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# Angular Momentum in Fission through Isomeric Ratio measurements with the Lohengrin spectrometer

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<sup>1</sup> CEA, DES, IRESNE, DER, SPRC, Cadarache, Physics Studies Laboratory, 13108 Saint-Paul-lès-Durance, France

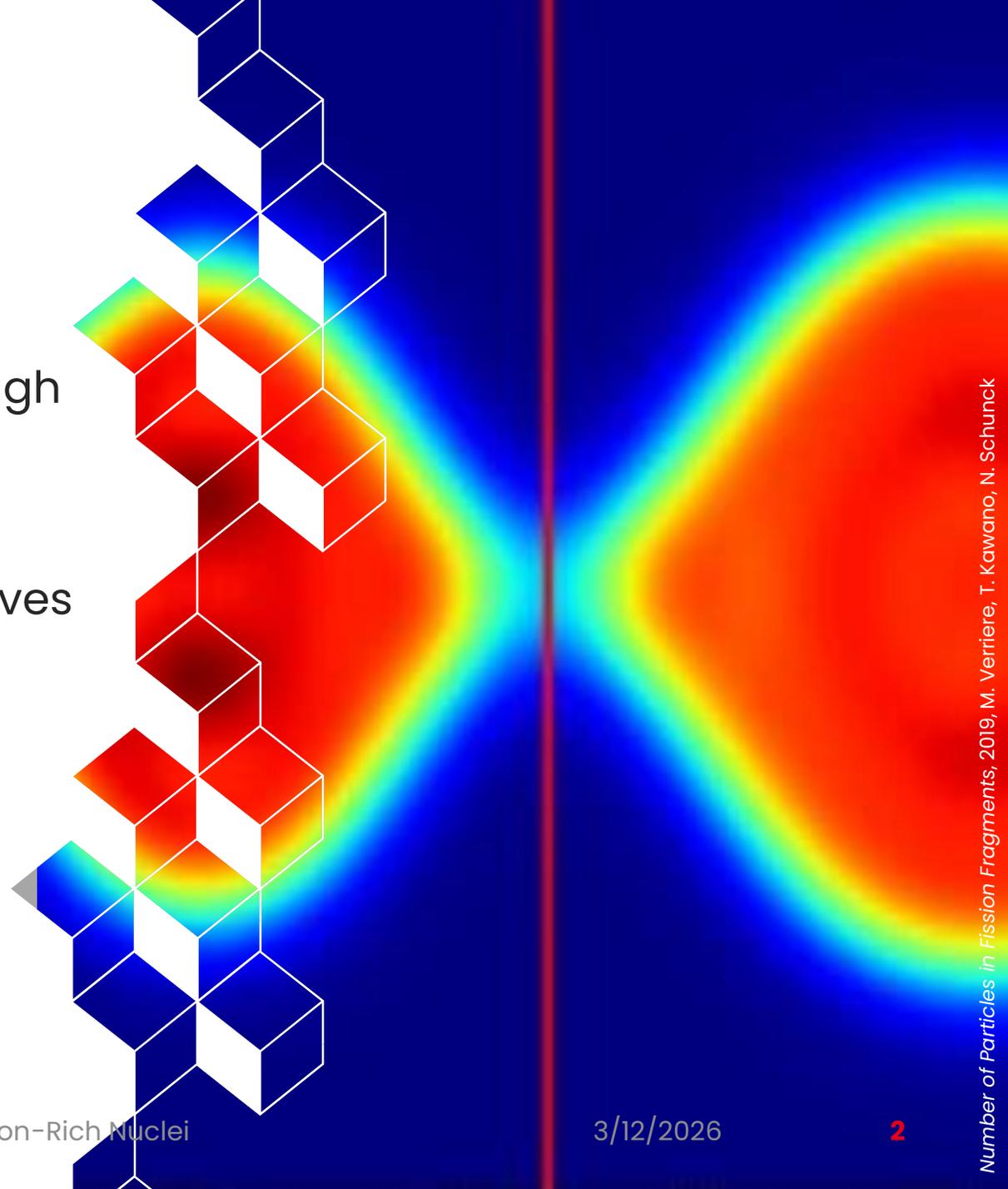
<sup>2</sup> National Technical University of Athens, Iroon polytechniou 9, Zografou Campus 15780, Athens, Greece

<sup>3</sup> Institut Laue-Langevin, 38042 Grenoble Cedex 9, France

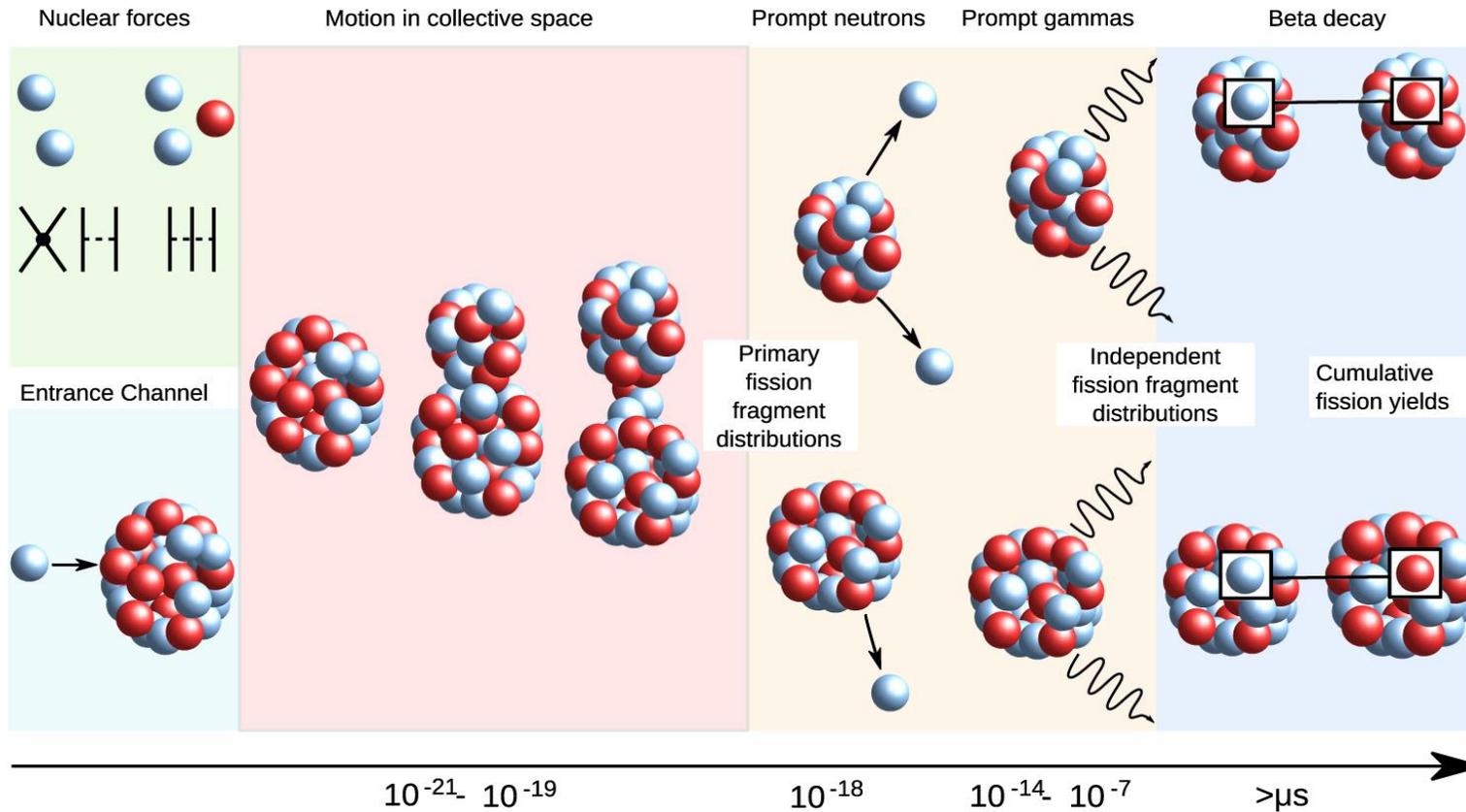
<sup>4</sup> LPSC, Université Grenoble-Alpes, CNRS/IN2P3, 38026 Grenoble Cedex, France

# Contents

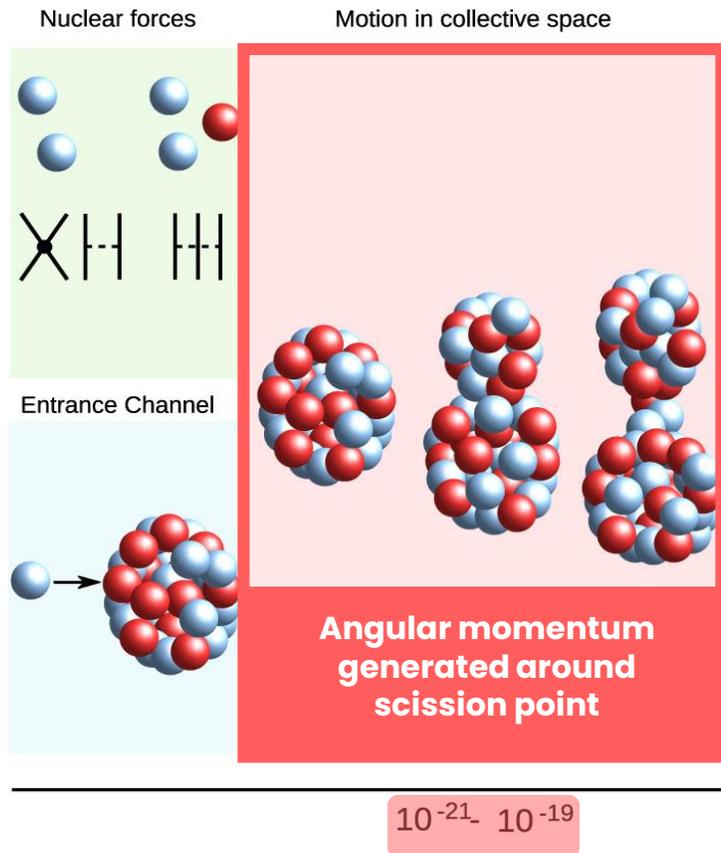
1. Motivation
2. **Isomeric Ratios of Fission Products** as observables of Angular Momentum through  $^{241}\text{Am}(2n_{\text{th}},f)$
3. Analysis Process
4. Preliminary results of  $^{100}\text{Nb}$  and perspectives



# Motivation: Angular momentum generation mechanism of Fission Fragments



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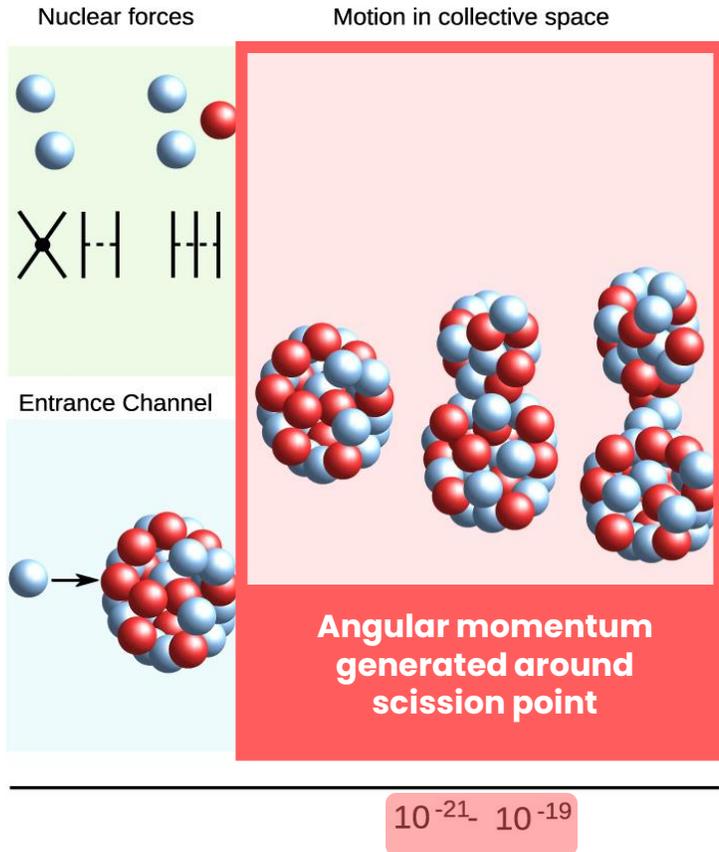
## Applications:



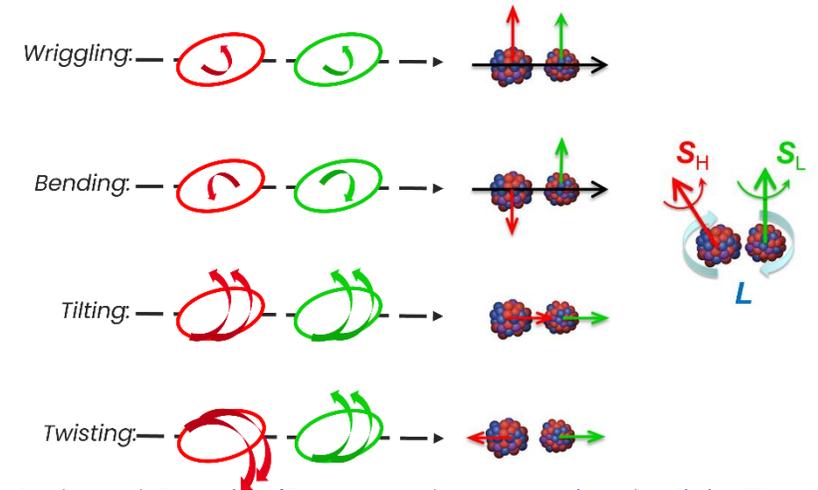
## Open questions:

- ❑ the general properties of the Angular Momenta (AM)
- ❑ the mechanism(s) behind their generation and correlation of the angular momentum of the two fission fragments

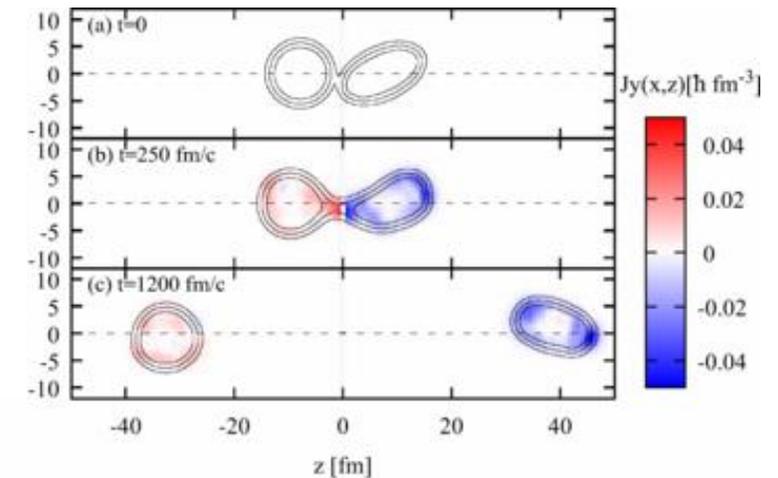
# Motivation: Angular momentum generation mechanism of Fission Fragments



## Competing theories

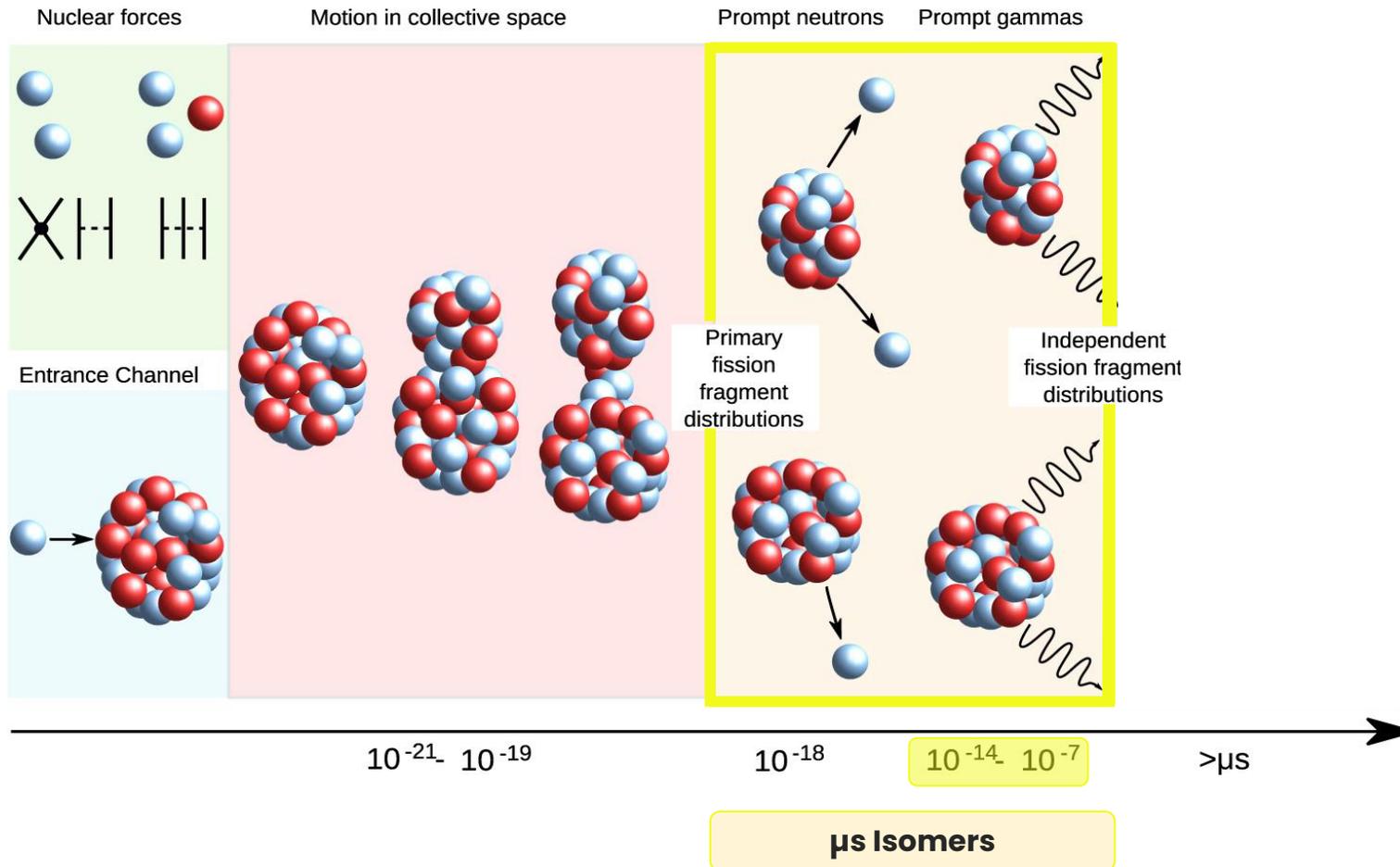


J. Randrup et al., Generation of Fragment Angular Momentum in Nuclear Fission, EPJ WoC, 284,04004 (2023)

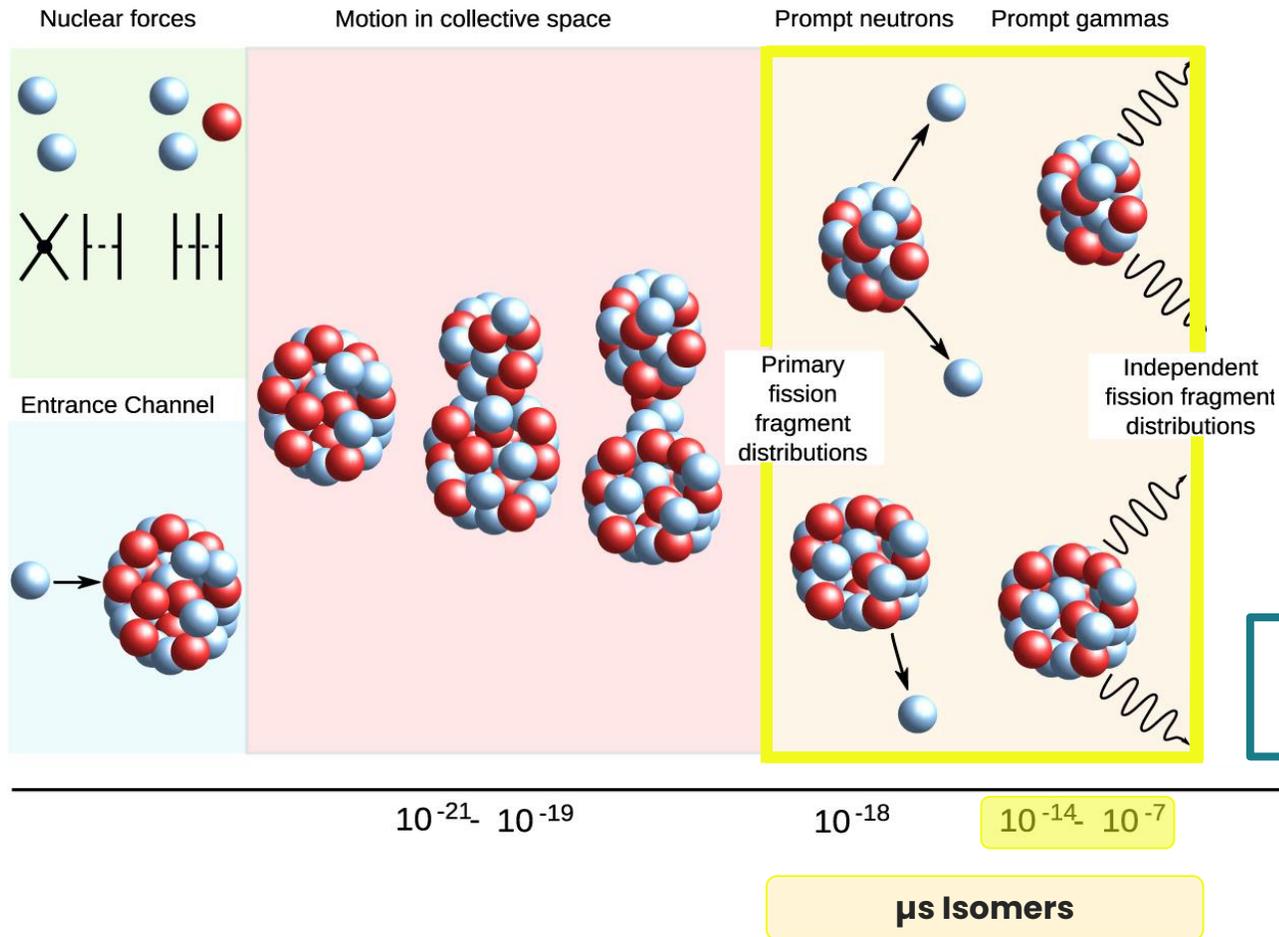


G. Scamps, Microscopic description of the torque acting on fission fragments, Phys. Rev. C 106, 054614 (2022)

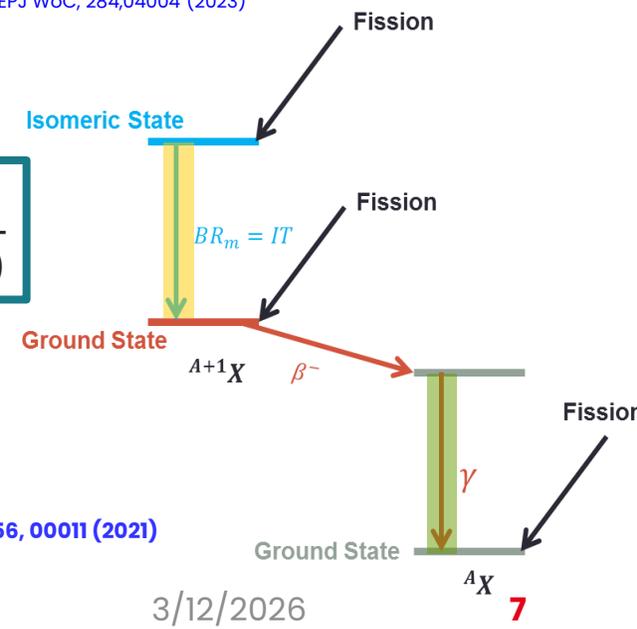
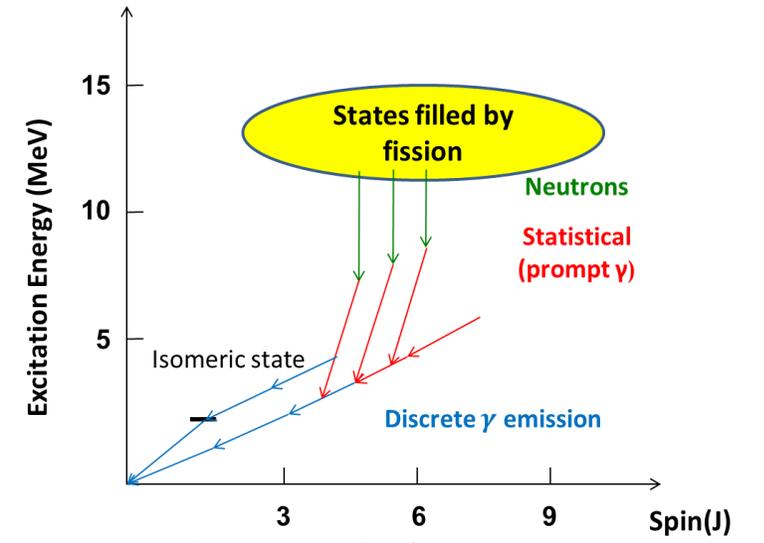
# Motivation: Angular momentum generation mechanism of Fission Fragments



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$$IR_{\text{exp}} = \frac{\tau_f(\text{IS})}{\tau_f(\text{IS}) + \tau_f(\text{GS})}$$



Investigation of fission product isomeric ratios and angular momenta of  $^{132}\text{Sn}$  populated in the  $^{241}\text{Pu}(\text{nth},\text{f})$  reaction, A. Chebboubi, EPJ Web of Conferences 256, 00011 (2021)

M. Bender et al., Future of nuclear fission theory, J. Phys. G: Nucl. Part. Phys. 47 113002 (2020)

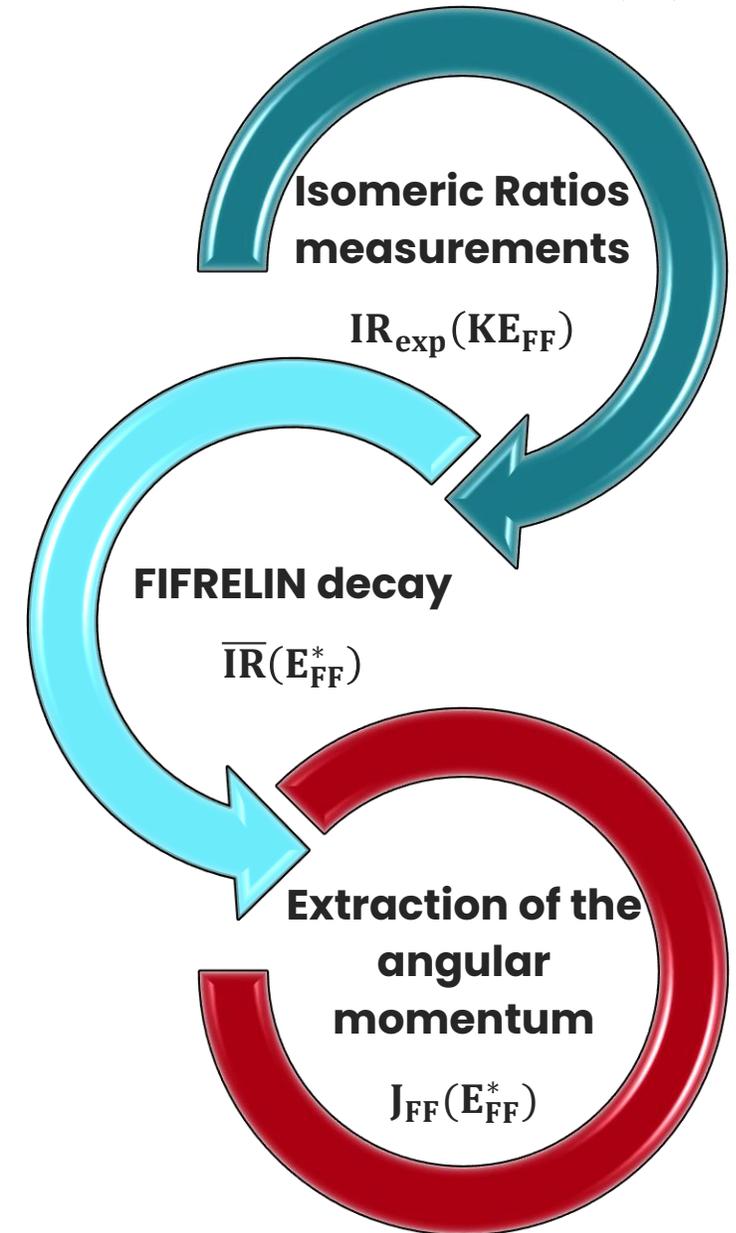
FISSION 2026: 7th Workshop on Nuclear Fission and Spectroscopy of Neutron-Rich Nuclei

