinement with all 7	2) Refinement with only 2
coefficients free	coefficients

TbFe3BO34\_free\_pcr TbFe3BO34\_fixed\_pcr

TbFe3BO34\_free\_sum TbFe3BO34\_fixed.sum

Fit not stable, large Fit stable, small error bars if error bars, peak at low angles? Fit stable, small error bars if C2(Fe1) = -C2(Fe2), low angle peak not longer there.

ly 2 3) Refinement with single crystal data

TbFe3BO34\_SingleCrystal\_free\_pcr

TbFe3BO34\_SingleCrystal\_free\_prf

TbFe3BO34\_PowderRefinement\_using SingleCrystal\_results\_prf

C1=0.39(5) C2=3.95(5) C3=0.00

Tb1 C1=-0.52(11) C2=8.66(7) C3=0.00 C1=0.00 C2=-8.56(6) C3=0.00 C1=0.00 C1=0.0

C1=0.00 C2=0.00 C3=-4.34(4)

C1=0.16(4) C2=-8.09(3) C3=0.00

C1 = -0.09(6) C2 = -0.07(6) C3 = -3.92(4)

Clemens Ritter

Fe1 C1=-0.21(17) C2=-3.46(72) C3=0.00

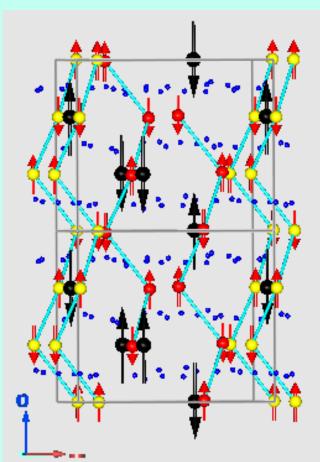
Fe2 C1=-0.45(26) C2=-0.04(11) C3=4.81(37)

ter FullProf School 2023

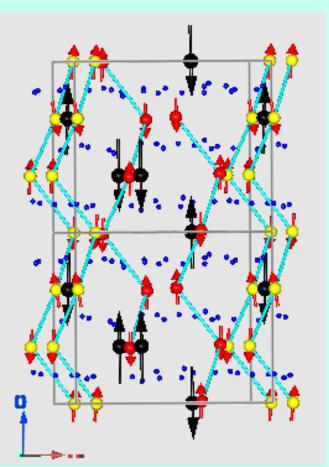
TbFe $_3(BO_3)_4$  Results of magnetic refinements



Powder, 2 Coefficients

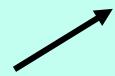


Single Crystal, 7 Coefficients



Bilbao Crystallographic Server: MAXMAGN

nttps://www.cryst.ehu.es/



Magnetic spacegroup: 2-  $P_c 3_2 21$  (#154.44)

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TbFe $_3(BO_3)_4$  Test of IREP2