3/12/2026

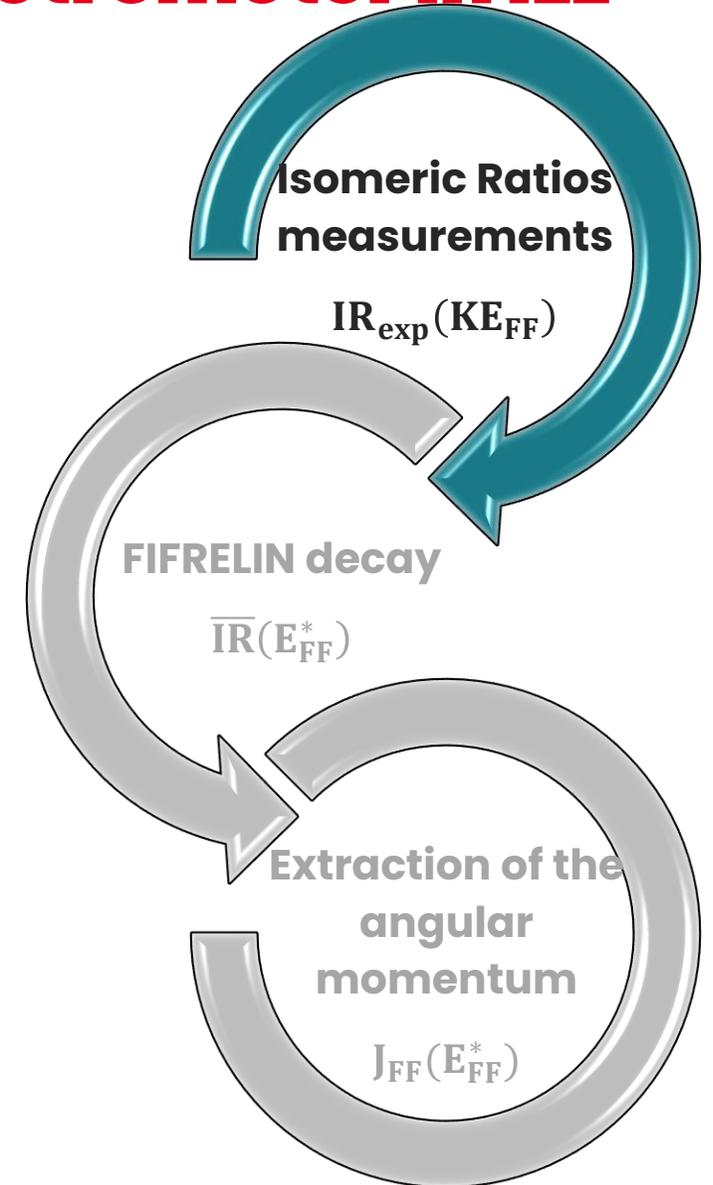
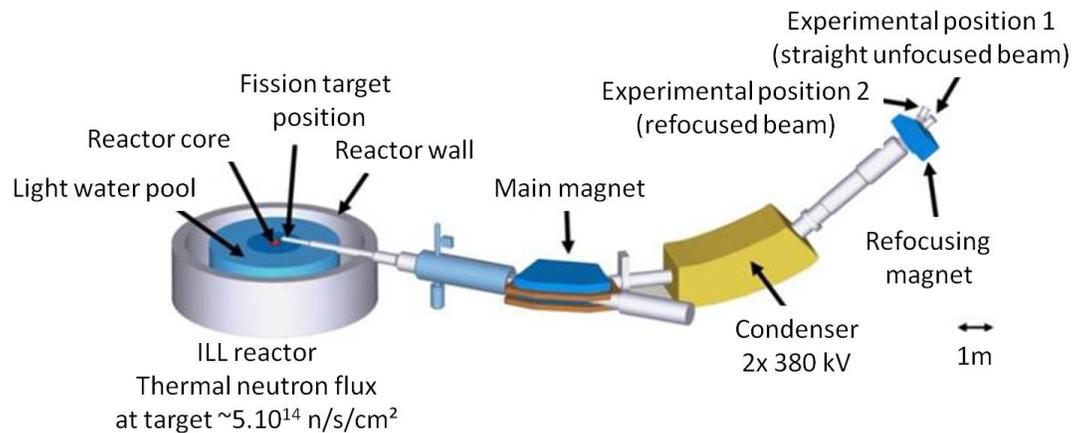
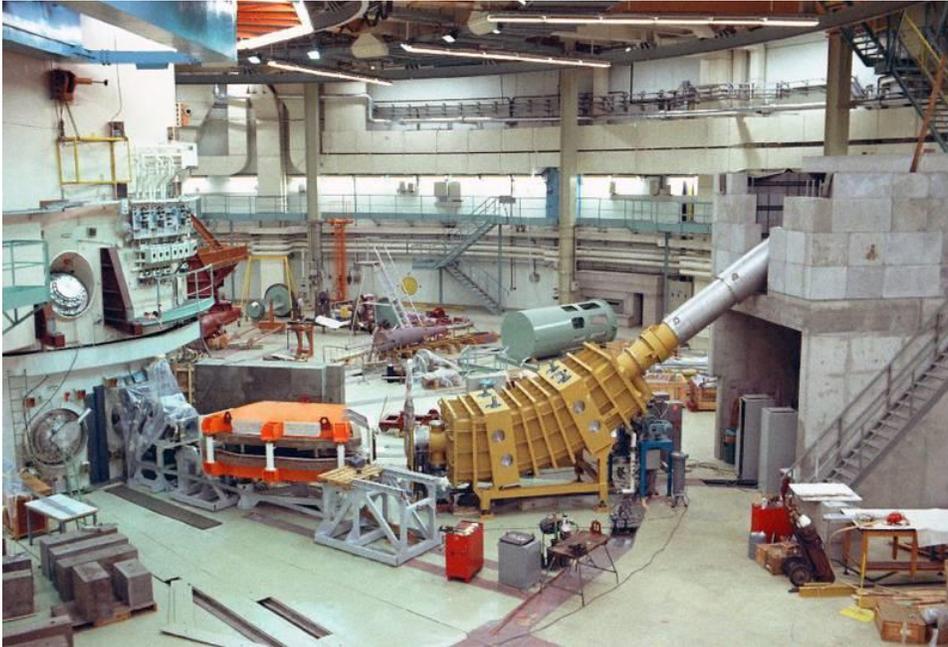
# Goals

Through experimental campaigns and interpretation, coupled with theoretical models, we aim to:

- constrain the different models of generation of the angular momentum of fission fragments
- find a possible correlation with excitation energy, through Isomeric Ratio (IR) measurements.



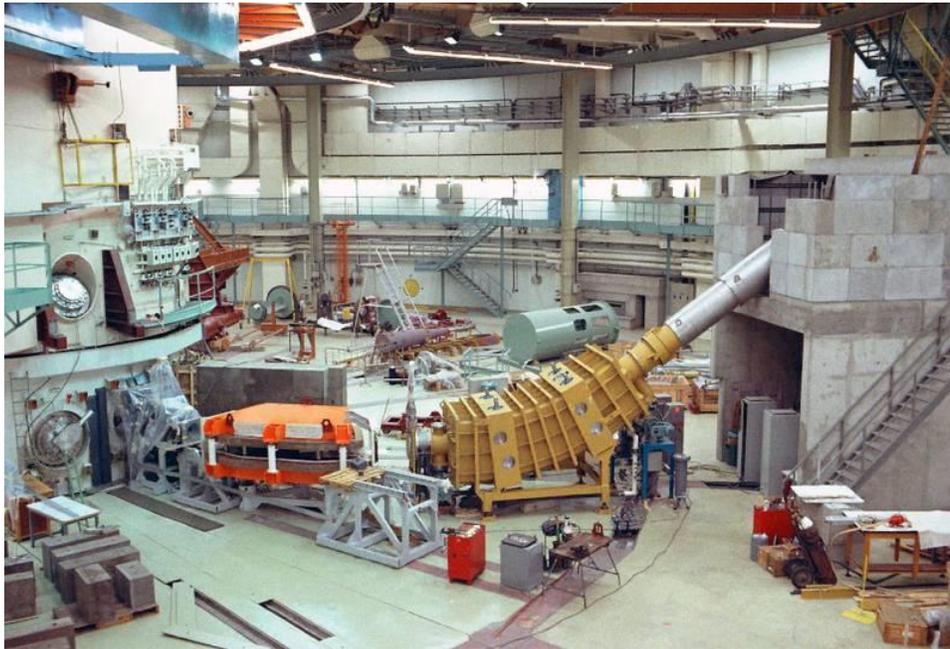
# Experimental setup of LOHENGRIN spectrometer in ILL



J. M. Daugas, Session 9, 3/11/2026, FISSION 2026 workshop

FISSION 2026: 7th Workshop on Nuclear Fission and Spectroscopy of Neutron-Rich Nuclei

# Experimental setup of LOHENGRIN spectrometer in ILL



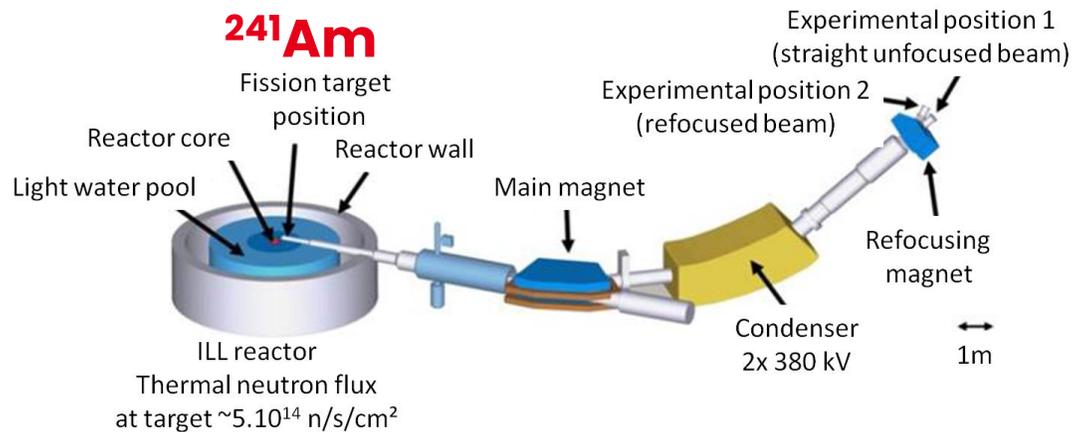
Choice of  $\vec{B}$   
and  $\vec{E}$



Selection of  
 $A/q$  and  $E_K/q$



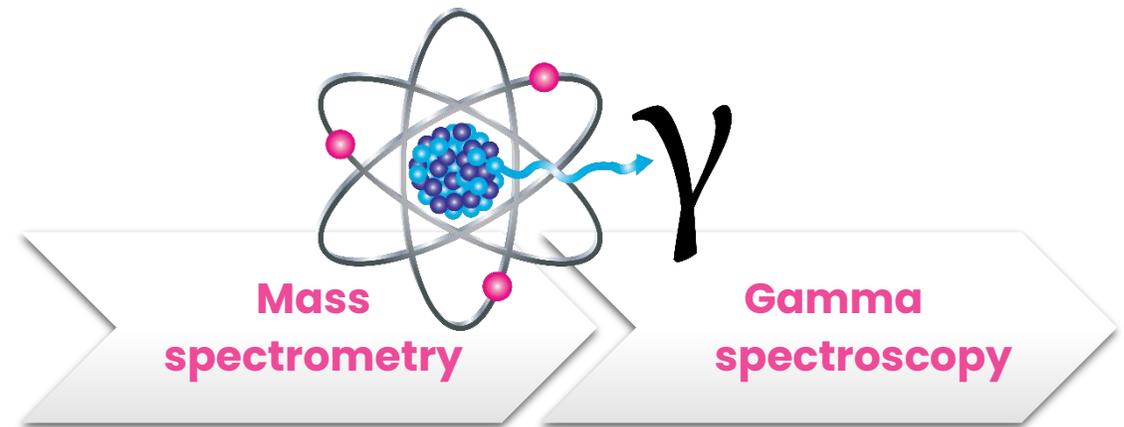
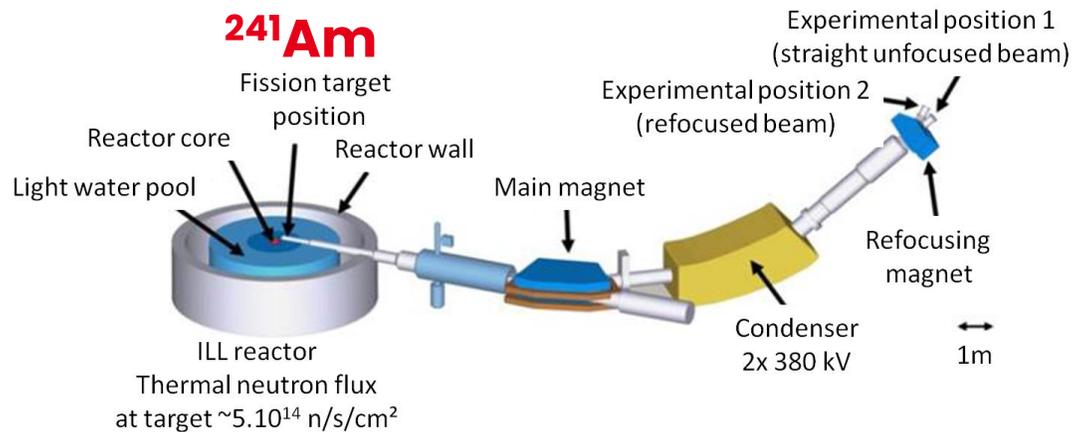
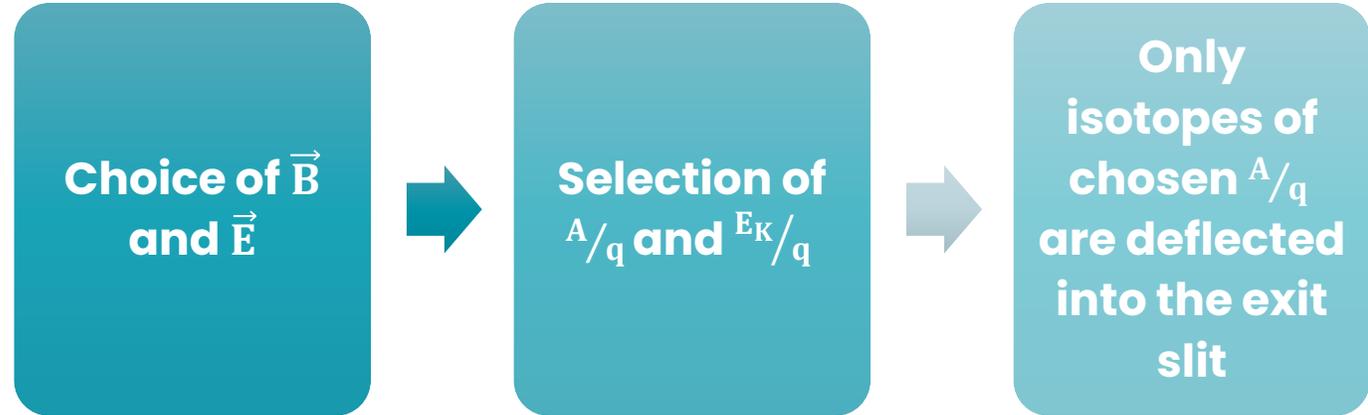
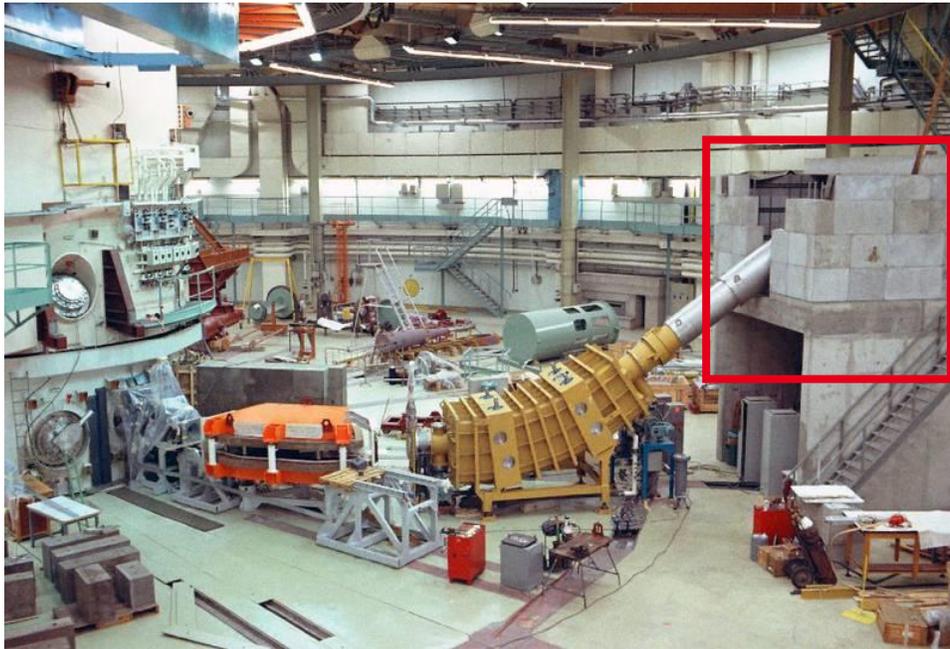
Only  
isotopes of  
chosen  $A/q$   
are deflected  
into the exit  
slit



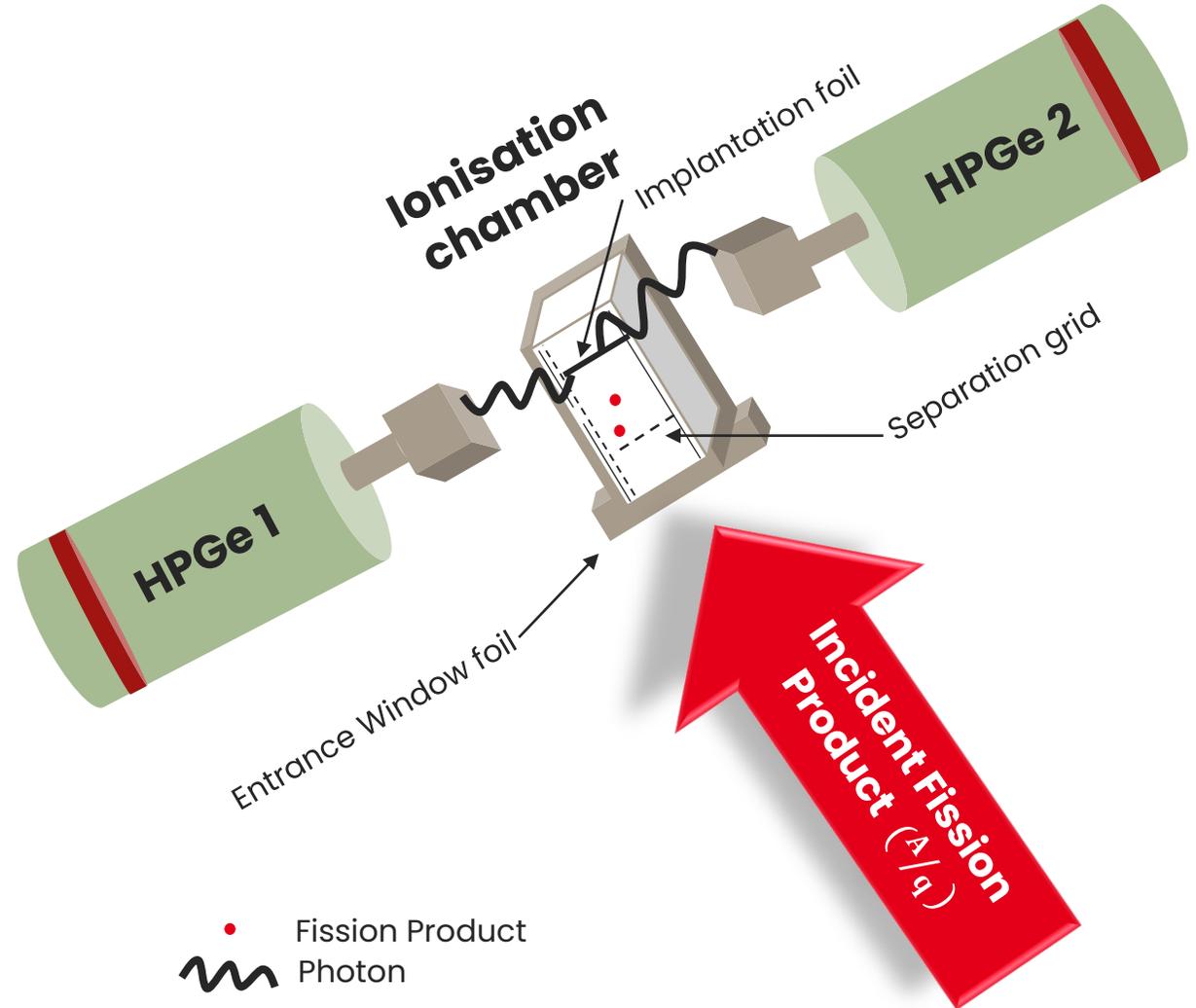
J. M. Daugas, Session 9, 3/11/2026, FISSION 2026 workshop

FISSION 2026: 7th Workshop on Nuclear Fission and Spectroscopy of Neutron-Rich Nuclei

# Experimental setup of LOHENGRIN spectrometer in ILL



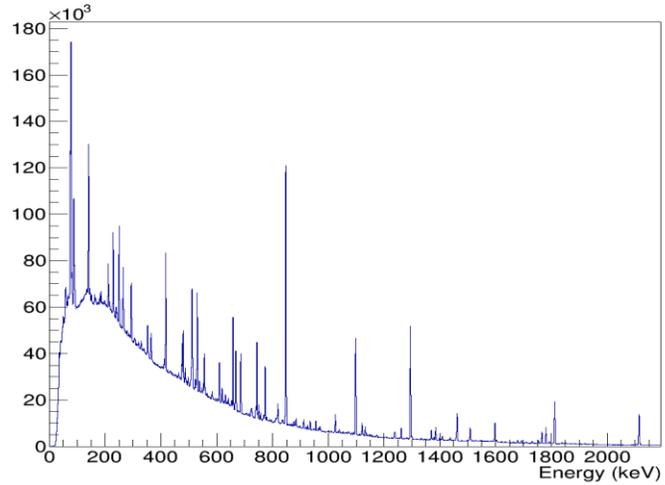
# Experimental setup of LOHENGRIN spectrometer in ILL



• Fission Product  
~ Photon

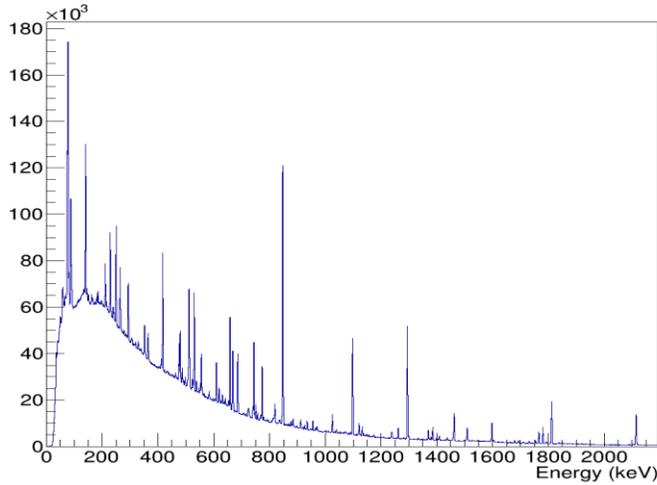
# Analysis - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

Delayed  $\gamma$ -Spectrum without selection



# Analysis - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

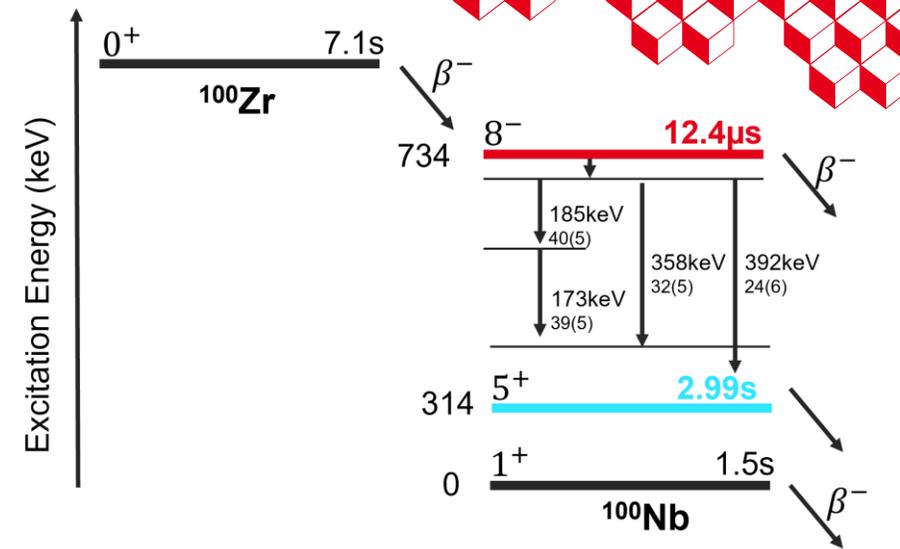
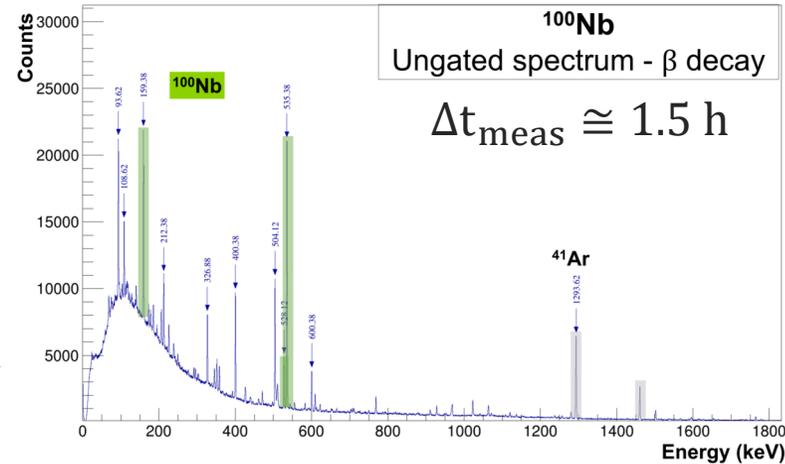
Delayed  $\gamma$ -Spectrum without selection



**LOHENGRIN: mass separator**

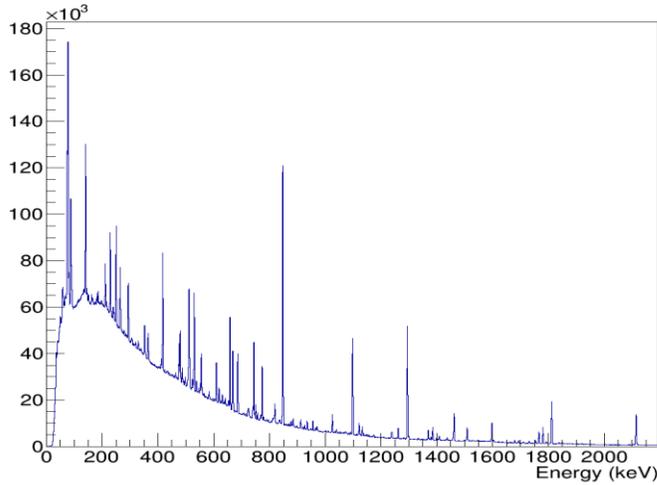
Ungated Spectrum

→  $\beta$ -decay states measurement →  $N_{\gamma}$ ,  $N_{\gamma}$



# Analysis - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

Delayed  $\gamma$ -Spectrum without selection

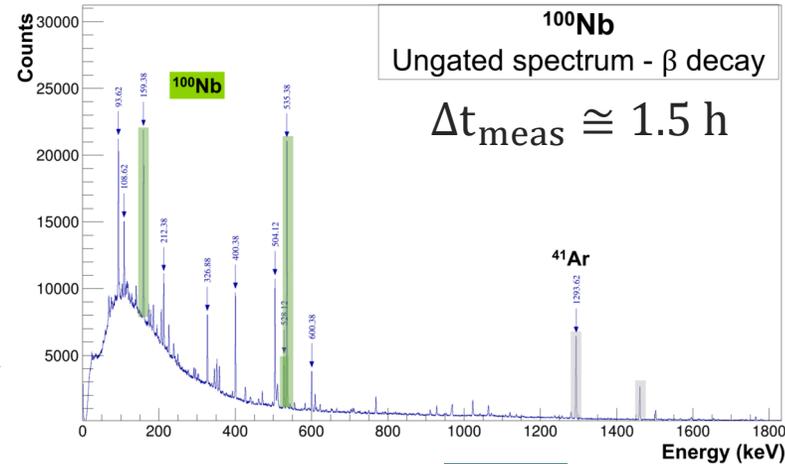


Mass, KE Filter

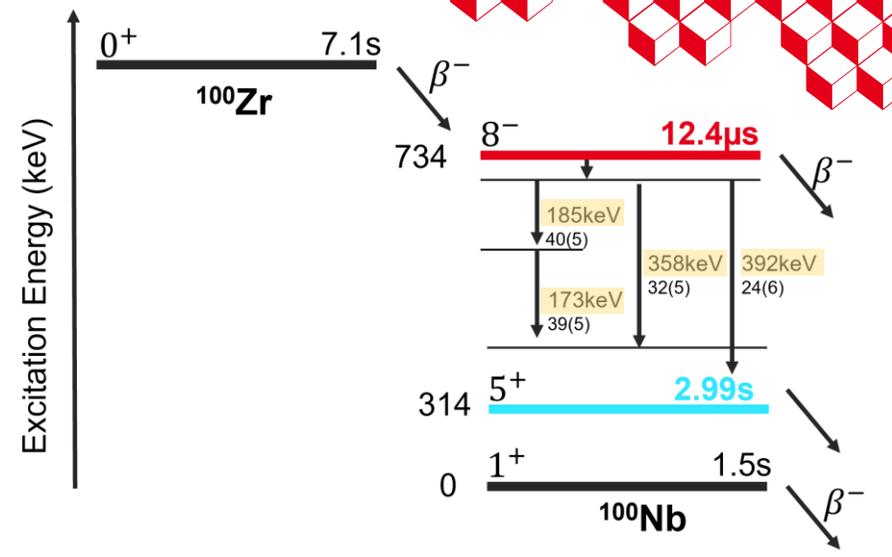
**LOHENGRIN: mass separator**

Ungated Spectrum

→  $\beta$ -decay states measurement →  $N_{\gamma}$ ,  $N_{\gamma}$

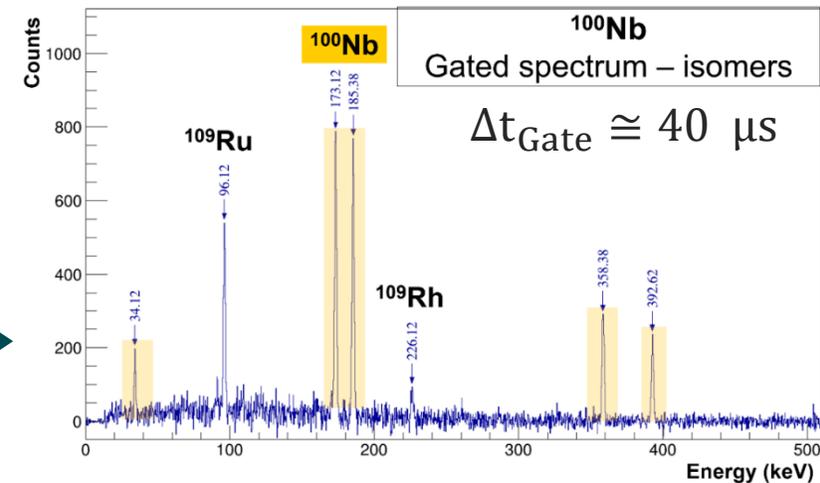


Time Coincidence Filter



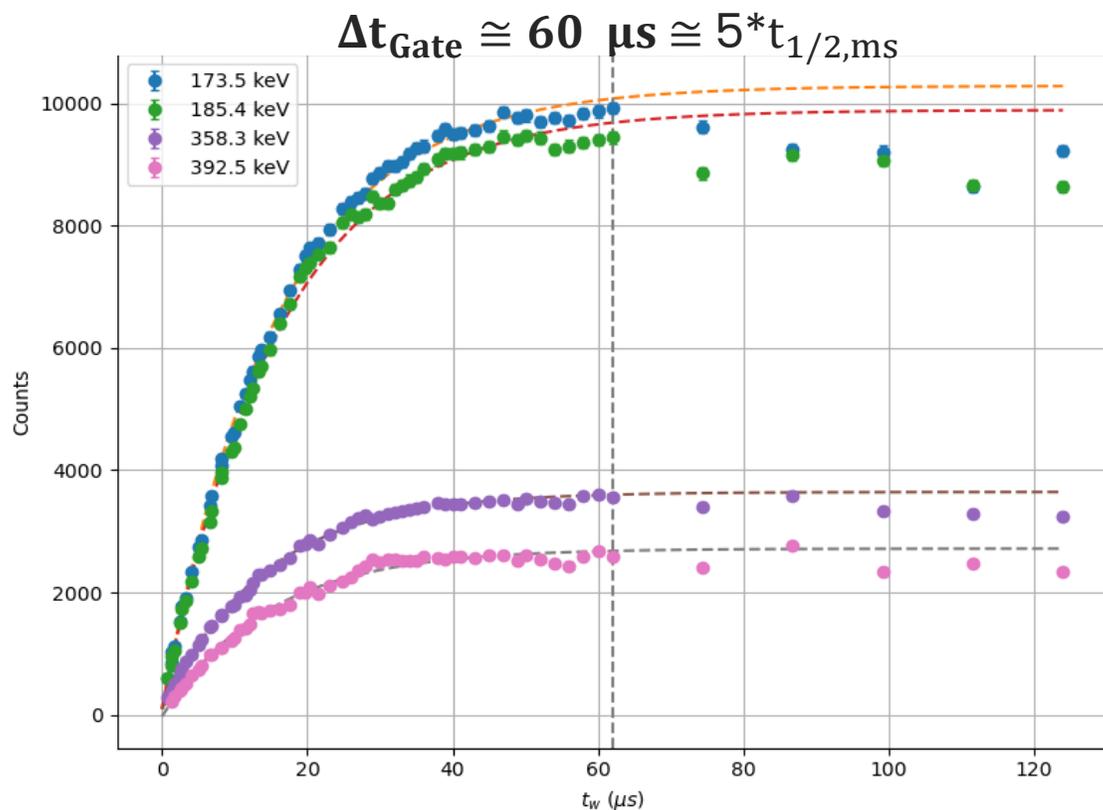
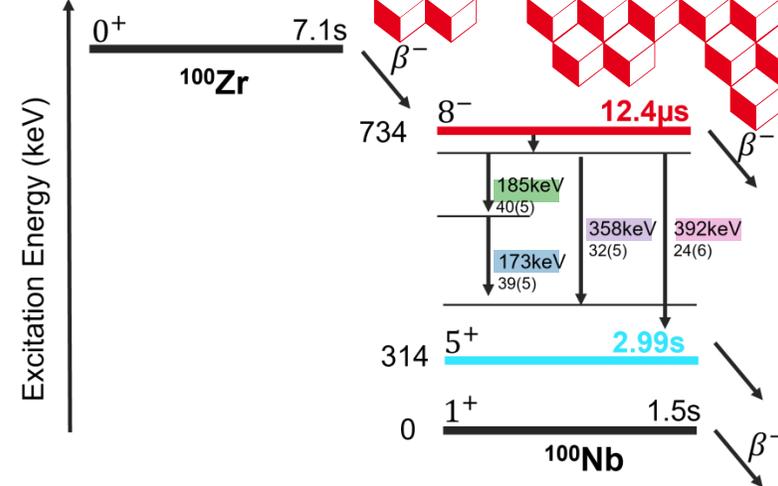
Gated Spectrum

**Time coincidence** between ionization chamber and HPGe detectors  
→  $\mu\text{s}$  Isomeric state measurement →  $N_{\gamma}$



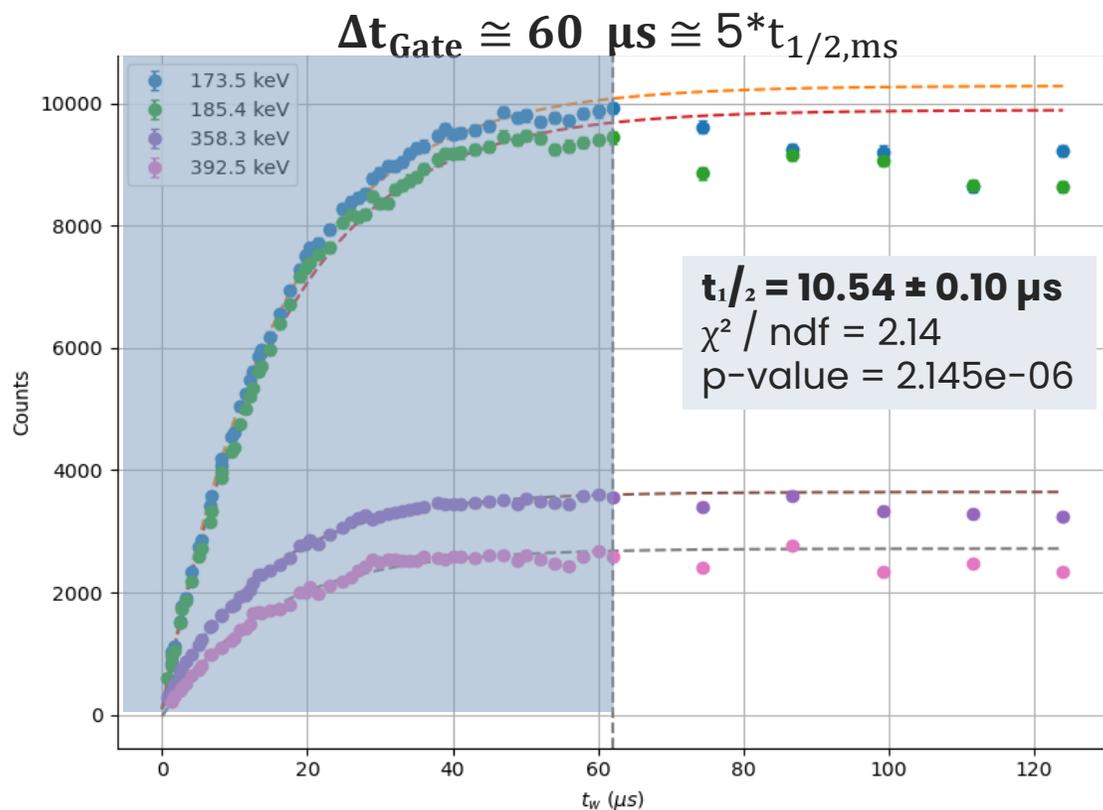
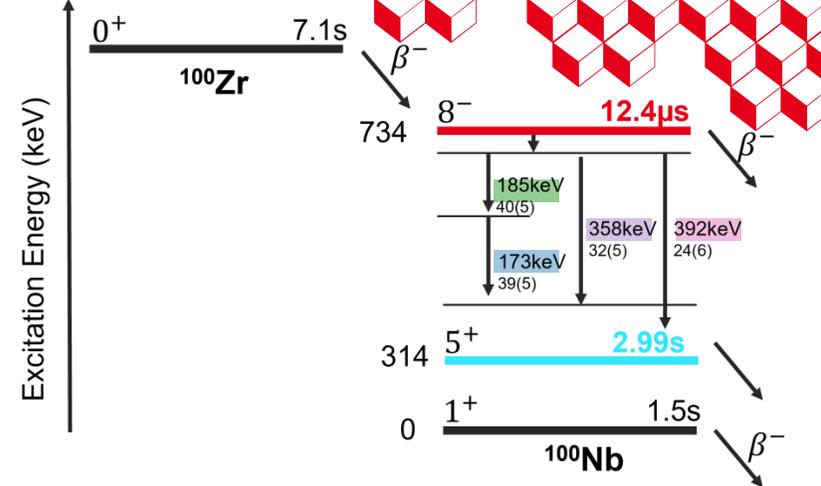
# Analysis - $\Delta t_{\text{Gate}}$ - $^{100}\text{Nb}$

For which  $\Delta t_{\text{Gate}}$  is the  $t_{1/2}$  of the isomeric level re-produced ?



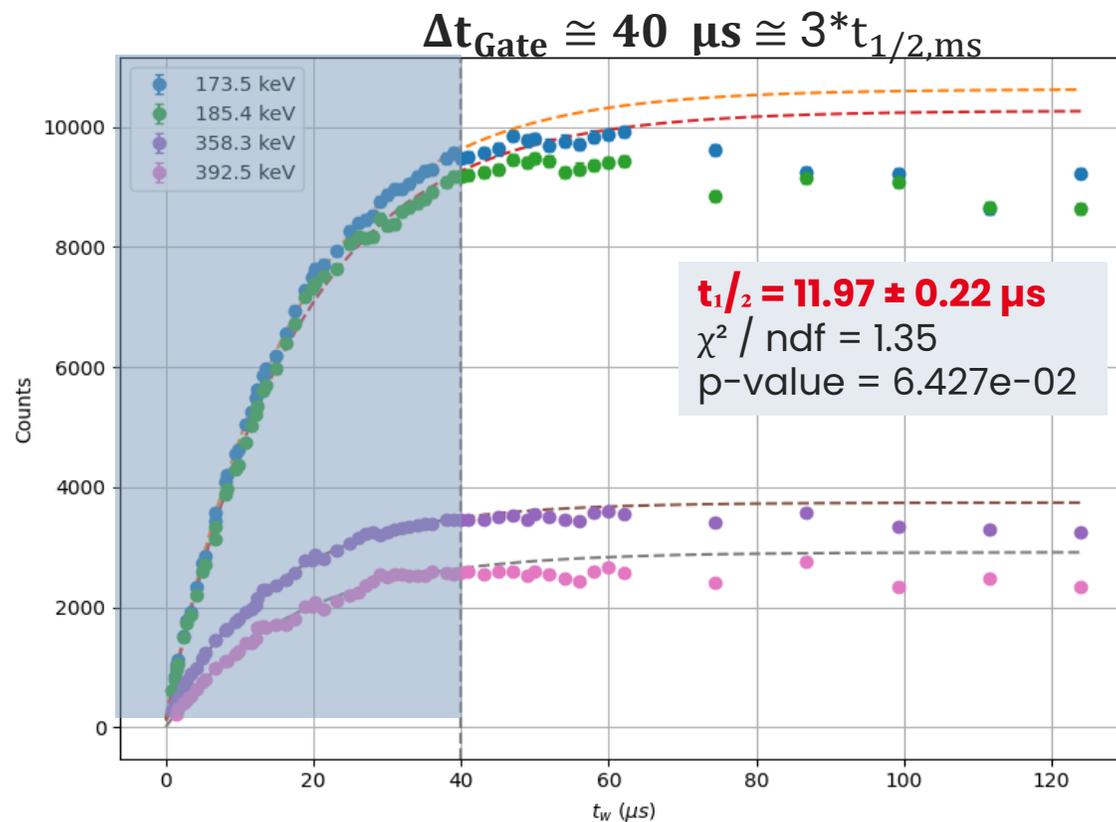
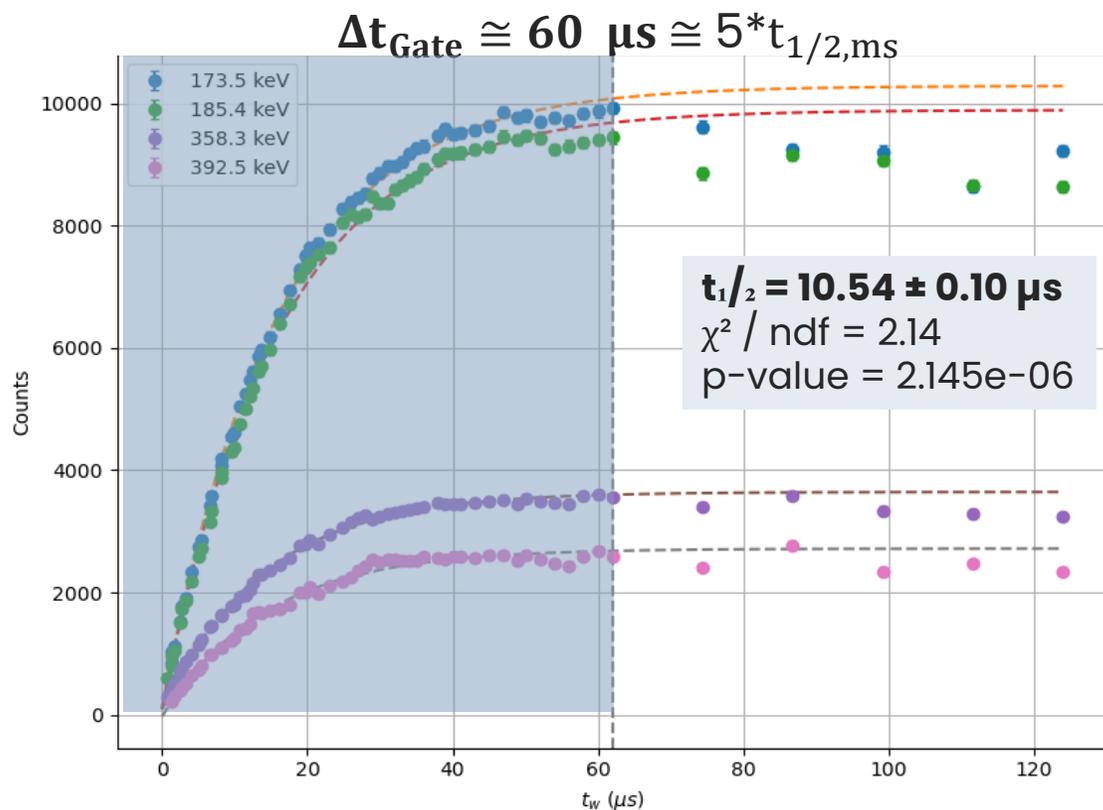
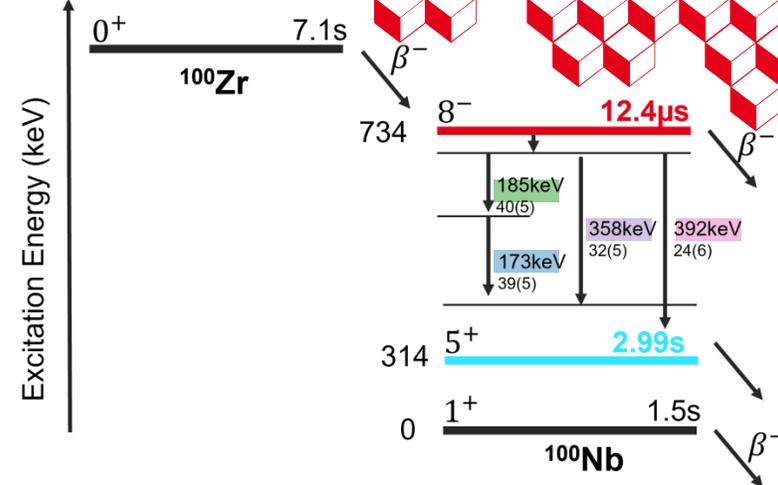
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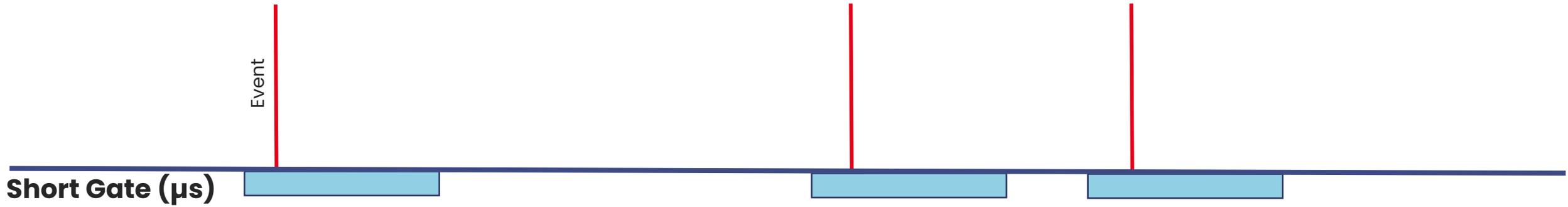


# Analysis - $\Delta t_{\text{Gate}}$ - $^{100}\text{Nb}$

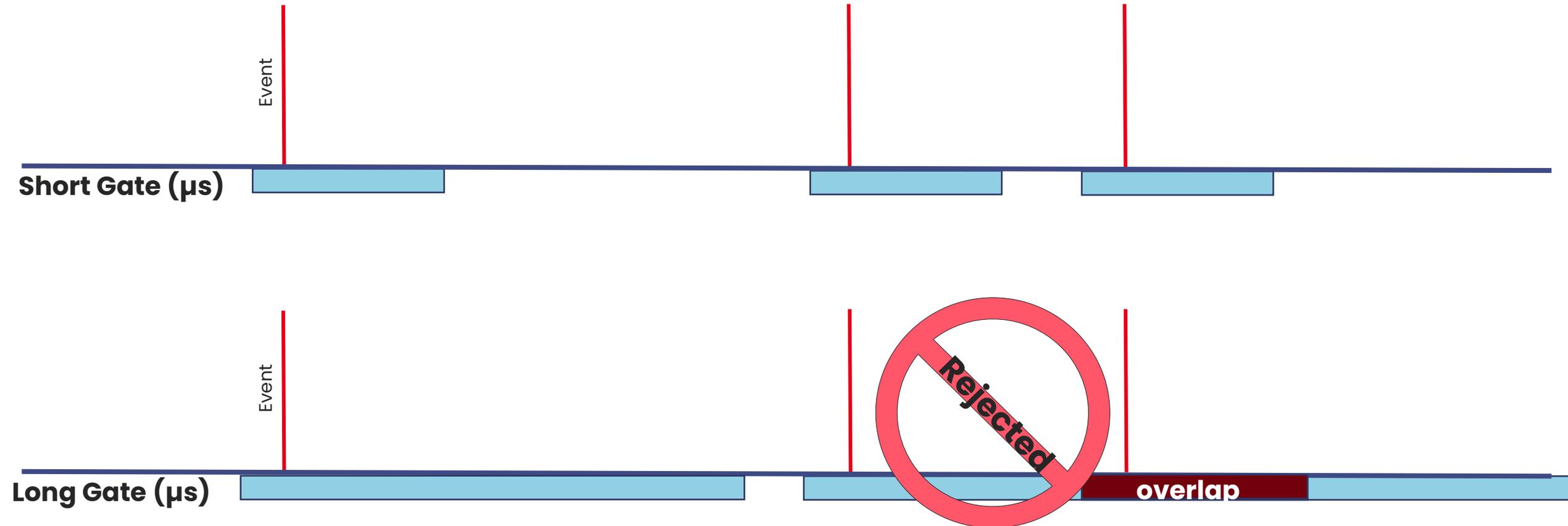
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# Analysis - $\Delta t_{\text{Gate}}$



# Analysis - $\Delta t_{\text{Gate}}$



Importance of data sorting

# Analysis - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

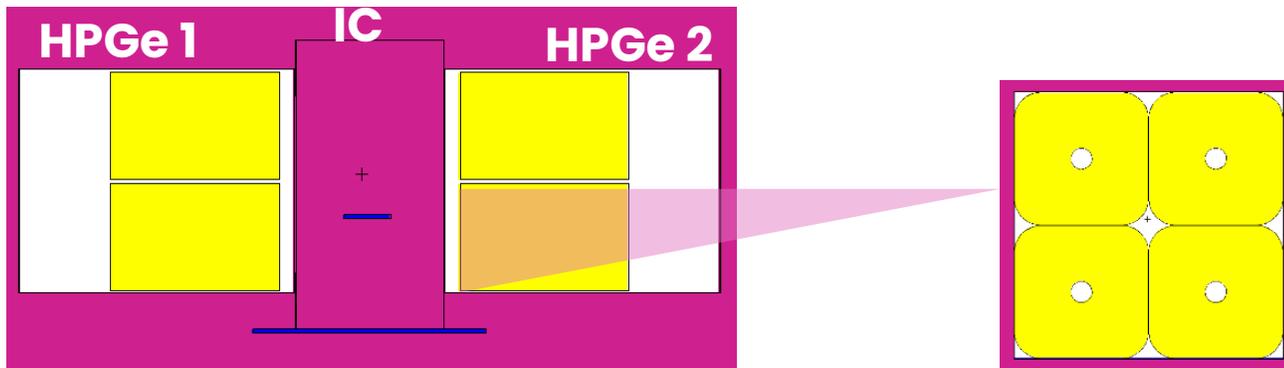
$$N_{\gamma} \rightarrow N_d = \frac{N_{\gamma}}{\epsilon_{\gamma} \times I_{\gamma} \times f_{\gamma}} \xrightarrow{\text{Average}} \overline{N_d} \xrightarrow{\text{Bateman Equation Resolution}} \tau_f \rightarrow \boxed{\text{IR}_{\text{exp}} = \frac{\tau_f(\text{IS})}{\tau_f(\text{IS}) + \tau_f(\text{GS})}}$$

IR: Isomeric Ratio, IS: Isomeric State, GS: Ground State

$\epsilon_{\gamma}$ : Detector Efficiency,

$f_{\gamma}$ : Correction factor for Sum-Effect,

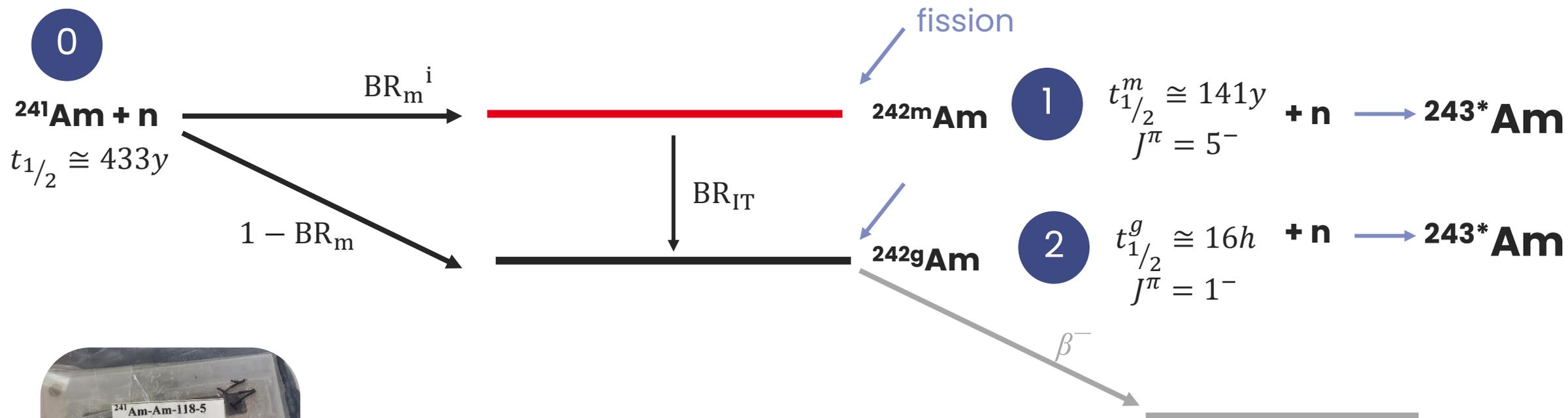
will be derived from **TrueCoinc**<sup>1</sup> and Monte Carlo Simulation **MCNP6**<sup>2</sup> (in progress).