<u>WinplotR</u>

TbFe3BO34_IREP2_pcr	TbFe3BO34_IREP2_sum	

TbFe3BO34\_IREP2\_prf

SYMM >	X, Y, Z									SYMM >	X, Y, Z								
BASR	1	1	0	0	0	0	0	0	0	BASR	1	-1	0	0	0	1	0	0	0
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
BASR	1	0	0	0	1	0	0	0	1	BASR	1	0	0	0	1	0	0	0	1
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
SYMM -Y+1, X-Y, Z+1/3						SYMM -	-Y+1, X	-Y, Z+1	1/3										
BASR	1	0	0	0	0	0	0	0	0	BASR	-1	-2	0	0	0	-1	0	0	0
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
BASR	0	-1	0	1	1	0	0	0	-1	BASR	0	-1	0	1	1	0	0	0	-1
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
	-X+Y, -	-X+1, Z	+2/3								-X+Y, -	X+1, Z-	-2/3						
BASR	Ó	-1	0	0	0	0	0	0	0	BASR	- <b>2</b>	-1	0	0	0	1	0	0	0
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
BASR	-1	-1	0	1	0	0	0	0	1	BASR	-1	-1	0	1	0	0	0	0	1
BASI	0	0	0	0	0	Ō	Ó	Ō	0	BASI	0	0	0	0	0	0	0	0	0
SYMM Y	y, x, -:	Z+1								SYMM Y	y, X, -Z	Z+1							
BASR	Ó	0	0	0	0	0	0	0	0	BASR	Ó	0	0	0	0	0	0	0	0
BASI	0	0	0	0	0	Ō	0	0	0	BASI	0	0	0	0	0	0	0	0	0
BASR	0	1	0	1	0	0	0	0	-1	BASR	0	-1	0	-1	0	0	0	0	1
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
SYMM >	X-Y>	/+1Z-	+2/3							SYMM >	Х-УУ	′+1Z+	-2/3						
BASR	Ó	o -	0	0	0	0	0	0	0	BASR	0	o _	0	0	0	0	0	0	0
BASI	0	0	0	0	0	0	0	0	0	BASI	0	0	0	0	0	0	0	0	0
BASR	1	Ō	Ö	-1	-1	Ō	Ö	Ö	-1	BASR	-1	0	Ō	1	1	Ō	Ō	0	1
BASI	Ō	Ö	Ō	0	Ō	Ö	Ö	Ö	Ō	BASI	0	0	0	0	0	0	0	0	0
SYMM -	-X+1 -	X+Y+1	-Z+4/3	}	•	_			_	SYMM -	-X+1 -	X+Y+1,	-Z+4/3	1	_	•	-	_	_
BASR	O O	0,	0	0	0	0	0	0	0	BASR	O O	0,	0	0	0	0	0	0	0
BASI	Ö	Ö	Ö	Ö	Ö	Ö	Ö	Ö	Ö	BASI	Ö	Ö	Ö	Ö	Ó	Ö	Ö	Ö	Ó
BASR	1	1	Ö	Ö	-1	Ö	o	Ō	1	BASR	-1	-1	Ö	Ö	1	Ö	Ö	Ö	-1
BASR	ō	ō	Ö	Ö	ō	Õ	Õ	Õ	ō	BASR	0	ō	Ö	Ö	ō	Ö	Ö	Ö	ō
			_	_		_				55.		_		_					_

IREP 2, e.g. no z-component for Fe1 and Tb

IREP 3

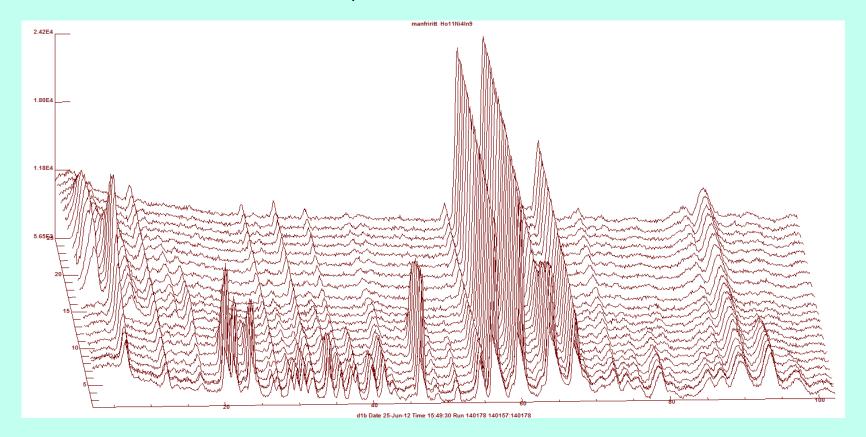
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#### More difficult case:

Ho<sub>11</sub>Ni<sub>4</sub>In<sub>9</sub>: several k-vectors

several magnetic sites

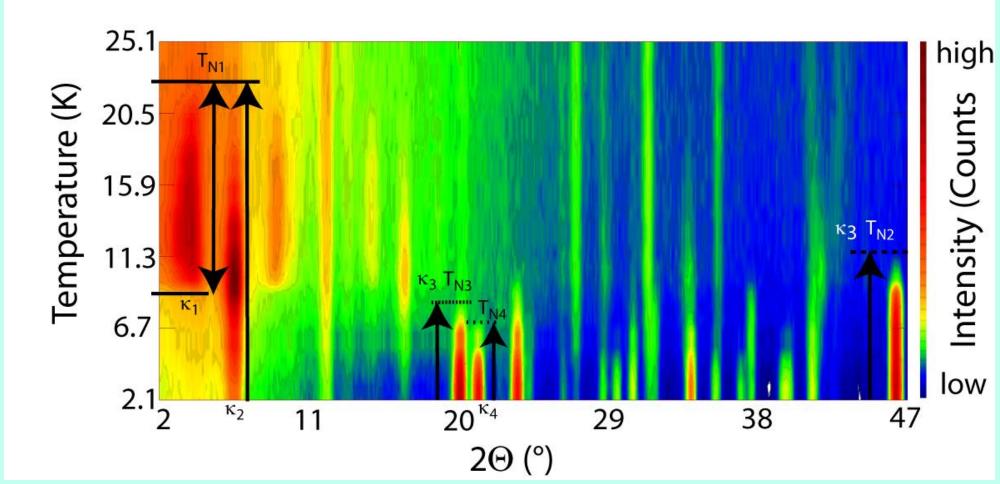
different T-dependencies





Ho<sub>11</sub>Ni<sub>4</sub>In<sub>9</sub>:

Cmmm, 5 Ho in crystallographic unit cell Ho1 on 8p, Ho2 and Ho3 on 4i Ho4 on 4g and Ho5 on 2a



 $\kappa_1 = [0\ 0.62\ 0], \ \kappa_2 = [0\ 1\ 0], \ \kappa_3 = [0\ 0\ \frac{1}{2}], \ \kappa_4 = [0\ 1\ \frac{1}{2}]$ 

At 16 K:  $\kappa_1$  only Ho2

 $\kappa_2$  only Ho5  $\kappa_4$  only Ho3 and only below  $T_{N4}$ 

16K\_pcr 16K\_prf

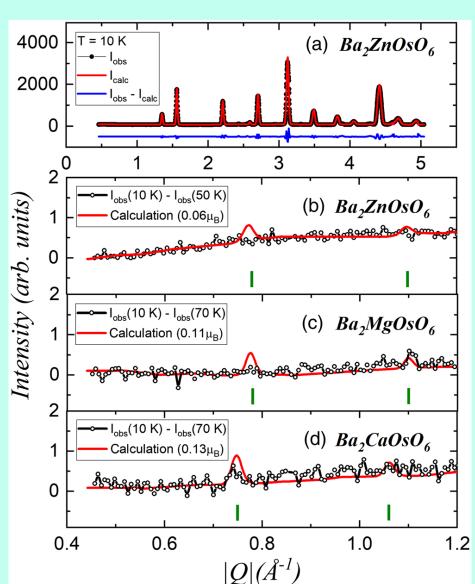


#### Simulation:

If you determine the symmetry allowed magnetic structures for a given magnetic propagation vector you can simulate the magnetic diffraction pattern and determine the upper limit of the magnetic moment.

Assuming type 1 antiferromagnetic order in this cubic system:

Absence of dipolar magnetic order in cubic  $5d^2$  Double Perovskites  $\rightarrow$  Strong indication for Octupolar order



10 h (10 K) - 10 h (50 K):  $\mu_{Os} < 0.06 \ \mu_{B}$ 



#### A short sumamary:

The different irreducible representations contain the information on the possible magnetic structures, which fulfill the symmetry restrictions of the crystal structure.

The basis vectors of these irreducible representations can be calculated using existing programs (BASIREPS, SARAH) and serve for the refinement of the magnetic structures.

Magnetic symmetry analysis limits dramatically the number of magnetic structures, which have to be taken into account for the refinement.

Magnetic symmetry analysis helps to avoid "wrong" solutions.

Magnetic symmetry analysis indicates the possible existence of a structural phase transition.

Magnetic symmetry analysis allows to simulate magnetic diffraction patterns to determine lower limits for magnetic moments.

Apart from BASIREPS there is as well the Program SARAh which does the same thing

SARAh: http://fermat.chem.ucl.ac.uk/spaces/willsgroup/web-software/sarah-refine-fullprof/