1. Sudar, S. (2002). 'TrueCoinc' software utility for calculation of the true coincidence correction (IAEA-TECDOC--1275). International Atomic Energy Agency (IAEA)

2. D.B. Pelowitz *et al.*, Tech. Rep. LA-CP-13-00634, Los Alamos National Laboratory, Los Alamos, NM, USA, 2013

# $^{241}\text{Am}(2n_{\text{th}},f)$ mechanism

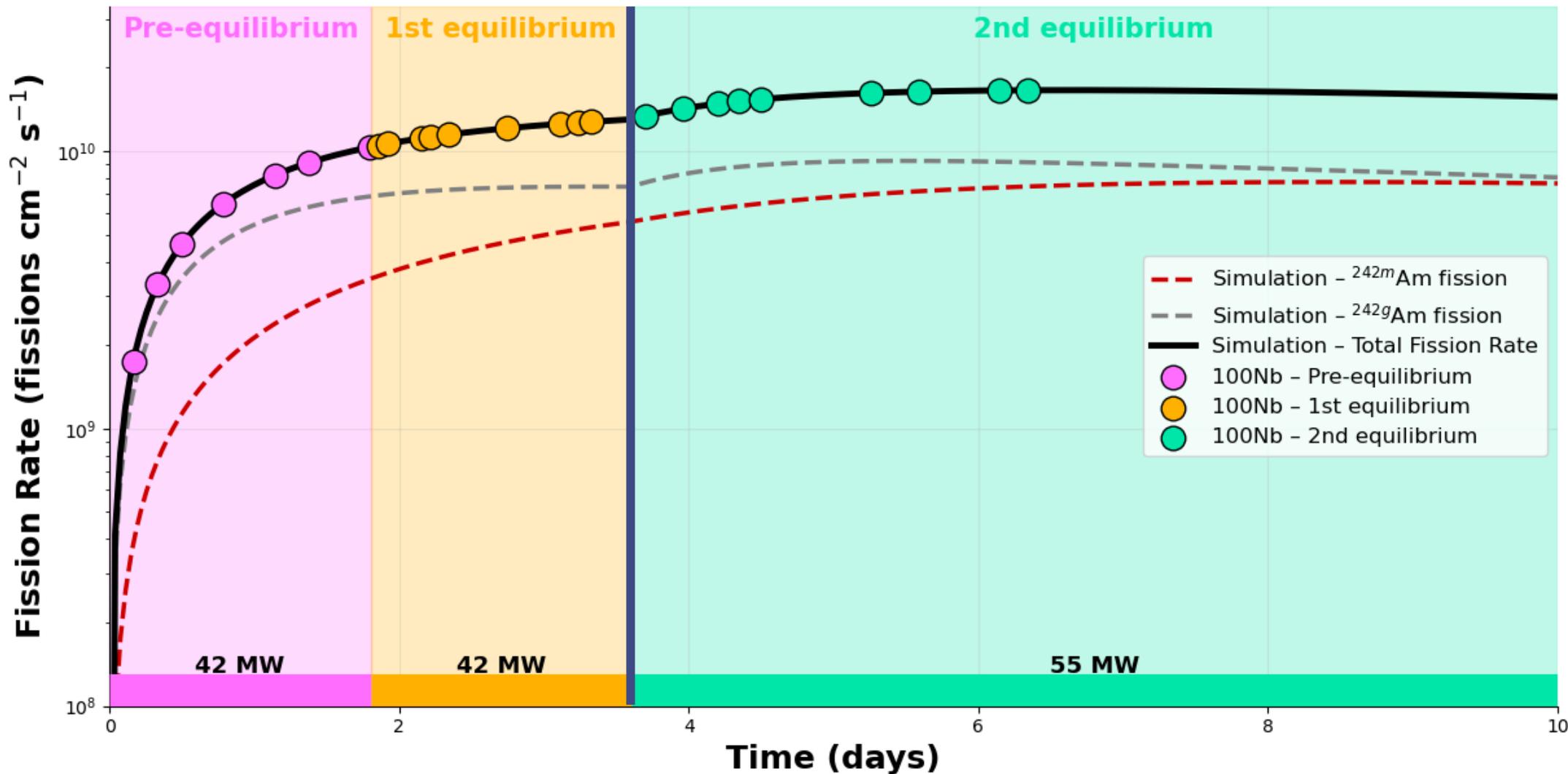


The branching ratio of  $^{241}\text{Am}$  is spectrum dependent

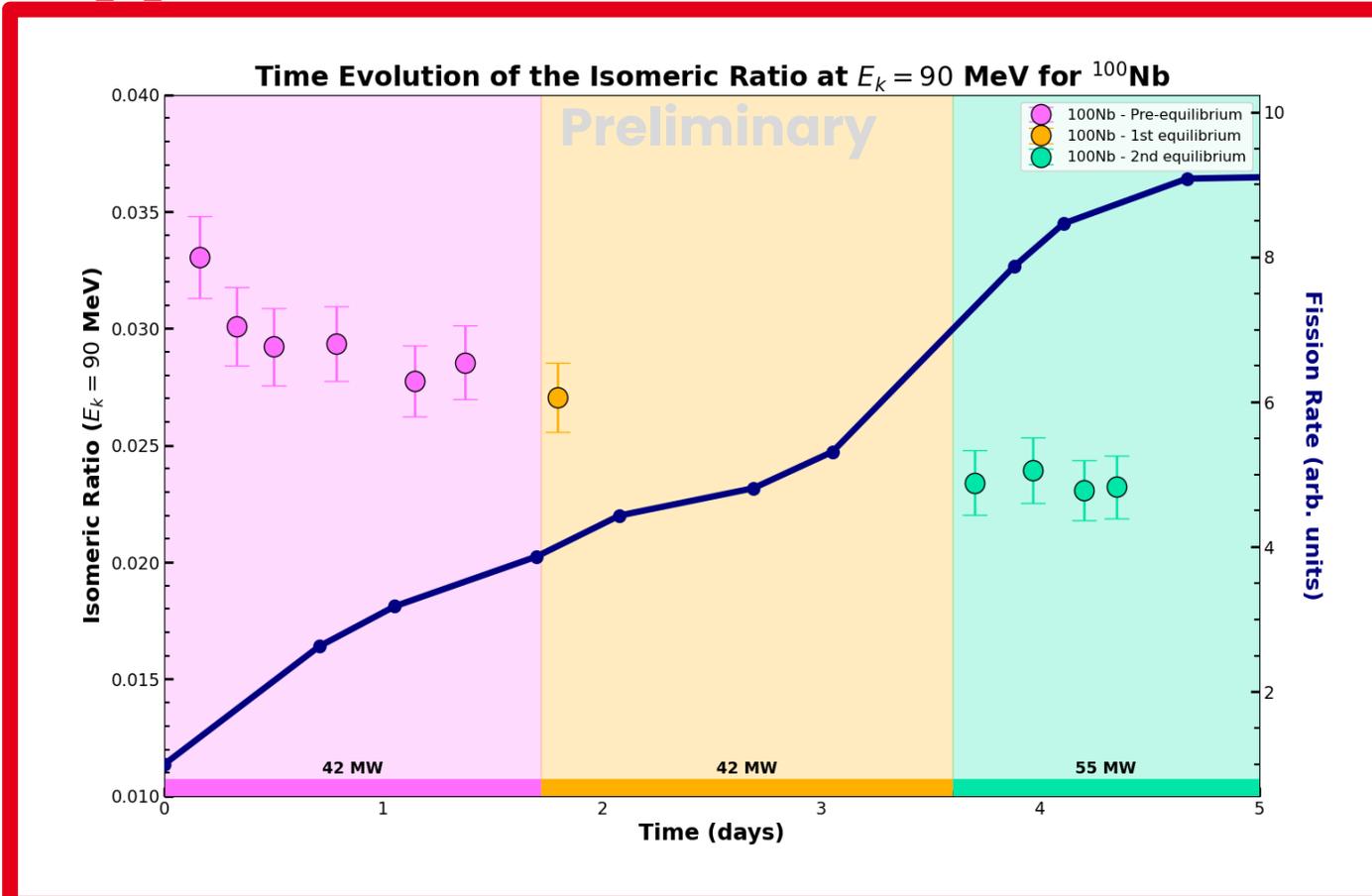
$$i : BR_m = 0.093$$

# Preliminary Results - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

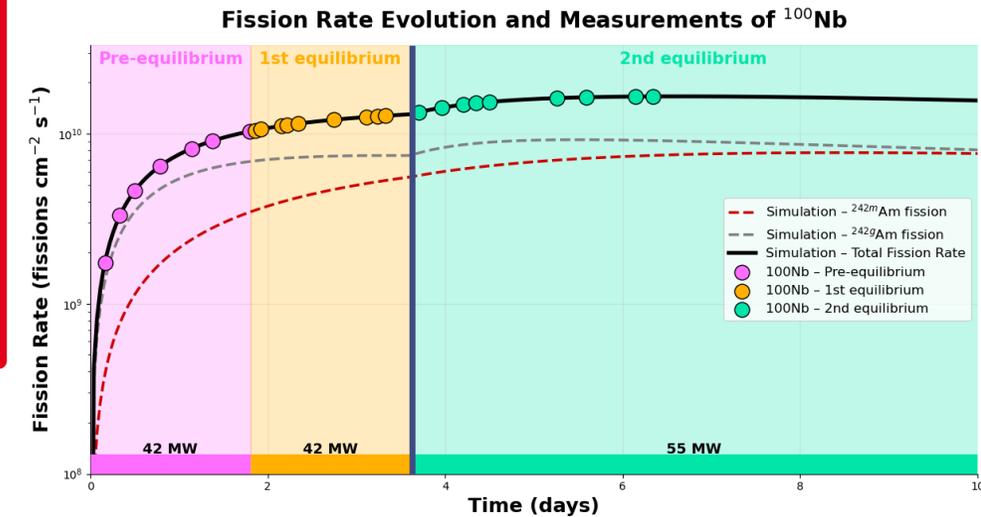
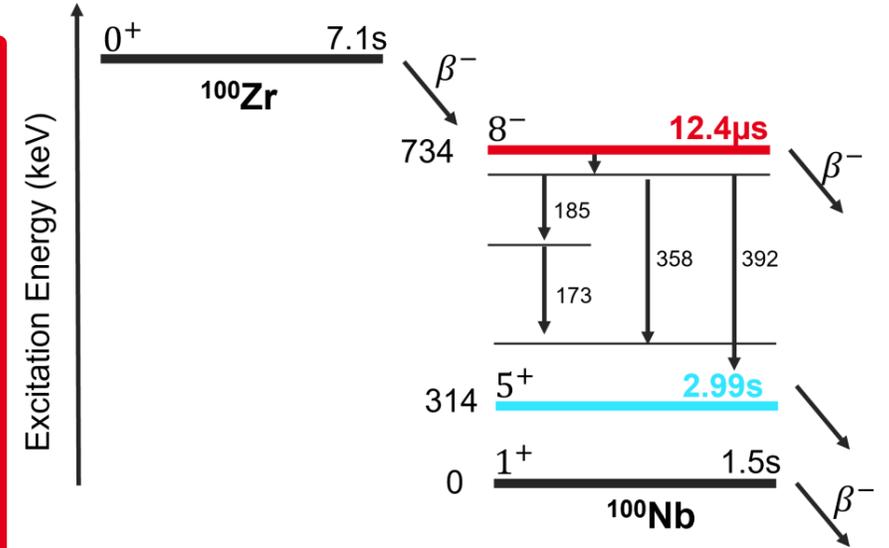
## Fission Rate Evolution and Measurements of $^{100}\text{Nb}$



# Preliminary Results - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$ IR(t)

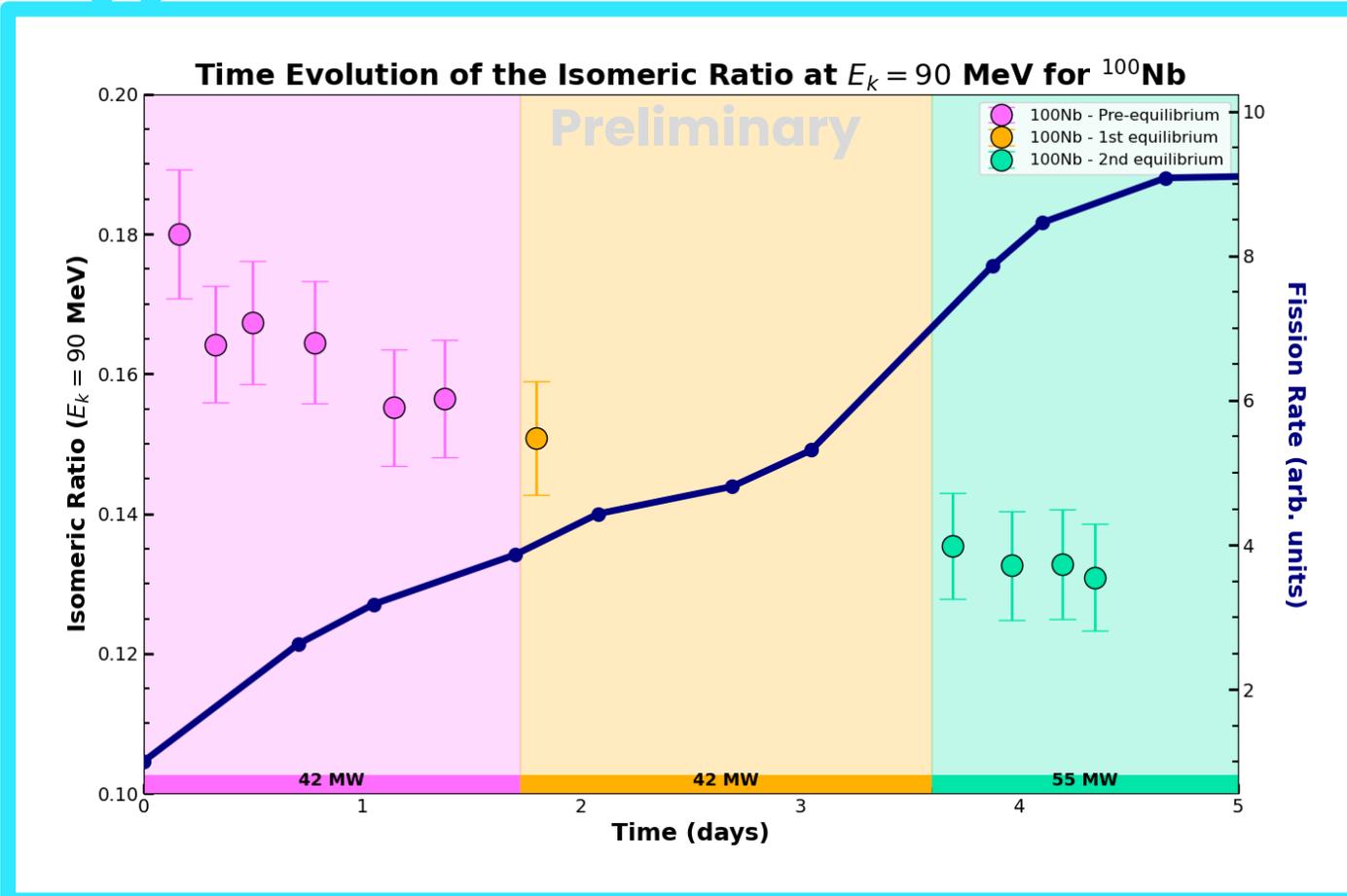


$J^\pi = 8^-$   
 $t_{1/2}^m \cong 12 \mu\text{s}$

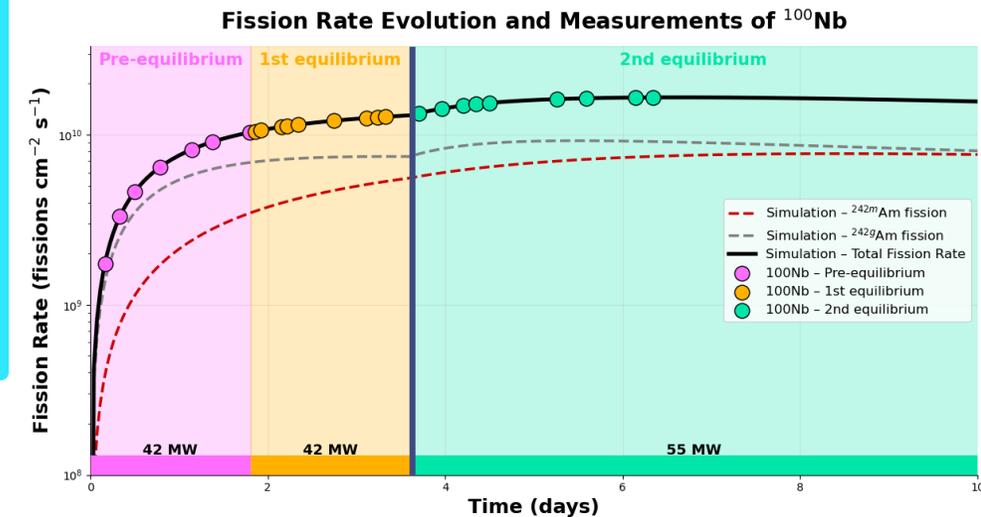
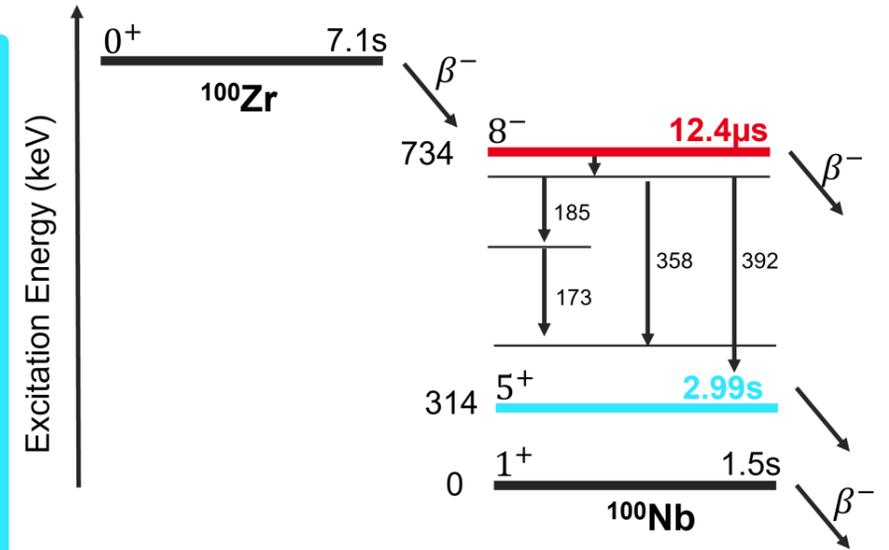


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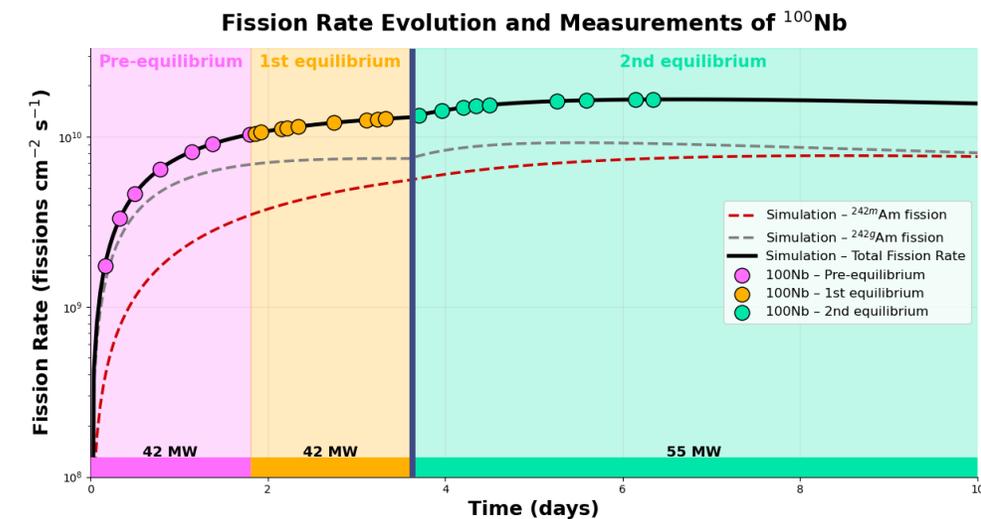
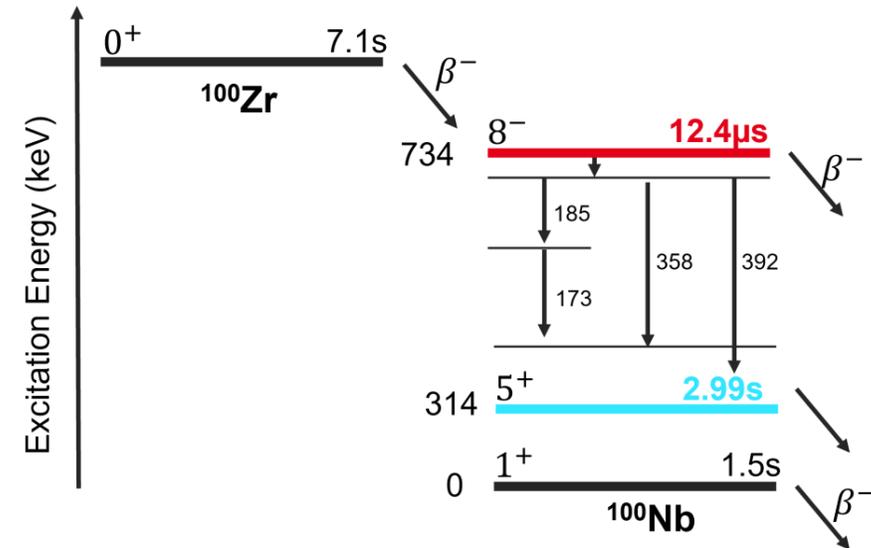
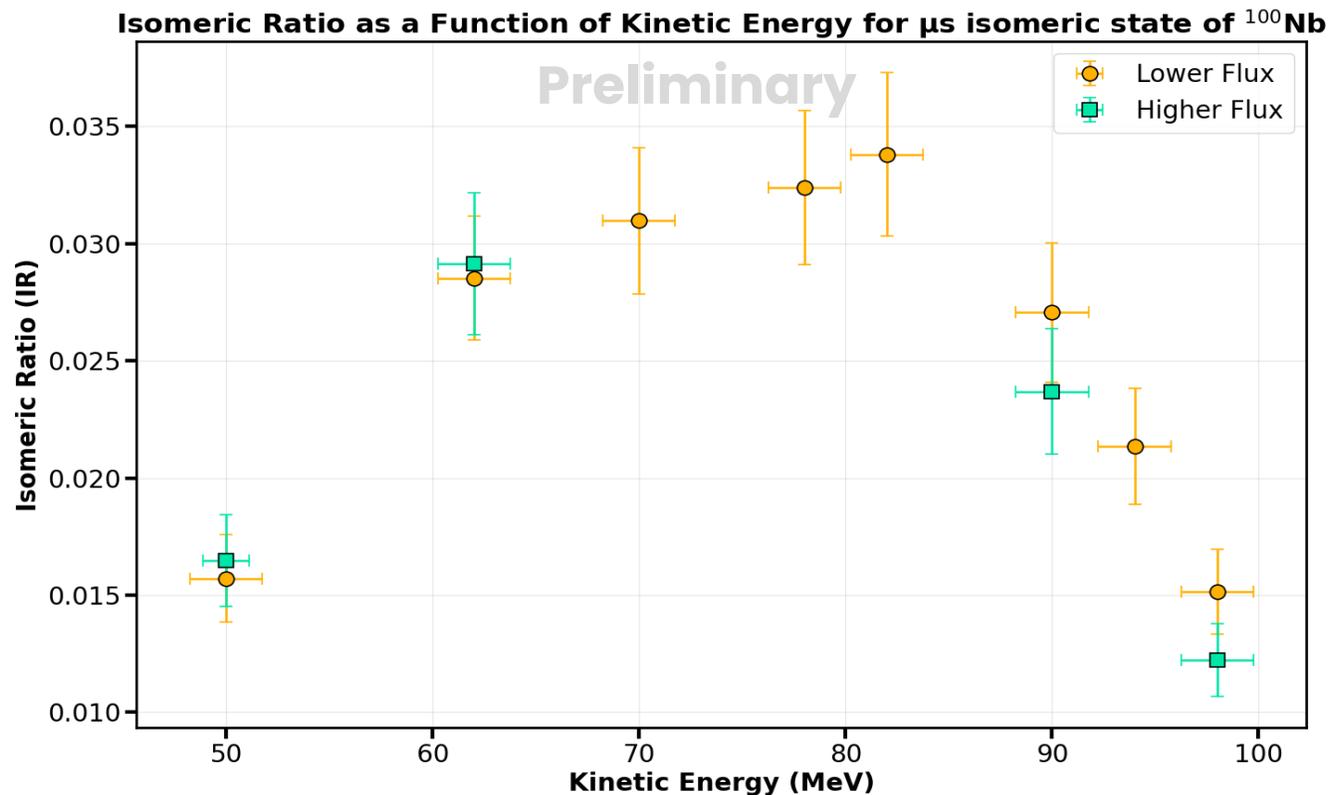
## IR(t)



$J^\pi = 5^+$   
 $t_{1/2}^m \cong 3$  s



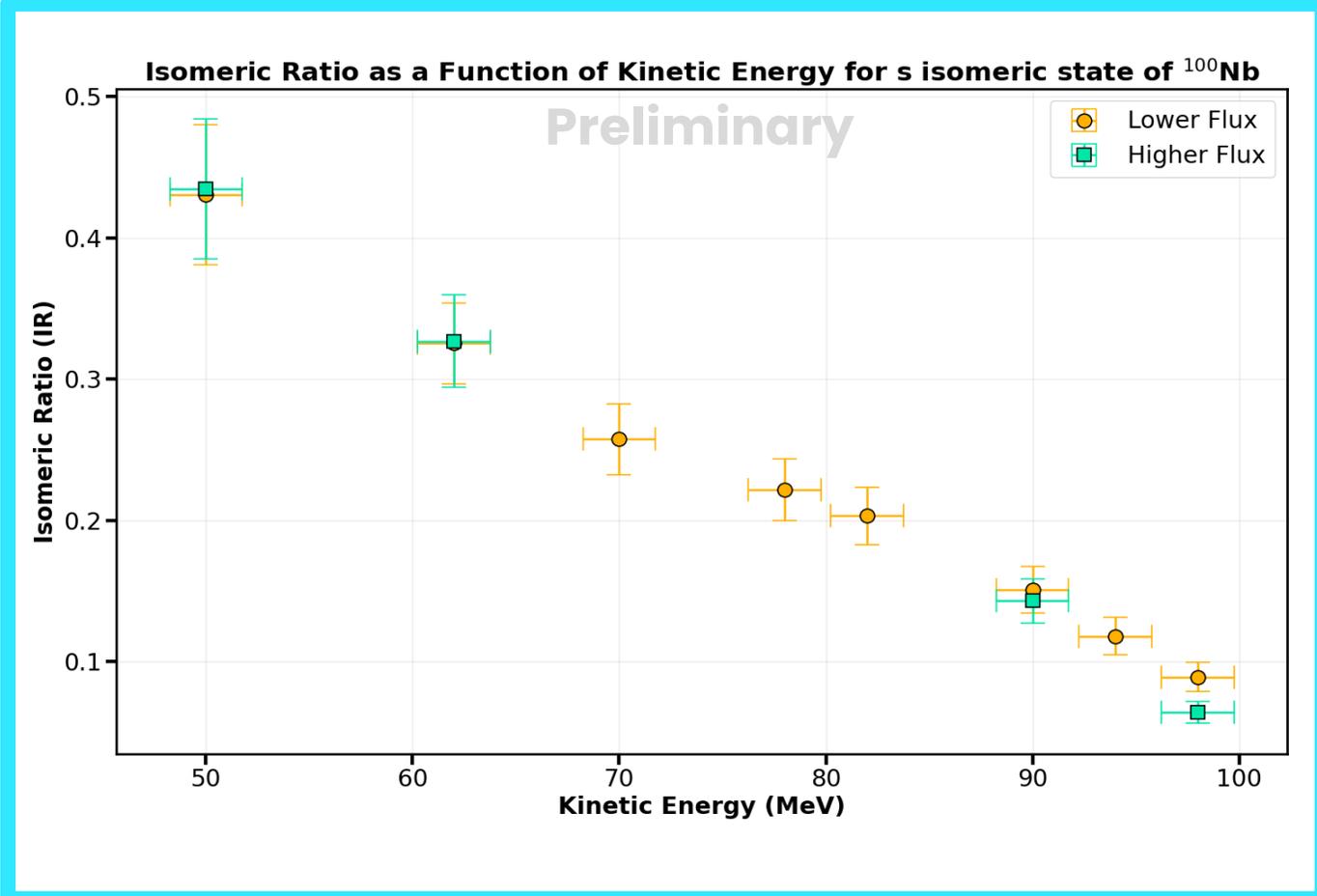
# Preliminary Results - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$ IR(KE)



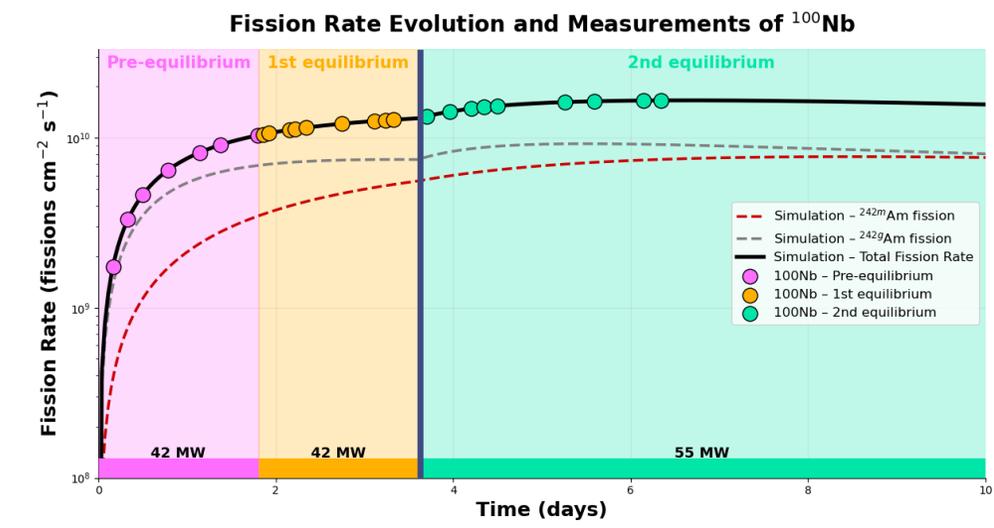
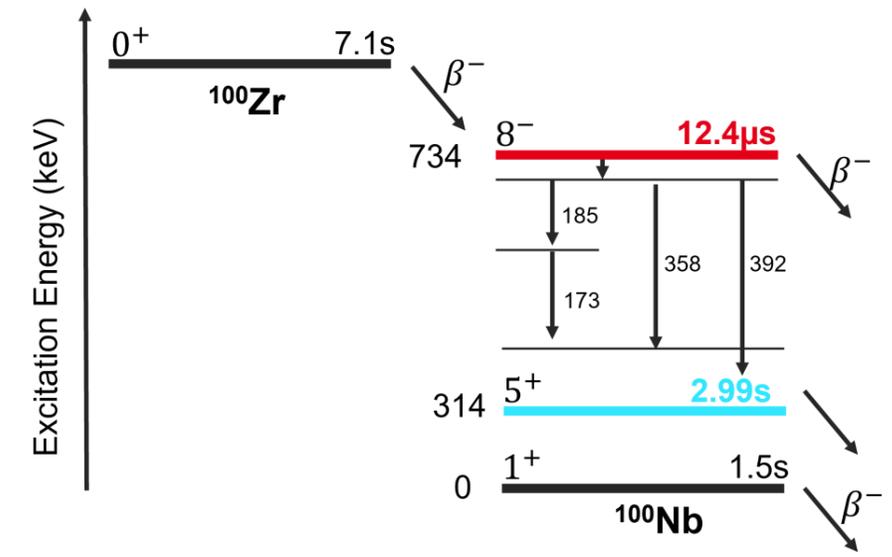
$J^\pi = 8^-$   
 $t_{1/2}^m \cong 12 \mu\text{s}$

# Preliminary Results - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

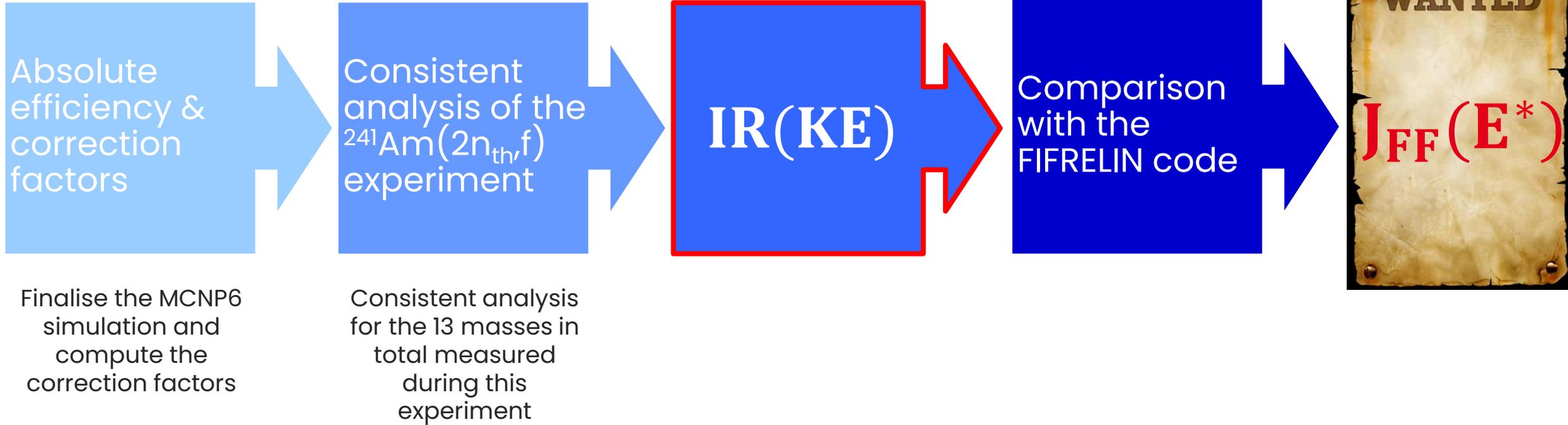
## IR(KE)



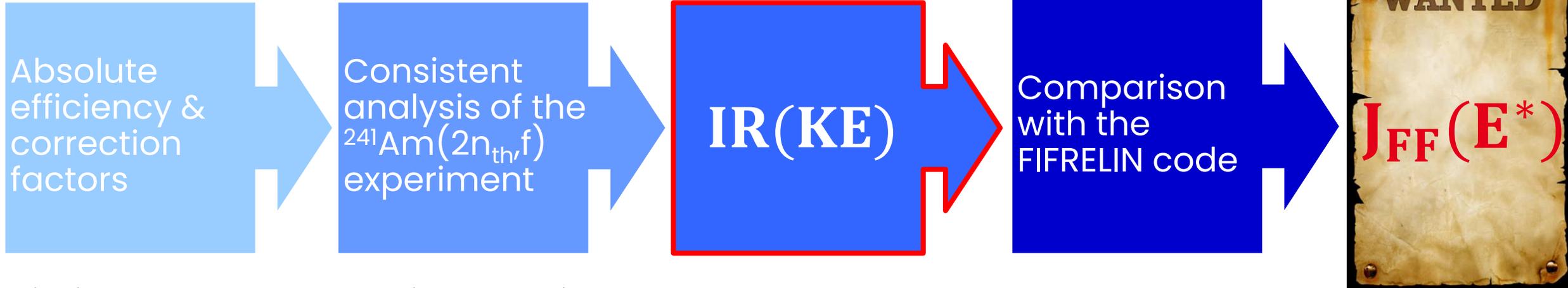
$J^\pi = 5^+$   
 $t_{1/2}^m \cong 3 \text{ s}$



# Future prospects



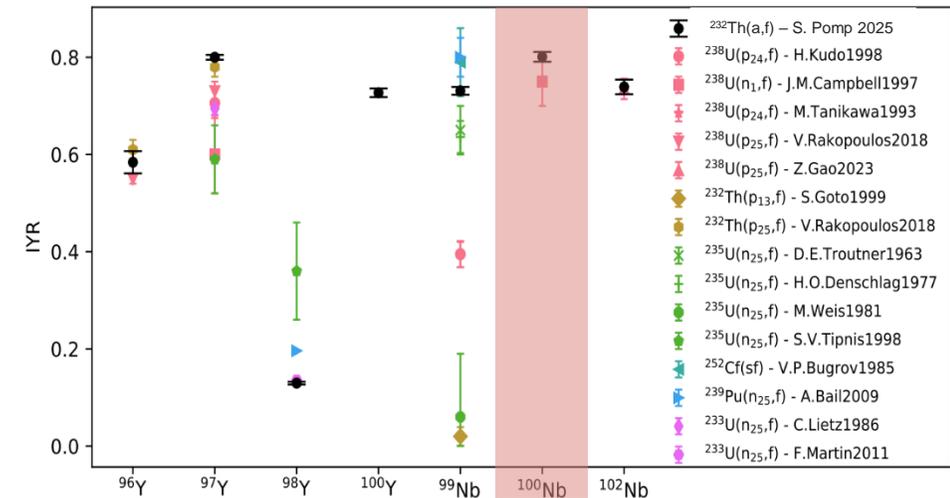
# Future prospects

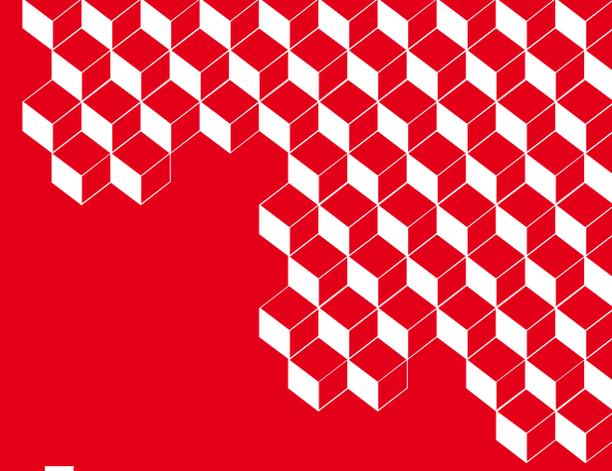


Finalise the MCNP6 simulation and compute the correction factors

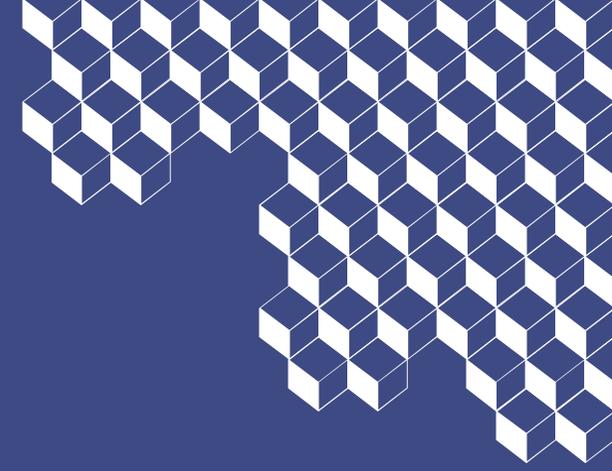
Consistent analysis for the 13 masses in total measured during this experiment

**+** Upcoming experiment  $^{237}\text{Np}(2n_{\text{th}},f)$  & Analysis of existing data for  $^{239,241}\text{Pu}(n_{\text{th}},f)$





**Thank you for your attention !**

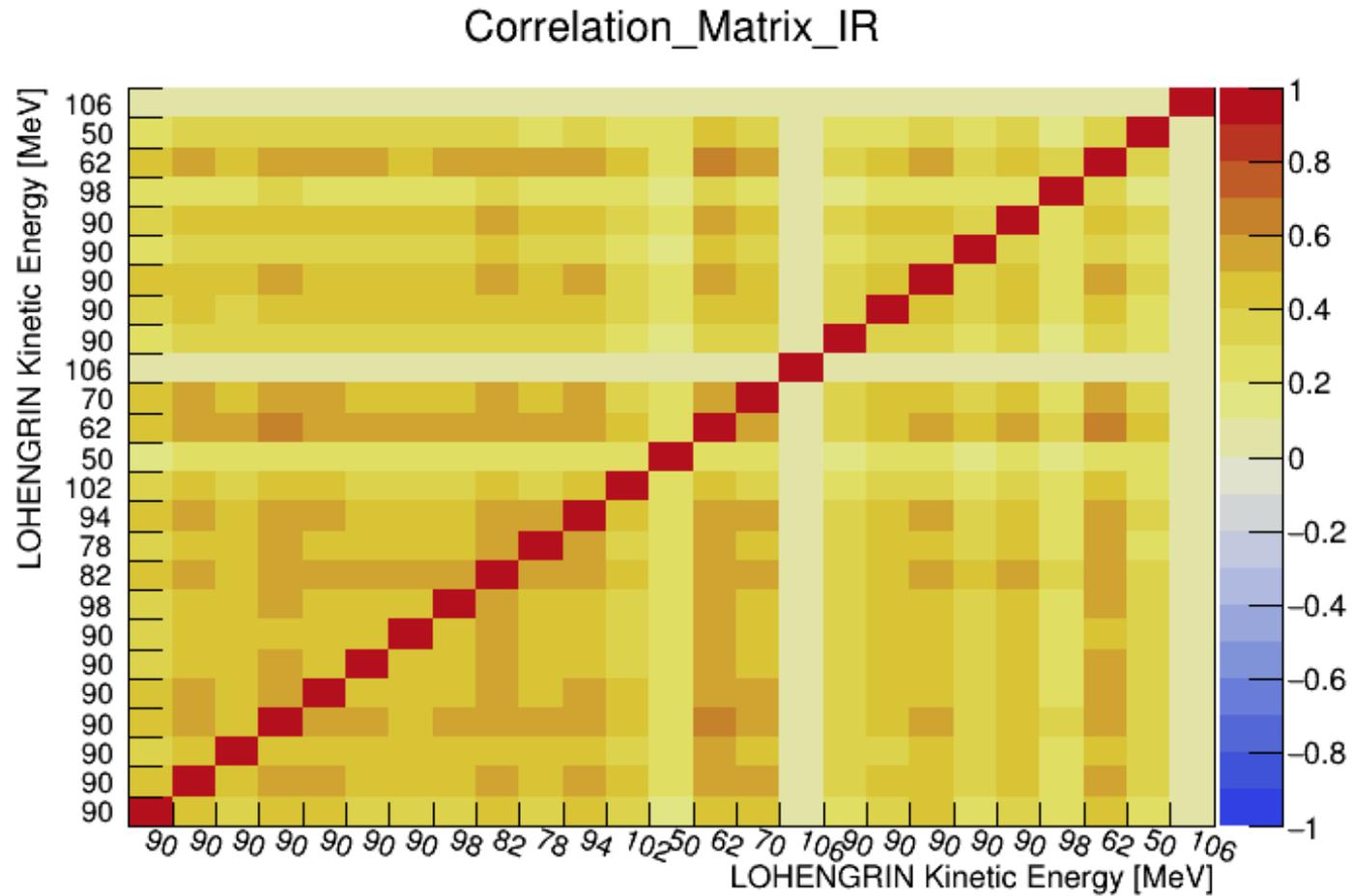


# // Back-up

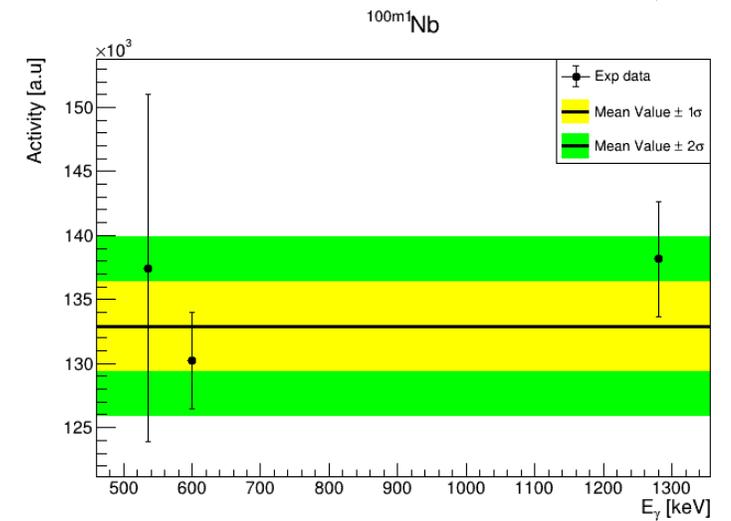
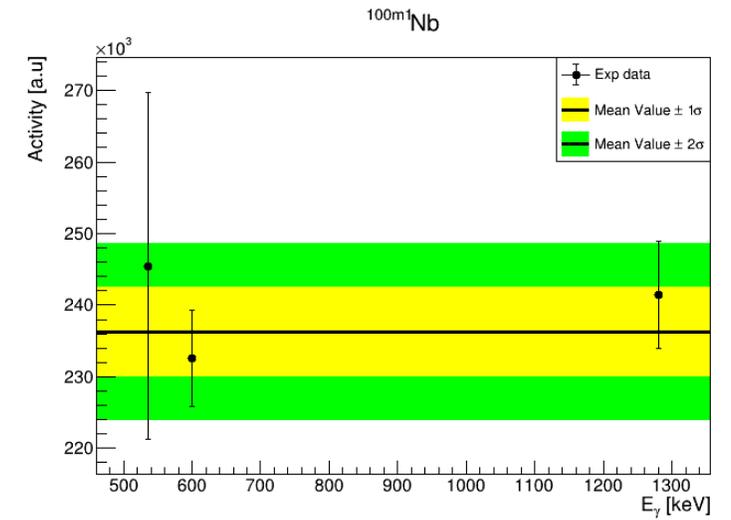
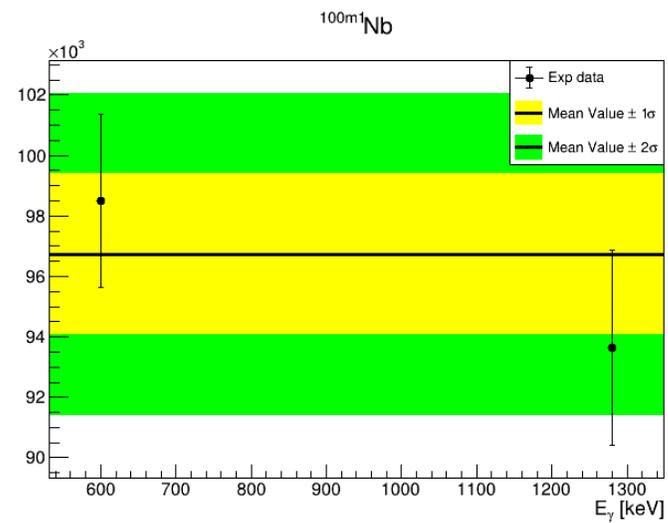
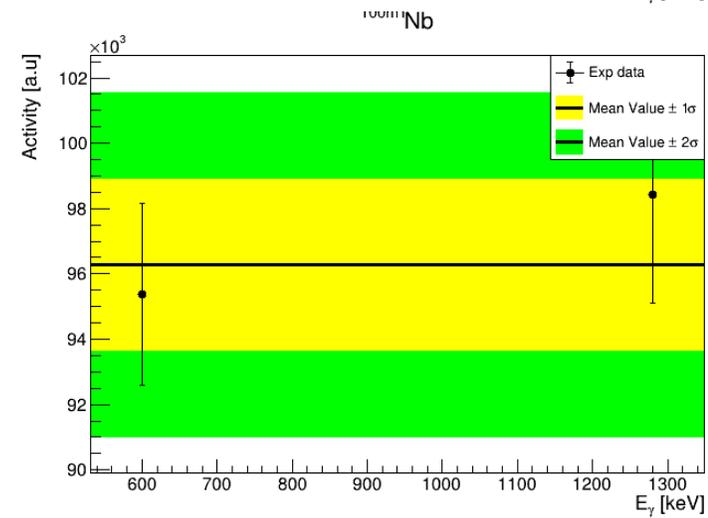
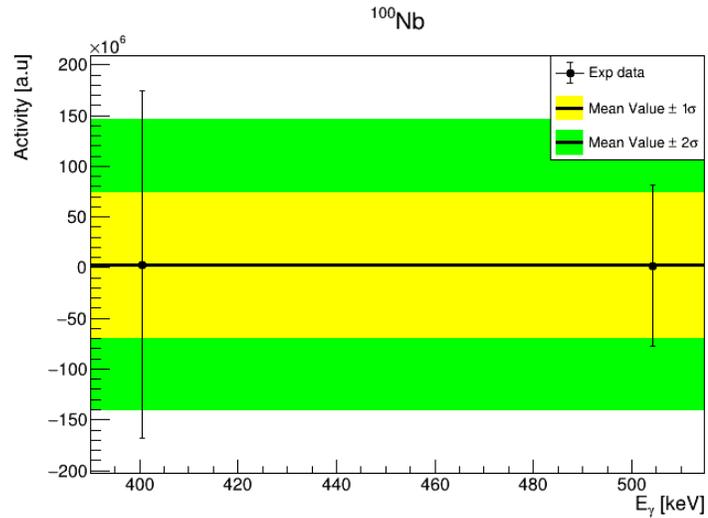
# Isotopes measured at Lohengrin experiment – Oct. 2025

| Isotope | $t_{1/2}$ ( $\mu$ s) level (NNDS) |
|---------|-----------------------------------|
| 88Br    | 5.3                               |
| 94Y     | 1.35                              |
| 98Y     | 0.78                              |
|         | 0.45                              |
|         | 6.9                               |
|         | 0.69                              |
| 99Y     | 8                                 |
| 100Nb   | 12.4                              |
| 109Ru   | 0.68                              |
| 126In   | 26                                |
| 128Sn   | 2.91                              |
| 129Sn   | 2.22                              |
|         | 3.40                              |
| 130Sn   | 1.61                              |
| 131Sb   | 1.1                               |
|         | 4.3                               |
|         | 91                                |
| 132Te   | 3.7                               |
|         | 28.1                              |

# Correlation Matrix IR(KE)

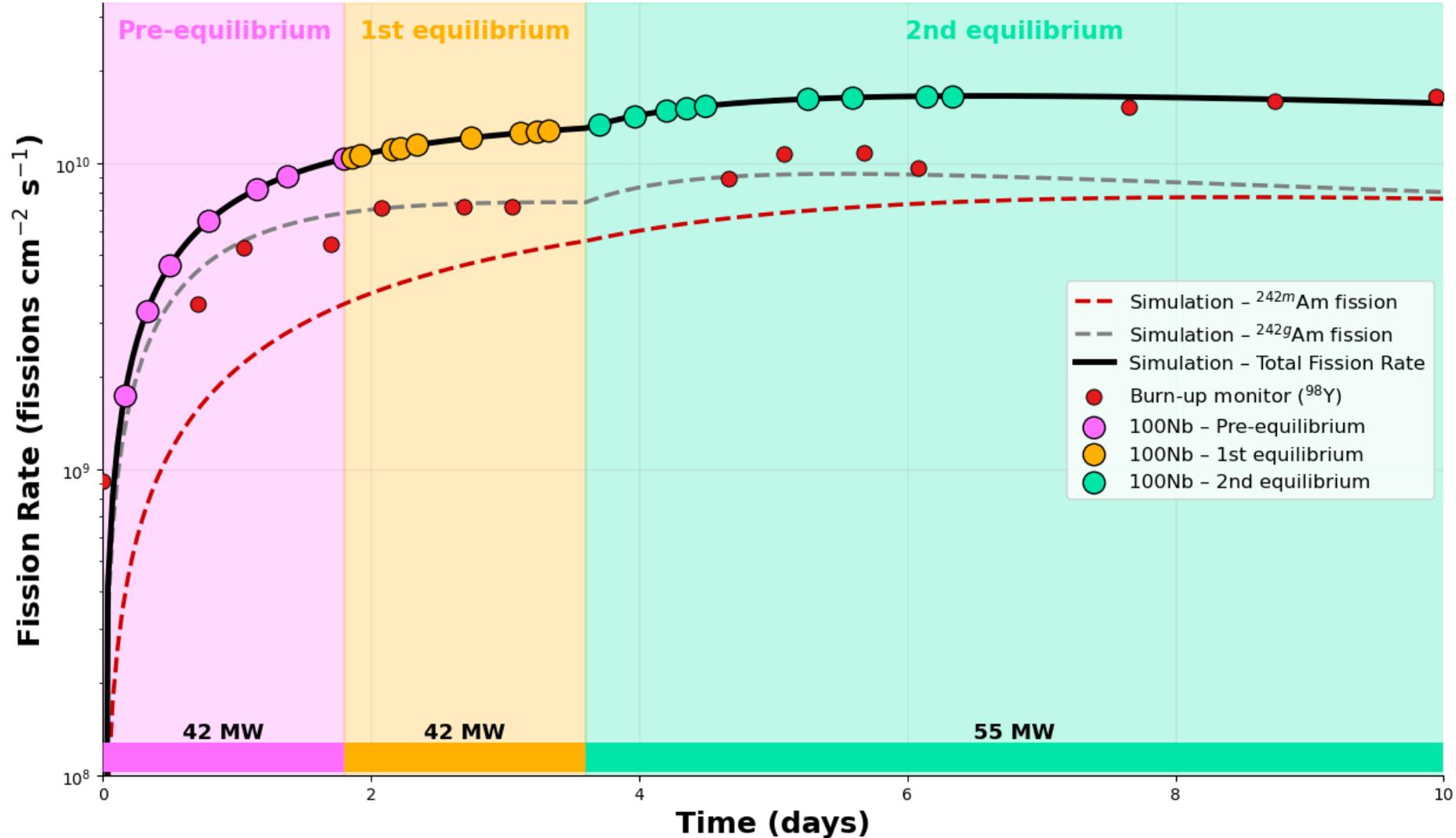


# Analysis



# Comparison with theoretical model (Bateman Eq.)

## Fission Rate Evolution and Measurements of $^{100}\text{Nb}$



# The Bateman Equation

$$\frac{dN_0}{dt} = -a_0 \times N_0$$

$$\frac{dN_1}{dt} = -a_1 \times N_1 + BR_m \times CS_{c,0} \times F \times N_0$$

$$\frac{dN_2}{dt} = -a_2 \times N_2 + (1 - BR_m) \times CS_{c,0} \times F \times N_0 + \underbrace{BR_{IT} \times \lambda_1 \times N_1}_{\text{negligible}}$$

Where

$$a_i = \lambda_i + (CS_{f,i} \times F) + (CS_{c,i} \times F)$$

$\lambda_i$  : due to the decay

$F$  : Flux

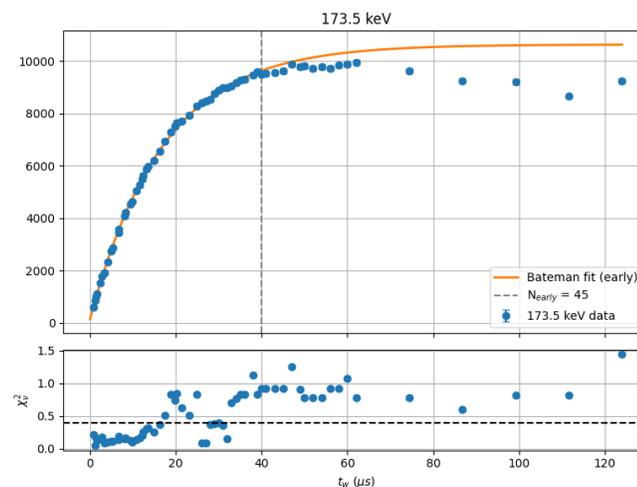
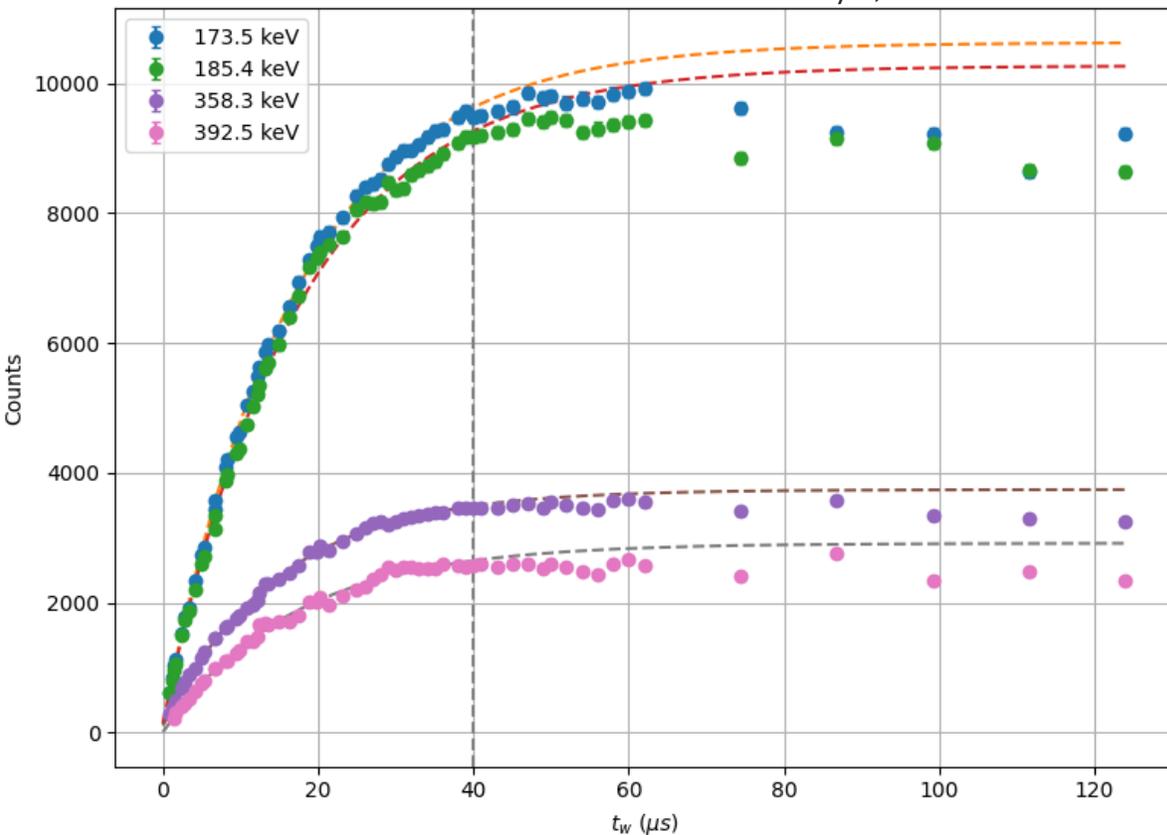
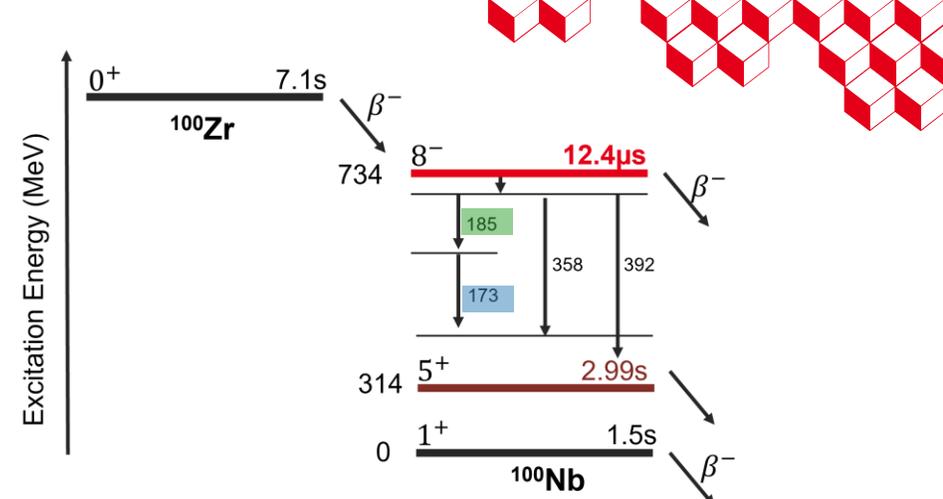
$CS_{f,i}$  : Cross Section due to fission

$CS_{c,i}$  : Cross Section due to capture

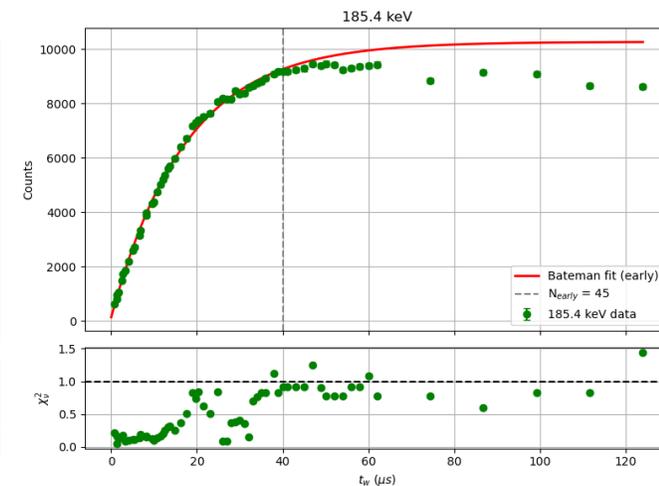
# Analysis - $\Delta t_{\text{Gate}} = {}^{100}\text{Nb}$

For which  $\Delta t_{\text{Gate}}$  is the  $t_{1/2}$  of the isomeric level re-produced?

$$\Delta t_{\text{Gate}} \cong 40 \mu\text{s} \cong 3 * t_{1/2, \text{ms}}$$



$t_{1/2} = 11.97 \pm 0.22 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 1.35$   
 p-value = 6.427e-02



$t_{1/2} = 12.27 \pm 0.21 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 3.06$   
 p-value = 9.947e-11

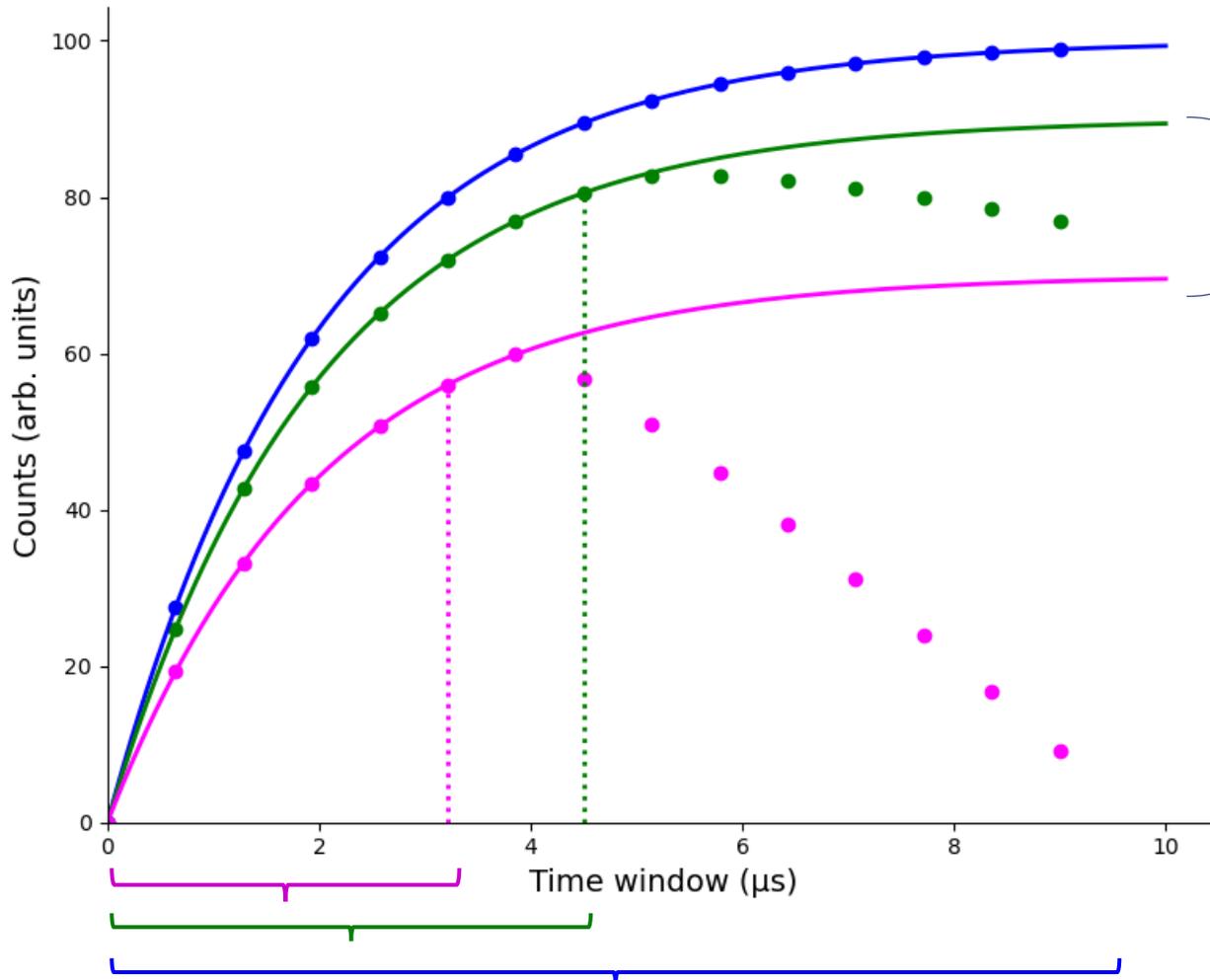
Importance of data shorting

# Analysis - $\Delta t_{\text{Gate}}$

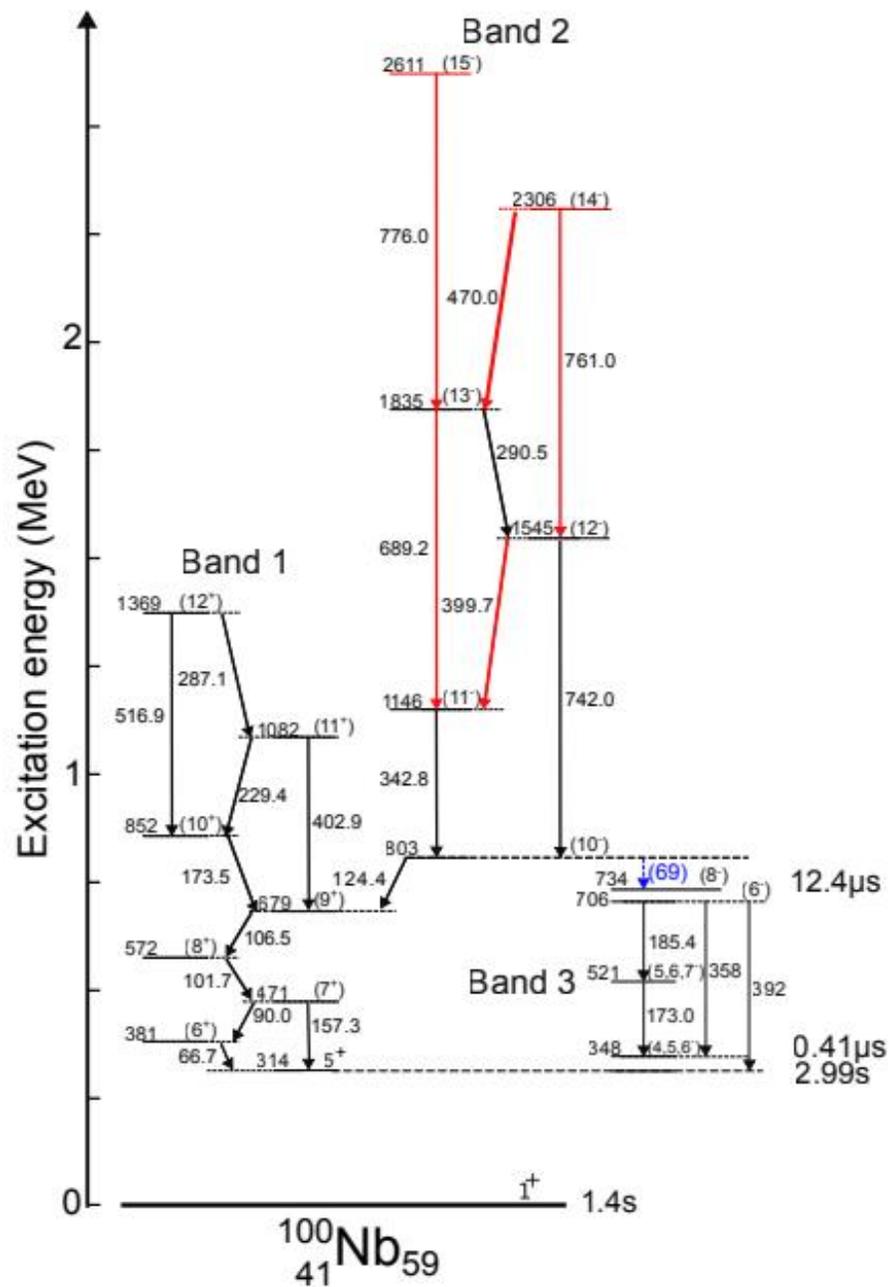
## Different cases:

Short half-life of isomeric level ( $< 10\mu\text{s}$ )  $\rightarrow$  no correction needed

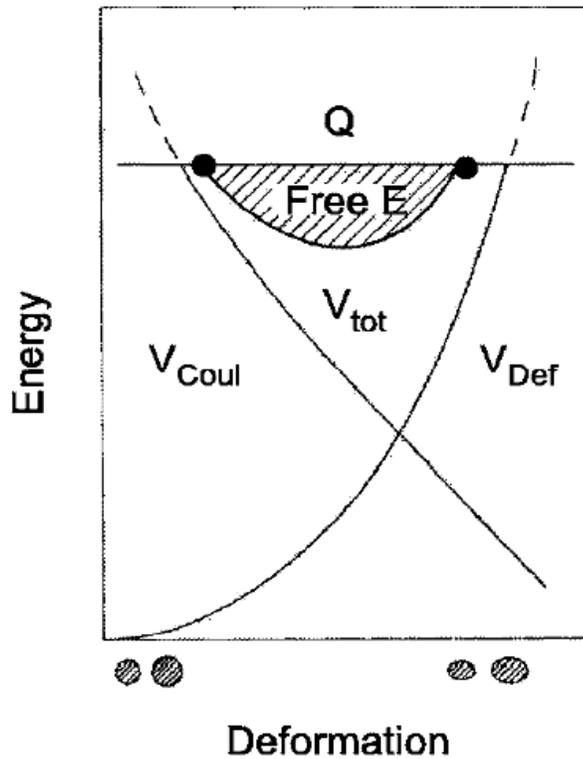
Longer half-life of isomeric level ( $> 10\mu\text{s}$ )  $\rightarrow$  correction needed  $\rightarrow f_{\text{gate}}$



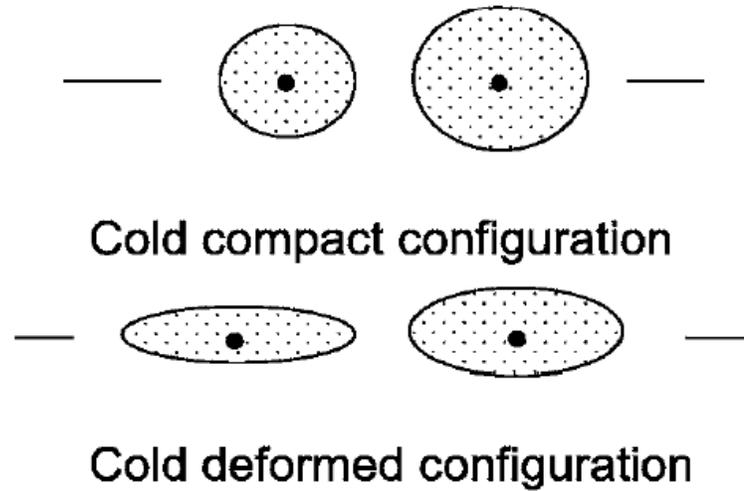
Fully controlled analysis



# Theory...



$$Q = E_{int} + V_{Coul} + V_{Def}$$



# Theory...

being maximal or minimal, respectively. This is because the **large Coulomb repulsion** in compact fission leads to high kinetic energies of the fragments while in deformed fission the smaller Coulomb repulsion entails low kinetic energies. In dedicated experiments a pronounced enhancement of **odd-even effect** in the charge yields both at the highest and the lowest kinetic energies of fragments was observed.<sup>17</sup> Since large **odd-even effects reveal low intrinsic excitation energies** the label “cold” given is justified.

$$Q = E_{int} + V_{Coul} + V_{Def}$$

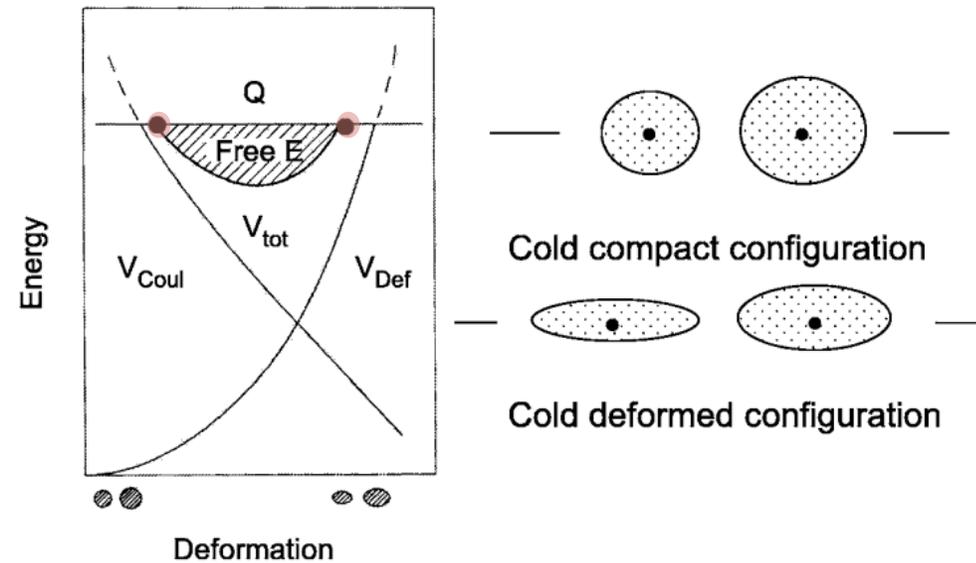
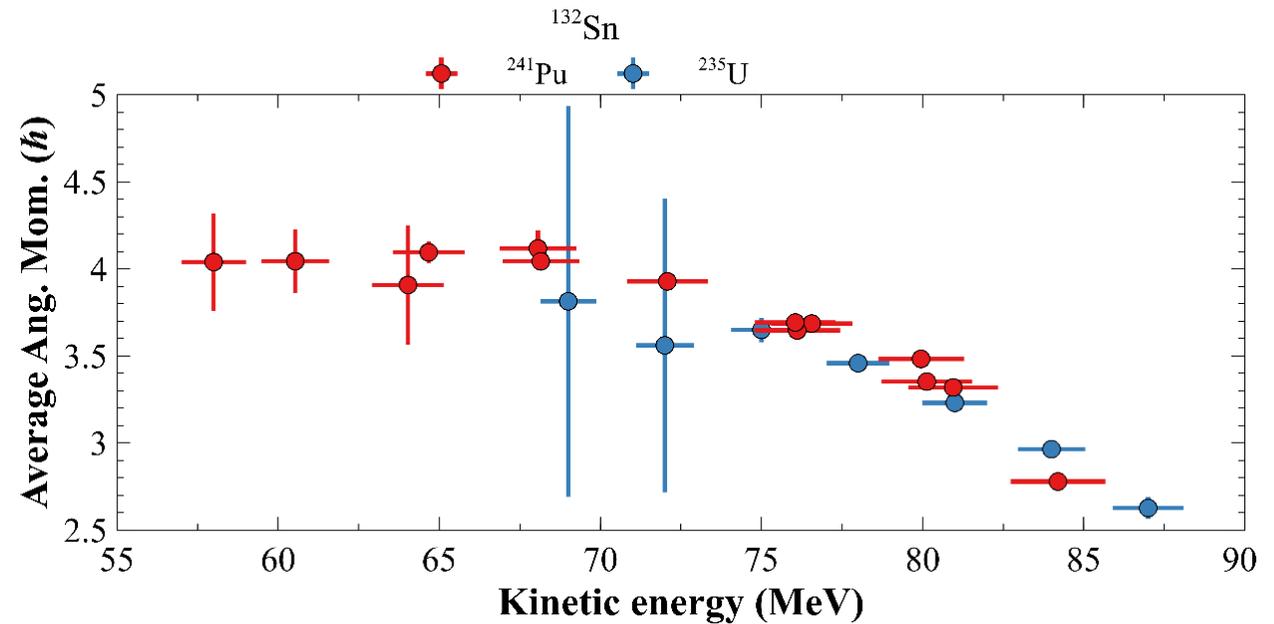
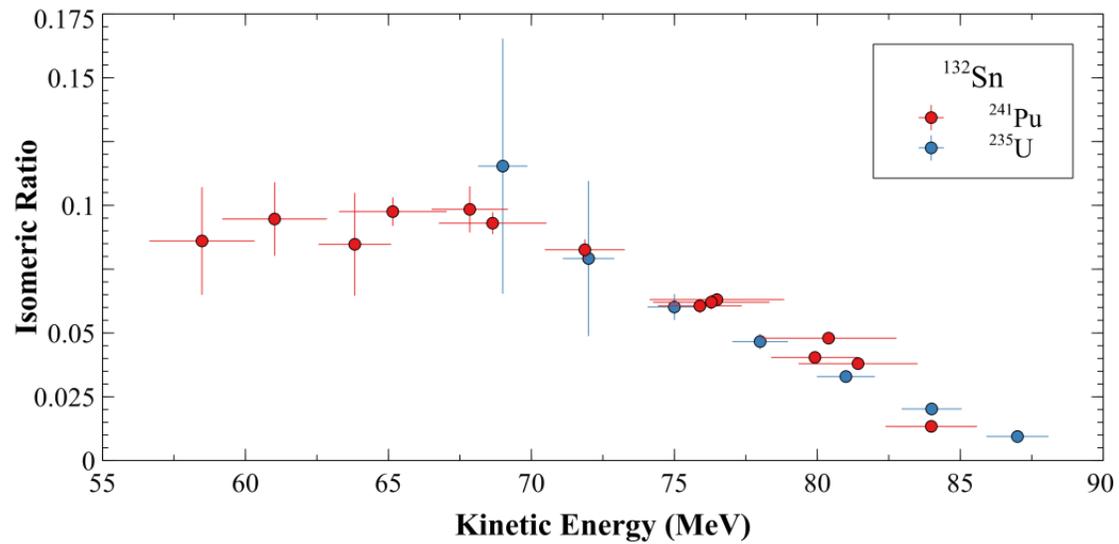


Fig. 4. Scheme of scission point model for fixed tip distance between fragments. The tip distance mimics the neck at rupture. The total potential energy  $V_{tot}$  is the sum of the Coulomb energy  $V_{Coul}$  and the deformation energy  $V_{Def}$ . The difference ( $Q - V_{tot}$ ) is called the free energy; it is available for intrinsic excitation of fragments at scission. No free energy is left at the two intersections of  $Q$  and  $V_{tot}$  as a function of fragment deformation, i.e. the fragments are cold.

# Investigation of fission product isomeric ratios and angular momenta of $^{132}\text{Sn}$ populated in the $^{241}\text{Pu}(n, f)$ reaction



# FIFRELIN: connecting experiment with theory

FIFRELIN (Fission FRagments Evaporation modeLING), used only as de-excitation code

Input parameters:  $A, Z, E^*, J^\pi$

For particular  $(E^*, J^\pi)$  combination: Models and experimental level schemes are used to create the nuclear schemes of the fission fragments

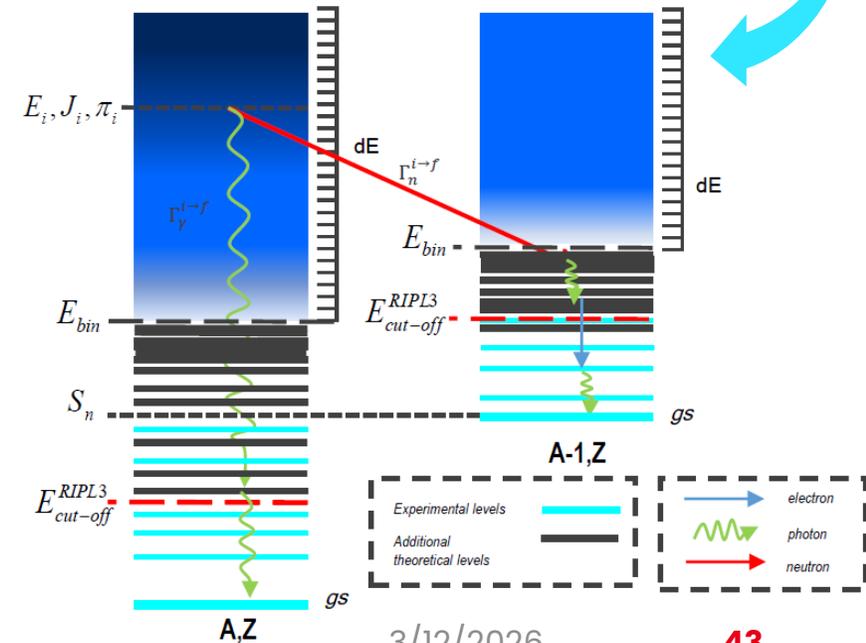
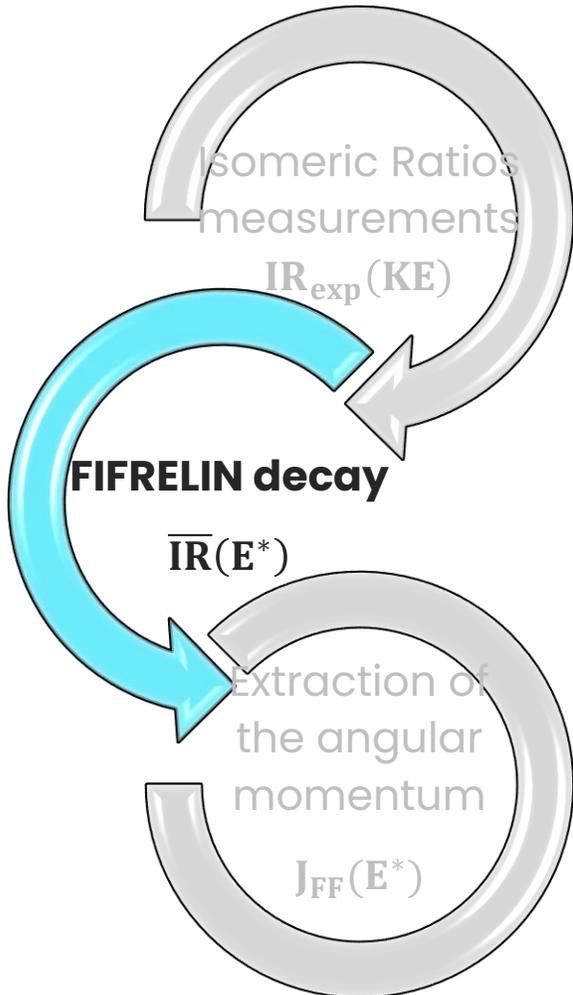
→ calculate the  $IR_{FIF}(E^*, J^\pi)$

For comparison with experimental results, they are averaged by:

$$IR_{FIF}(E^*, J_{cutoff}) = \sum_J \sum_\pi P(\pi)P(J)IR_{FIF}(E^*, J^\pi), \text{ with:}$$

- $P(J) \propto (2J + 1) \exp\left(-\frac{(J+\frac{1}{2})^2}{J_{cutoff}^2}\right)$ , spin cut-off
- $J_{cutoff}$  is a free parameter
- $P(\pi) = \frac{1}{2}$

**GOAL: extract the  $J_{cutoff}(KE)$  to reproduce experimental data**



O. Litaize, O. Serot, L. Berge, The European Physical Journal A 51, 177 (2015)

# Analysis

$$N_Y \rightarrow N_d = \frac{N_Y}{\epsilon_Y \times I_Y \times f_Y} \xrightarrow{\text{Average}} \overline{N}_d$$

$$\overline{N}_d \xrightarrow[\text{Resolution}]{\text{Bateman Equation}} \tau_f \xrightarrow[\text{ToF}]{f_{\text{gate}} \pm \sigma_{f_{\text{gate}}}}$$

$$\text{IR}_{\text{exp}} = \frac{\tau_f(\text{IS})}{\tau_f(\text{IS}) + \tau_f(\text{GS})}$$

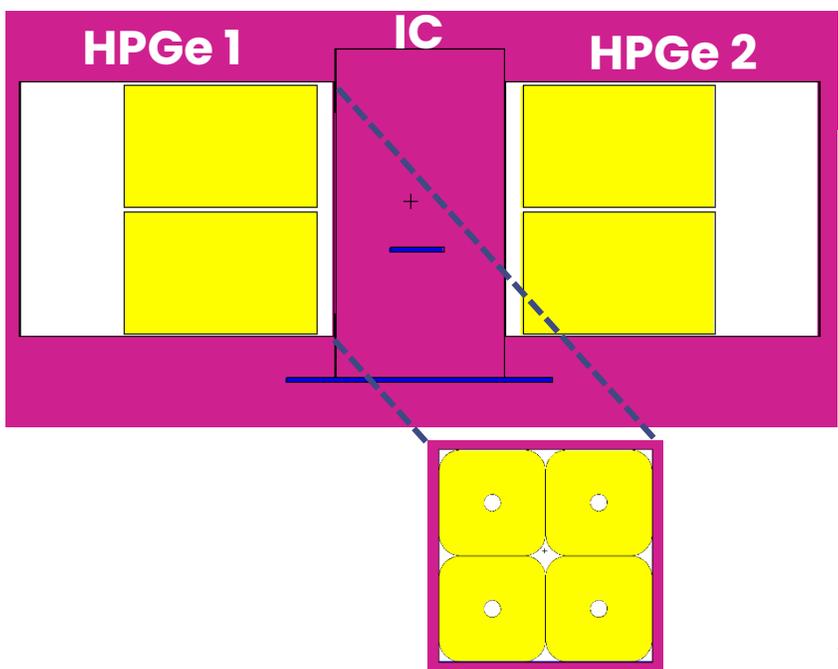
$\epsilon_Y$ : Detector Efficiency,

$f_Y$ : Correction factor for Sum-Effect,

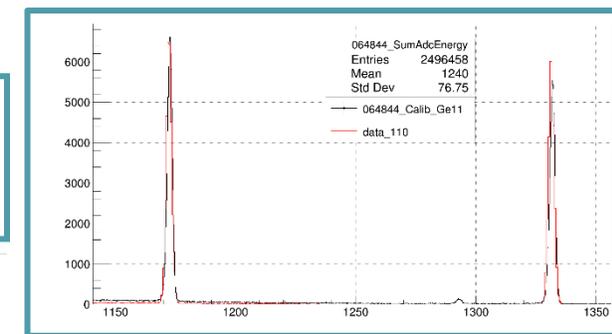
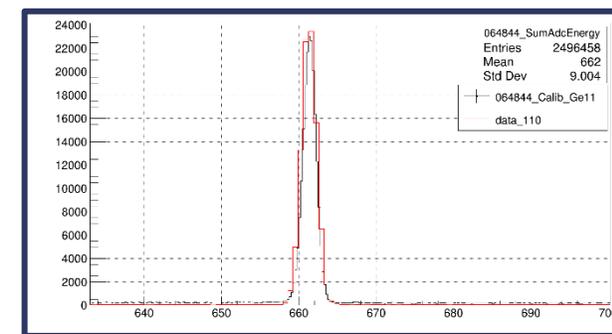
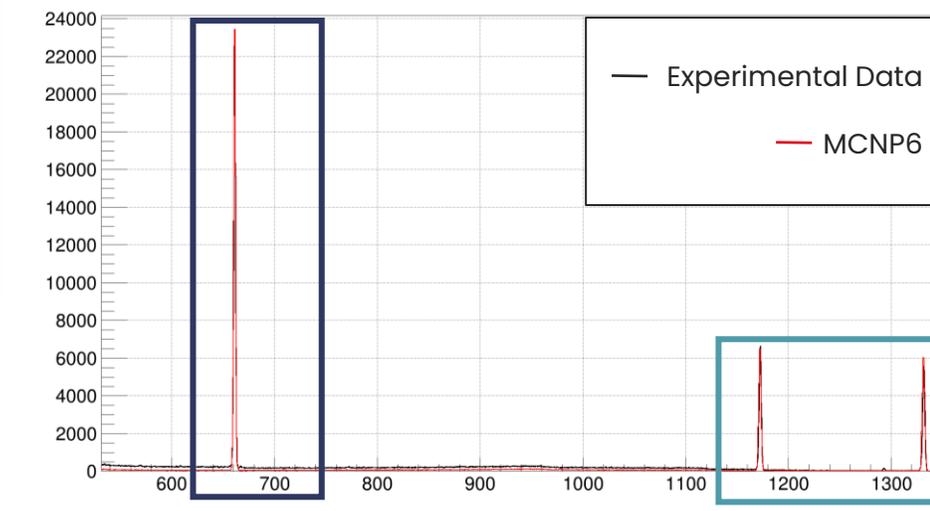
$f_{\text{gate}}$ : Correction factor for coincidence time gate,

will be derived from **TrueCoinc**<sup>1</sup> and Monte Carlo Simulation **MCNP6**<sup>2</sup> (in progress):

IR: Isomeric Ratio, IS: Isomeric State, GS: Ground State



## <sup>60</sup>Co+<sup>137</sup>Cs efficiency run



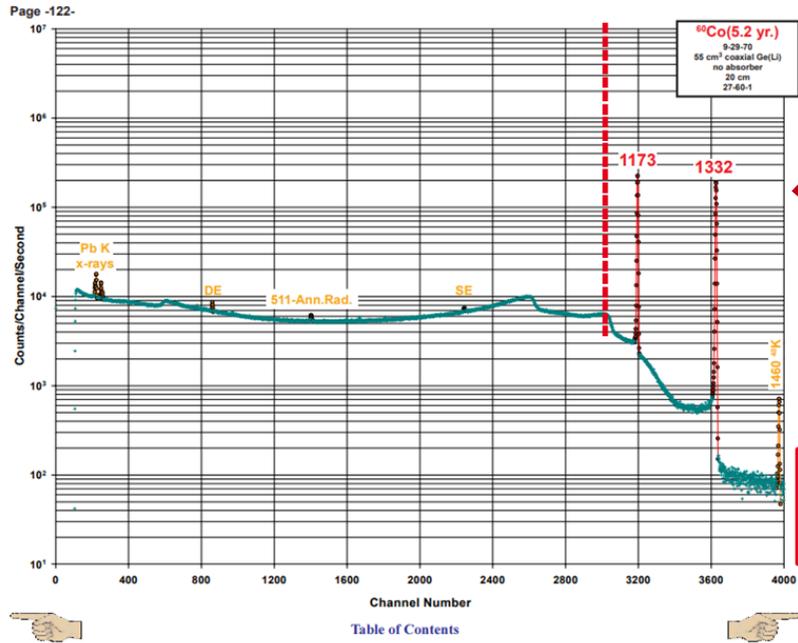
1. Sudar, S. (2002). 'TrueCoinc' software utility for calculation of the true coincidence correction (IAEA-TECDOC--1275). International Atomic Energy Agency (IAEA)

2. D.B. Pelowitz et al., Tech. Rep. LA-CP-13-00634, Los Alamos National Laboratory, Los Alamos, NM, USA, 2013

# Analysis: Sum-effect - TrueCoinc

Expected output:

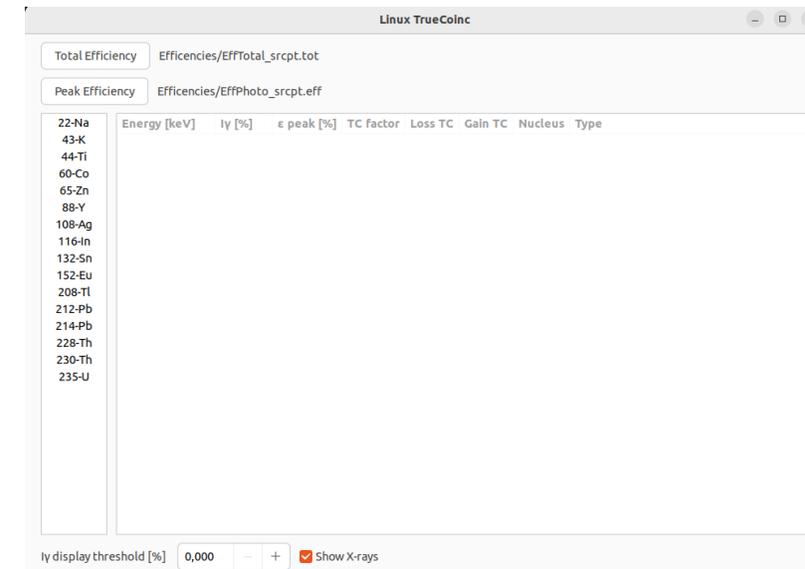
## HPGe: Co-60



<https://www.mirion.com/discover/knowledge-hub/articles/education/true-coincidence-summing>

$$\epsilon_{\gamma}^{ph} = \frac{N_{\gamma}^{ph}}{N_{emit}^{\gamma}}, \quad \epsilon_{\gamma}^{tot} = \frac{N_{\gamma}^{tot}}{N_{emit}^{\gamma}}$$

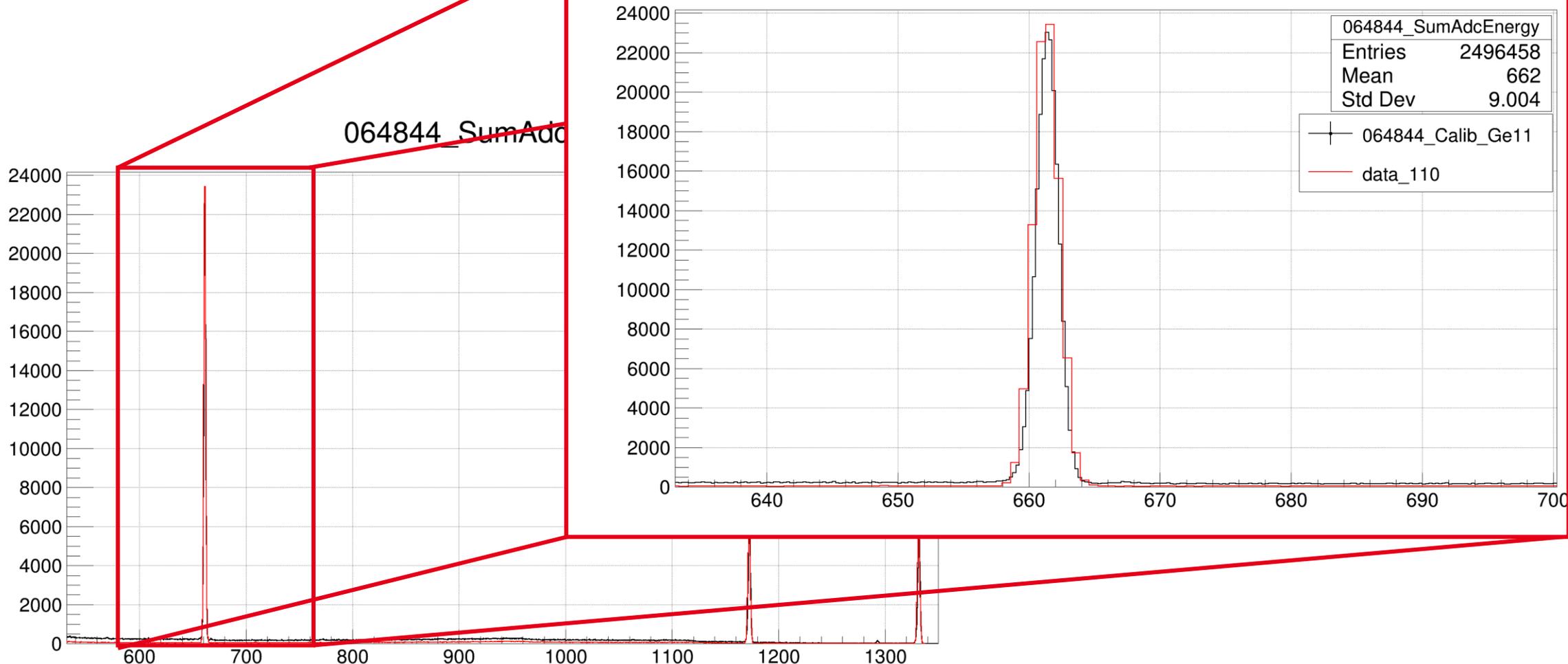
Compute the Sum-effect correction factors using **TrueCoinc**  
 Input: ENSDF files for decay



# MCNP6 - <sup>60</sup>Co + <sup>137</sup>Cs run

NPS 1e7

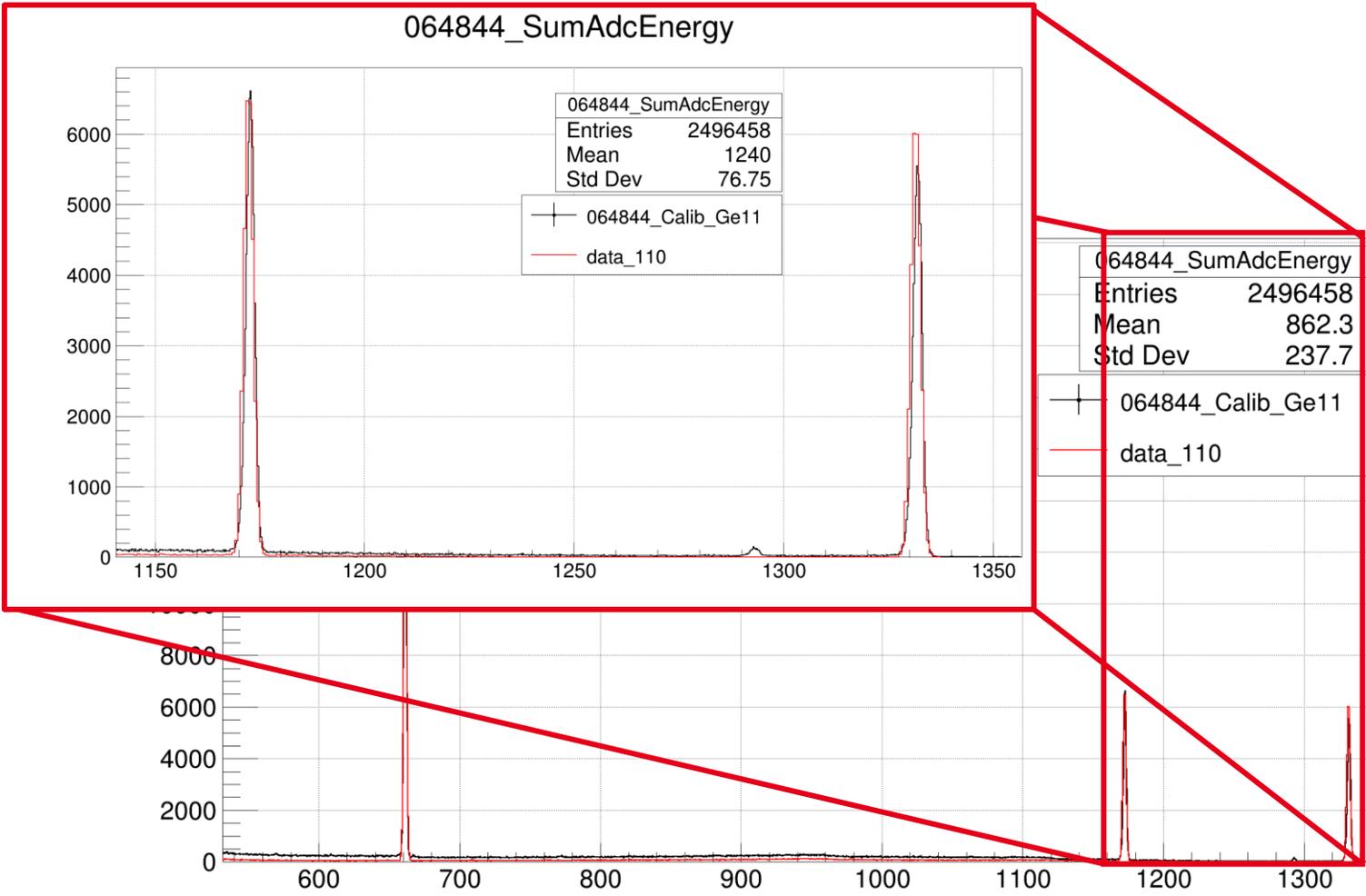
064844\_SumAdcEnergy



**points → bins: Gaussian Energy Broadening (GEB) + same binning with experimental**  
 GEB parameters based on specific crystal response

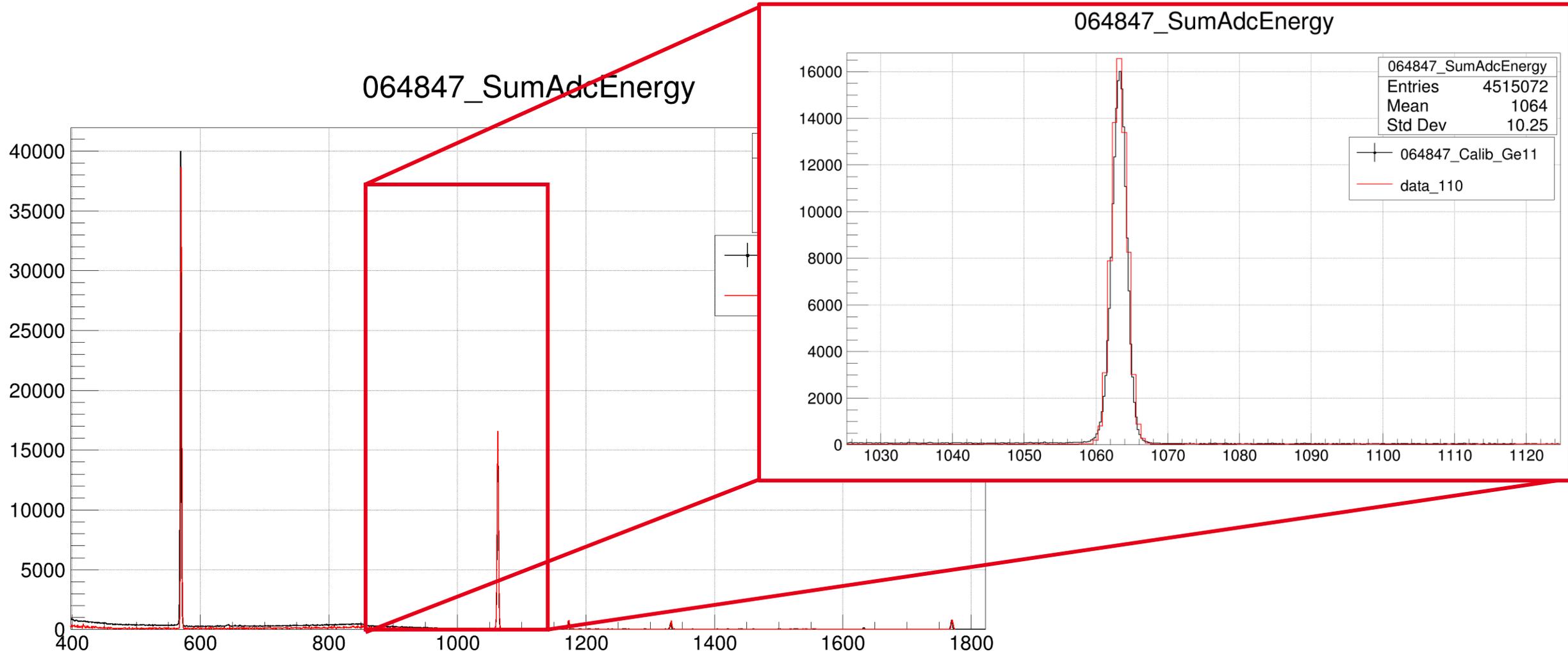
GEB a b c which defines **FWHM**:  $FWHM(E) = a + b \sqrt{E + cE^2}$ , <sup>134</sup>Te + <sup>96</sup>Y runs used

# MCNP6 – $^{60}\text{Co} + ^{137}\text{Cs}$ run



# MCNP6 – $^{60}\text{Co} + ^{207}\text{Bi}$ run

NPS 1e8

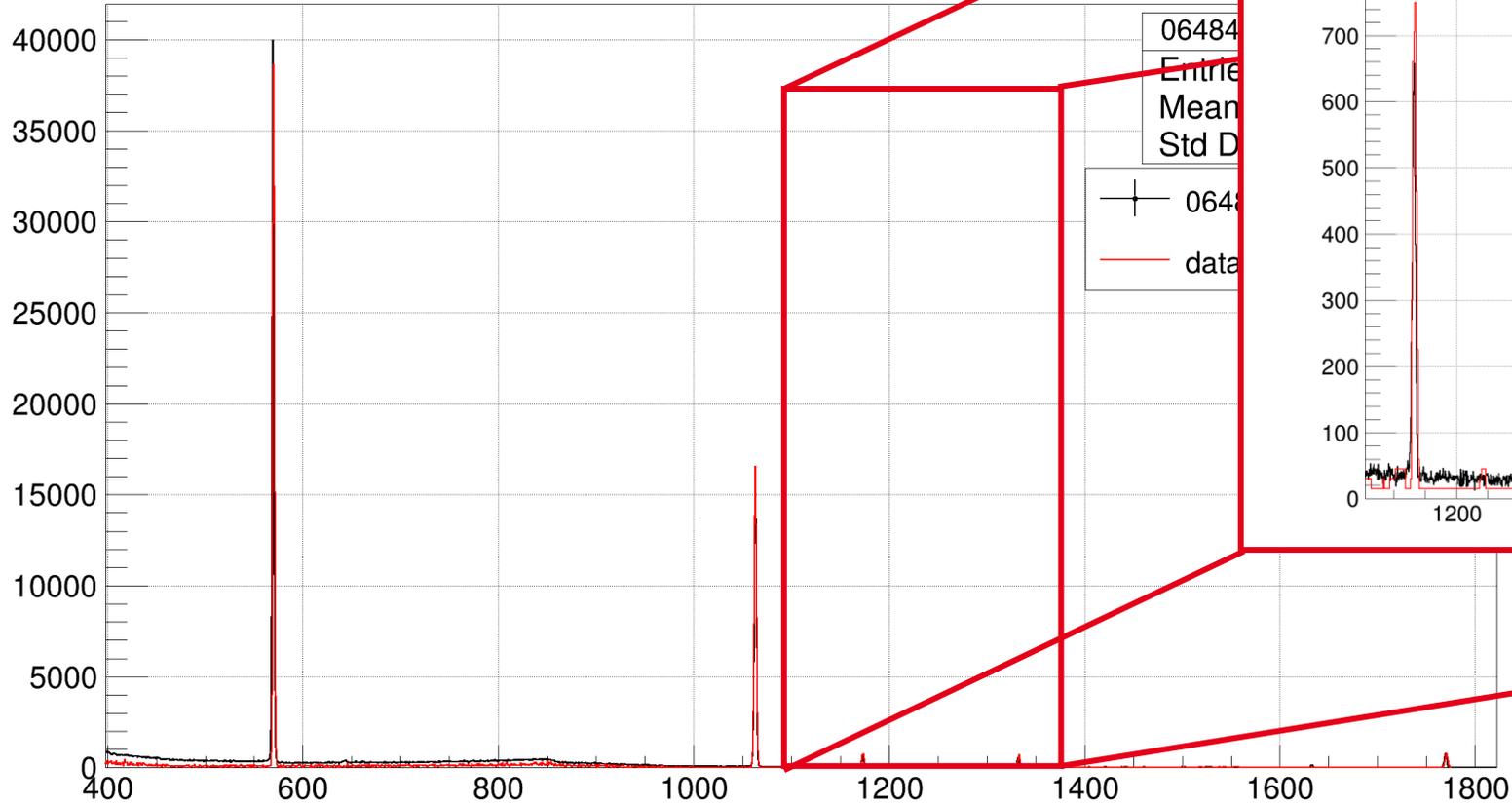


# MCNP6 – $^{60}\text{Co} + ^{207}\text{Bi}$ run

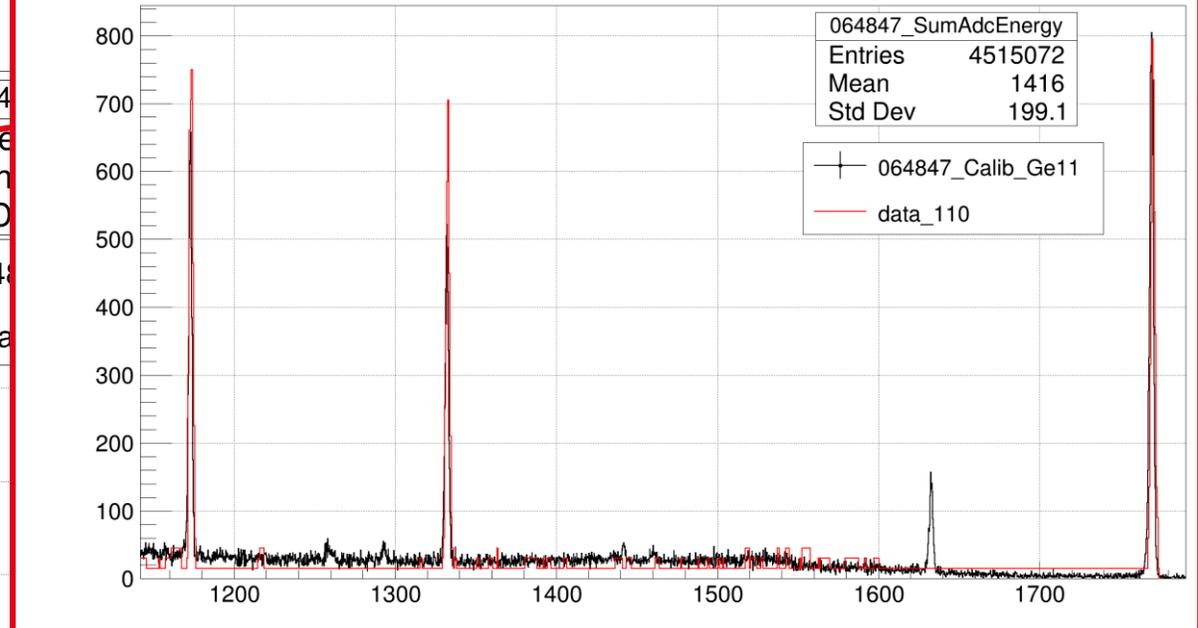
NPS 1e8



### 064847\_SumAdcEnergy



### 064847\_SumAdcEnergy



# Upcoming



| Nuclide                        | Energy [keV] | $J^\pi$ | $T_{1/2}$ or Width Abund. [mole fract.] |
|--------------------------------|--------------|---------|---|
| $^{241}_{95}\text{Am}_{146}$   | 0.0          | 5/2-    | 432.6 y 6                               |
| $^{241m1}_{95}\text{Am}_{146}$ | 2200         |         | 1.2 $\mu$ s 3                           |

| Nuclide                      | Energy [keV] | $J^\pi$   | $T_{1/2}$ or Width Abund. [mole fract.] |
|------------------------------|--------------|-----------|---|
| $^{237}_{93}\text{Np}_{144}$ | 0.0          | 5/2+      | $2.144 \times 10^6$ y 7                 |
| $^{237}_{93}\text{Np}_{144}$ | 945.20 10    | 11/2,13/2 | 0.71 $\mu$ s 4                          |

| Nuclide                        | Energy [keV]        | $J^\pi$ | $T_{1/2}$ or Width Abund. [mole fract.] |
|--------------------------------|---------------------|---------|---|
| $^{239}_{94}\text{Pu}_{145}$   | 0                   | 1/2+    | 24110 y 30                              |
| $^{239}_{94}\text{Pu}_{145}$   | 391.584 3           | 7/2-    | 193 ns 4                                |
| $^{239m1}_{94}\text{Pu}_{145}$ | $31. \times 10^2$ 2 | (5/2+)  | 7.5 $\mu$ s 10                          |

| Nuclide                        | Energy [keV] | $J^\pi$ | $T_{1/2}$ or Width Abund. [mole fract.] |
|--------------------------------|--------------|---------|---|
| $^{241}_{94}\text{Pu}_{147}$   | 0.0          | 5/2+    | 14.329 y 29                             |
| $^{241}_{94}\text{Pu}_{147}$   | 161.6853 9   | 1/2+    | 0.88 $\mu$ s 5                          |
| $^{241m1}_{94}\text{Pu}_{147}$ | 2200         |         | 20.5 $\mu$ s 22                         |

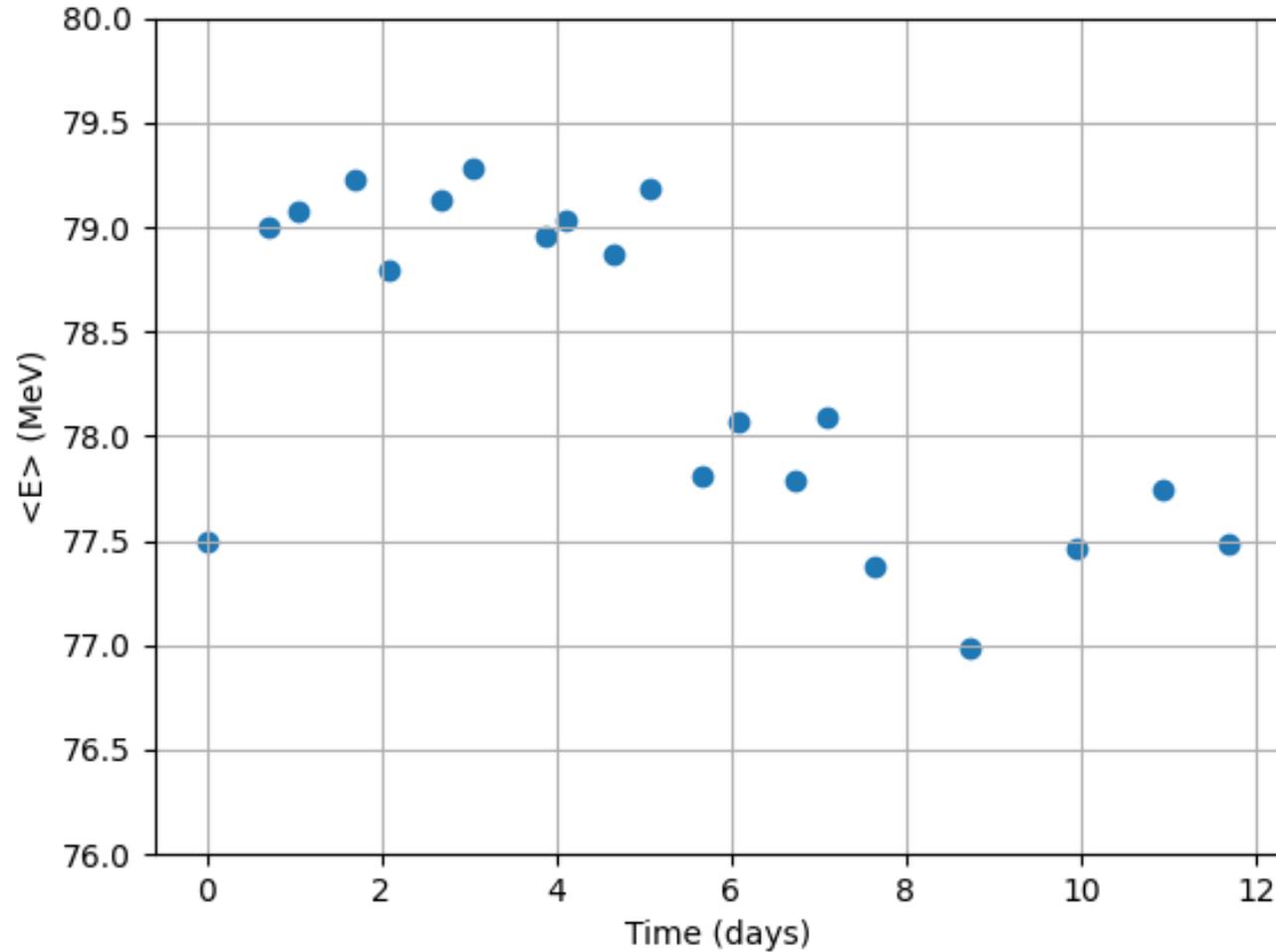
# New output:

| t_Gate ( $\mu$ s) | Coincidence Entries-Ge_5 | Duration (s) | Duration ( $\mu$ s) | Entries/t_Gate | Rate (Hz) |
|-------------------|--------------------------|--------------|---------------------|----------------|-----------|
| 12.4              | 58470                    | 1.74E-05     | 1.74E+04            | 4.72E+03       | 3.36E+09  |
| 24.8              | 96360                    | 2.98E-05     | 2.98E+04            | 3.89E+03       | 3.23E+09  |
| 26                | 99944                    | 3.10E-05     | 3.10E+04            | 3.84E+03       | 3.22E+09  |
| 28                | 105815                   | 3.30E-05     | 3.30E+04            | 3.78E+03       | 3.21E+09  |
| 29                | 108726                   | 3.40E-05     | 3.40E+04            | 3.75E+03       | 3.20E+09  |
| 30                | 111684                   | 3.50E-05     | 3.50E+04            | 3.72E+03       | 3.19E+09  |
| 31                | 114665                   | 3.60E-05     | 3.60E+04            | 3.70E+03       | 3.19E+09  |
| 32                | 117598                   | 3.70E-05     | 3.70E+04            | 3.67E+03       | 3.18E+09  |
| 34                | 123232                   | 3.90E-05     | 3.90E+04            | 3.62E+03       | 3.16E+09  |
| 35                | 126204                   | 4.00E-05     | 4.00E+04            | 3.61E+03       | 3.16E+09  |
| 36                | 128944                   | 4.10E-05     | 4.10E+04            | 3.58E+03       | 3.14E+09  |
| 38                | 134597                   | 4.30E-05     | 4.30E+04            | 3.54E+03       | 3.13E+09  |
| 39                | 137503                   | 4.40E-05     | 4.40E+04            | 3.53E+03       | 3.13E+09  |
| 40                | 140332                   | 4.50E-05     | 4.50E+04            | 3.51E+03       | 3.12E+09  |
| 41                | 143035                   | 4.60E-05     | 4.60E+04            | 3.49E+03       | 3.11E+09  |
| 42                | 145812                   | 4.70E-05     | 4.70E+04            | 3.47E+03       | 3.10E+09  |
| 43                | 148626                   | 4.80E-05     | 4.80E+04            | 3.46E+03       | 3.10E+09  |
| 44                | 151395                   | 4.90E-05     | 4.90E+04            | 3.44E+03       | 3.09E+09  |
| 45                | 154118                   | 5.00E-05     | 5.00E+04            | 3.42E+03       | 3.08E+09  |
| 46                | 156883                   | 5.10E-05     | 5.10E+04            | 3.41E+03       | 3.08E+09  |
| 47                | 159625                   | 5.20E-05     | 5.20E+04            | 3.40E+03       | 3.07E+09  |
| 48                | 162326                   | 5.30E-05     | 5.30E+04            | 3.38E+03       | 3.06E+09  |
| 49                | 165048                   | 5.40E-05     | 5.40E+04            | 3.37E+03       | 3.06E+09  |
| 50                | 167823                   | 5.50E-05     | 5.50E+04            | 3.36E+03       | 3.05E+09  |
| 52                | 173226                   | 5.70E-05     | 5.70E+04            | 3.33E+03       | 3.04E+09  |
| 54                | 178638                   | 5.90E-05     | 5.90E+04            | 3.31E+03       | 3.03E+09  |
| 56                | 184156                   | 6.10E-05     | 6.10E+04            | 3.29E+03       | 3.02E+09  |
| 58                | 189722                   | 6.30E-05     | 6.30E+04            | 3.27E+03       | 3.01E+09  |
| 60                | 195126                   | 6.50E-05     | 6.50E+04            | 3.25E+03       | 3.00E+09  |
| 62                | 200454                   | 6.70E-05     | 6.70E+04            | 3.23E+03       | 2.99E+09  |

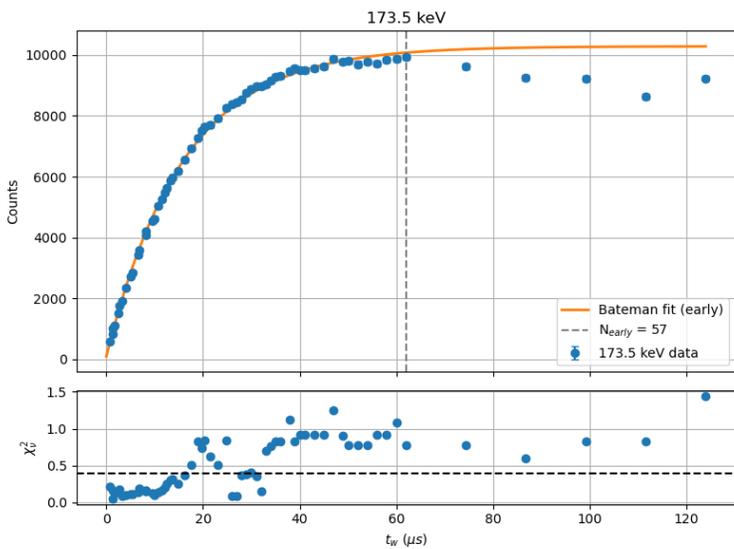
x3

x4

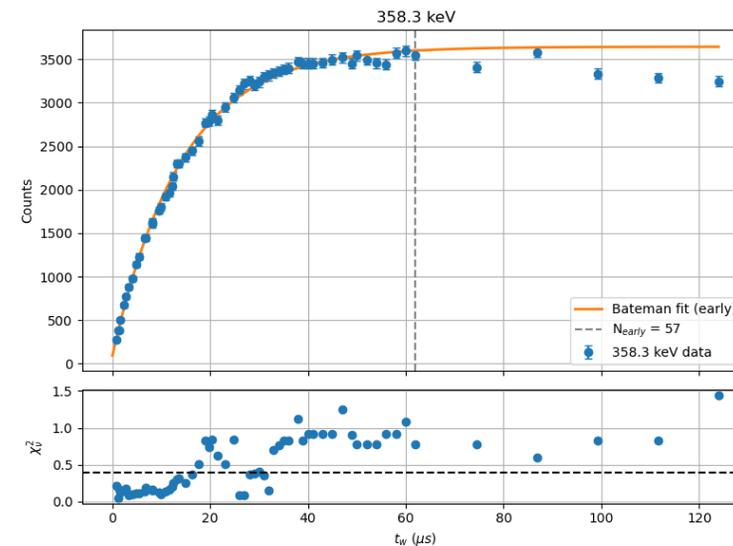
# Target evolution



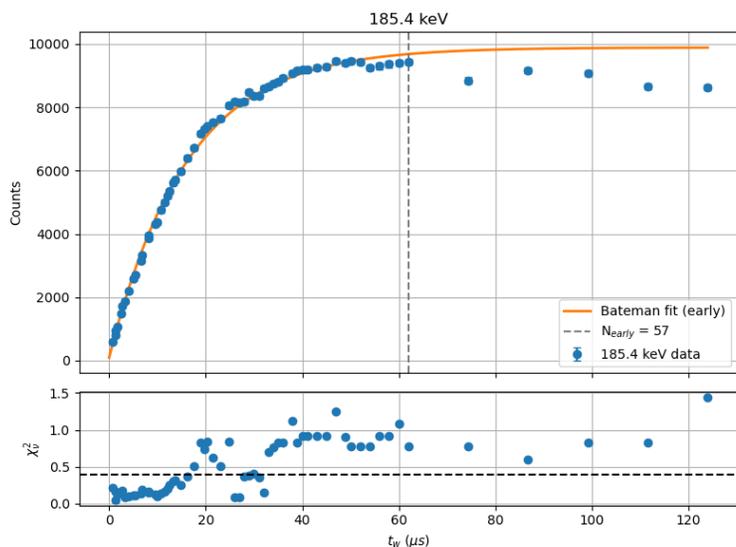
# Analysis - $\Delta t_{\text{Gate}} = {}^{100}\text{Nb}$



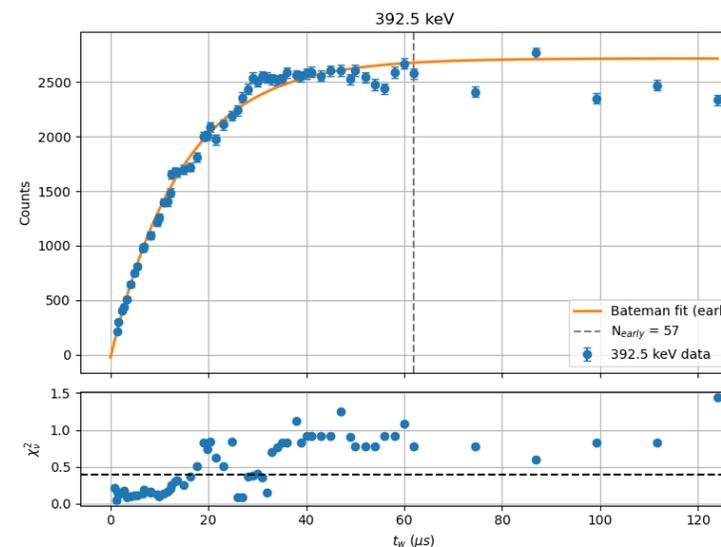
$t_{1/2} = 11.06 \pm 0.10 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 2.14$   
 p-value = 2.145e-06



$t_{1/2} = 9.83 \pm 0.15 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 1.31$   
 p-value = 6.427e-02

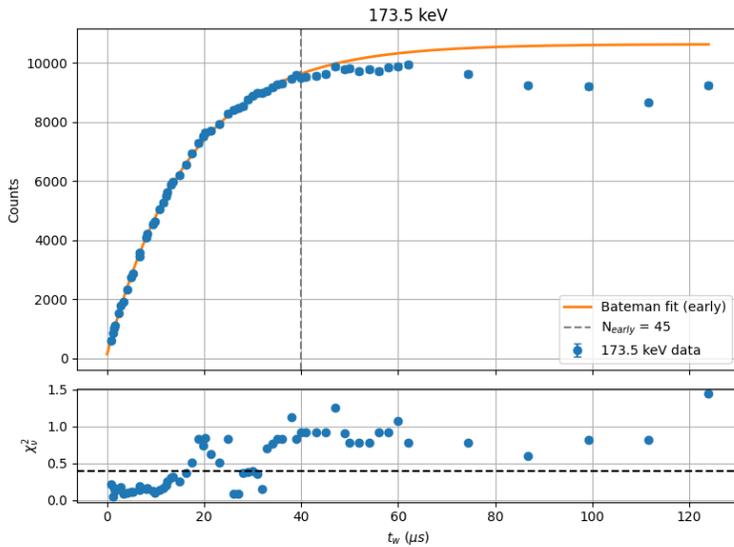


$t_{1/2} = 11.15 \pm 0.10 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 3.91$   
 p-value = 2.987e-16

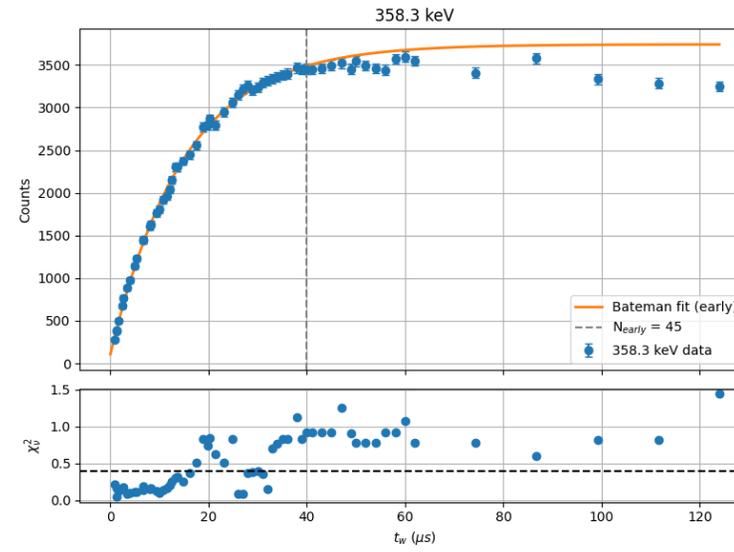


$t_{1/2} = 10.09 \pm 0.18 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 3.34$   
 p-value = 4.996e-15

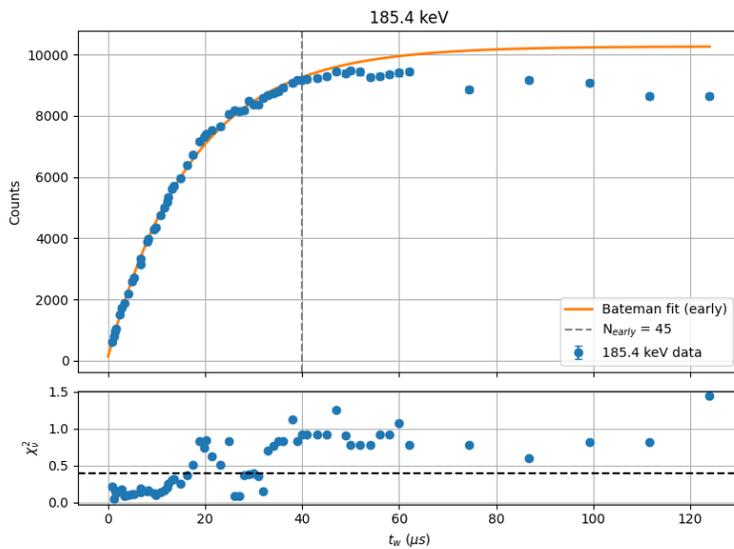
# Analysis - $\Delta t_{\text{Gate}} = {}^{100}\text{Nb}$



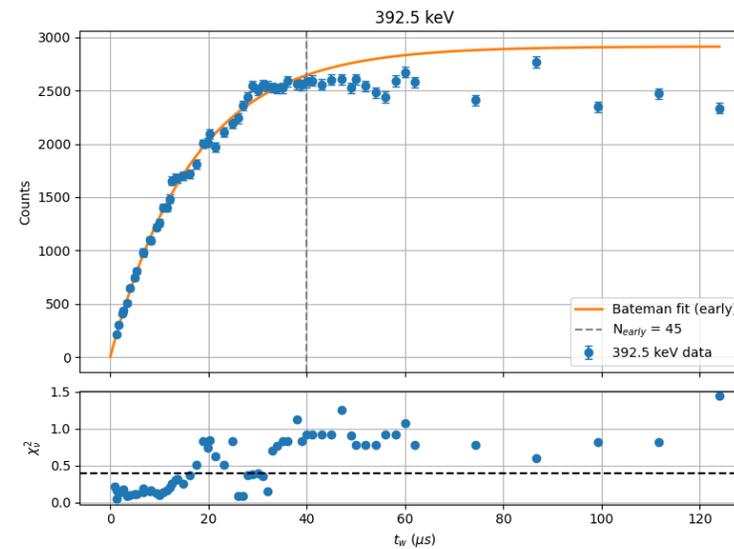
$t_{1/2} = 11.79 \pm 0.16 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 1.35$   
 p-value = 6.427e-02



$t_{1/2} = 10.39 \pm 0.22 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 1.16$   
 p-value = 2.163e-01



$t_{1/2} = 12.00 \pm 0.16 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 3.06$   
 p-value = 9.947e-11



$t_{1/2} = 11.64 \pm 0.31 \mu\text{s}$   
 $\chi^2 / \text{ndf} = 2.11$   
 p-value = 5.271e-5

VS

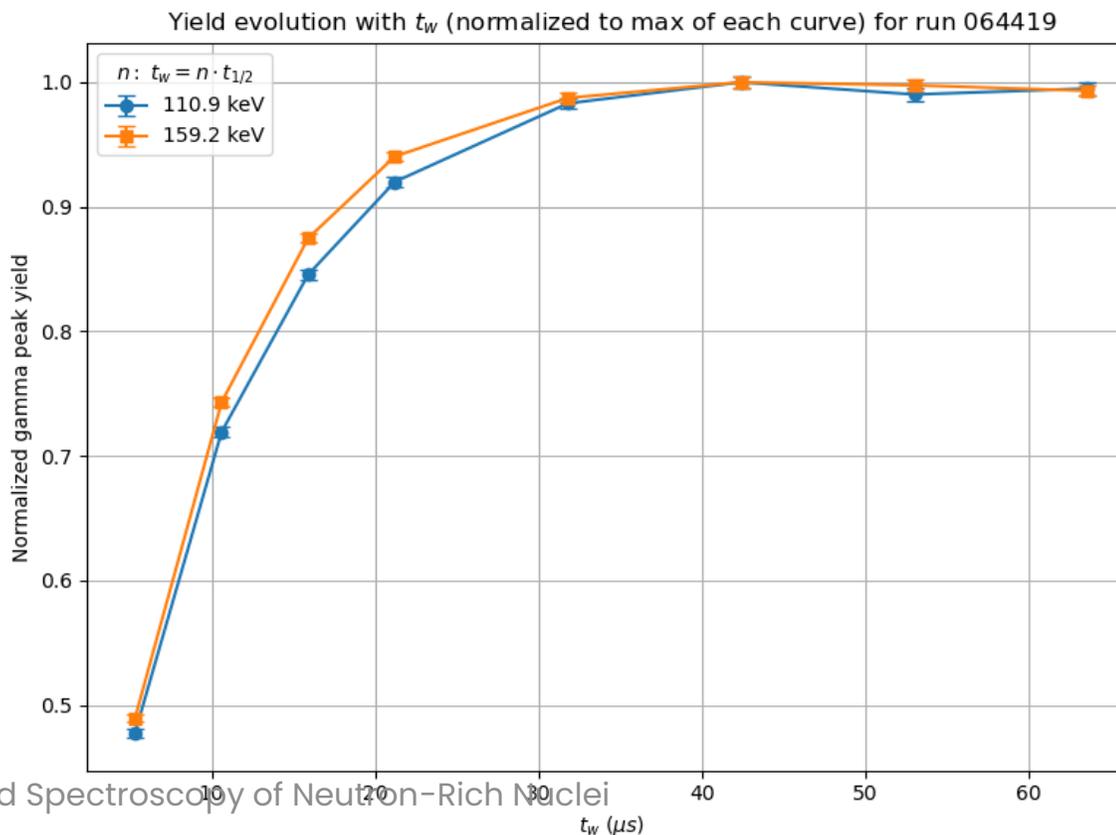
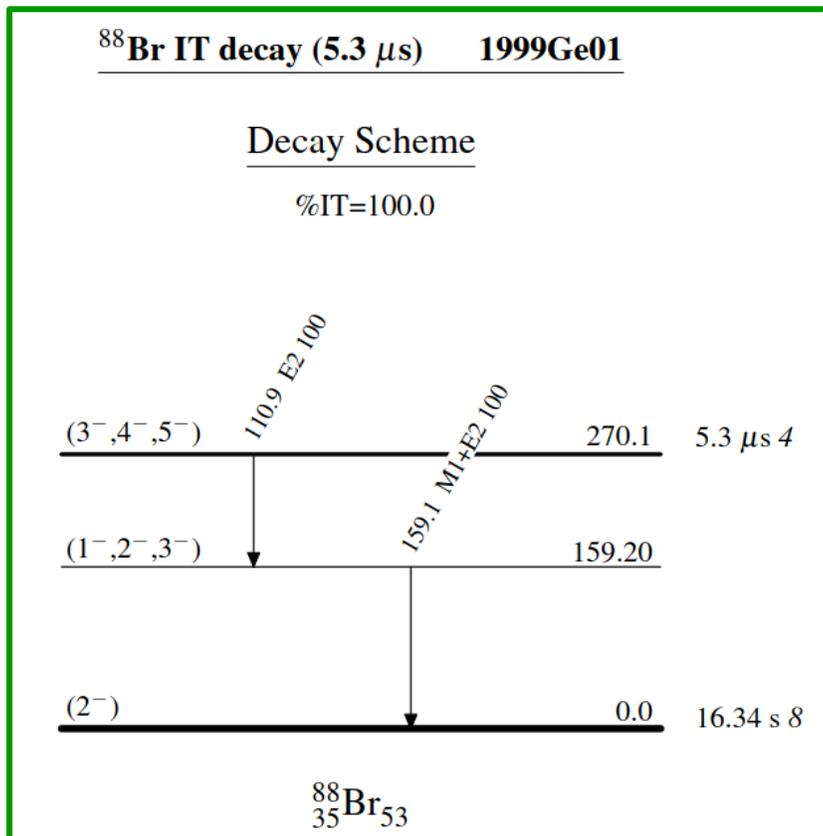
Fit experimental data with function  $f = A + B * (1 - e^{-\lambda t})$ :

110.9 keV:  $\lambda = 1.2019\text{e-}01 \pm 1.0444\text{e-}03$   **$t_{1/2} = 5.77 \pm 0.05$**

Quality of the fit:  $\chi^2 / \text{ndf} = 1.71$ , p-value =  $1.5\text{e-}01$

159.2 keV:  $\lambda = 1.2877\text{e-}01 \pm 1.0513\text{e-}03$   **$t_{1/2} = 5.38 \pm 0.04$**  😊

Quality of the fit:  $\chi^2 / \text{ndf} = 1.81$ , p-value =  $1.2\text{e-}01$





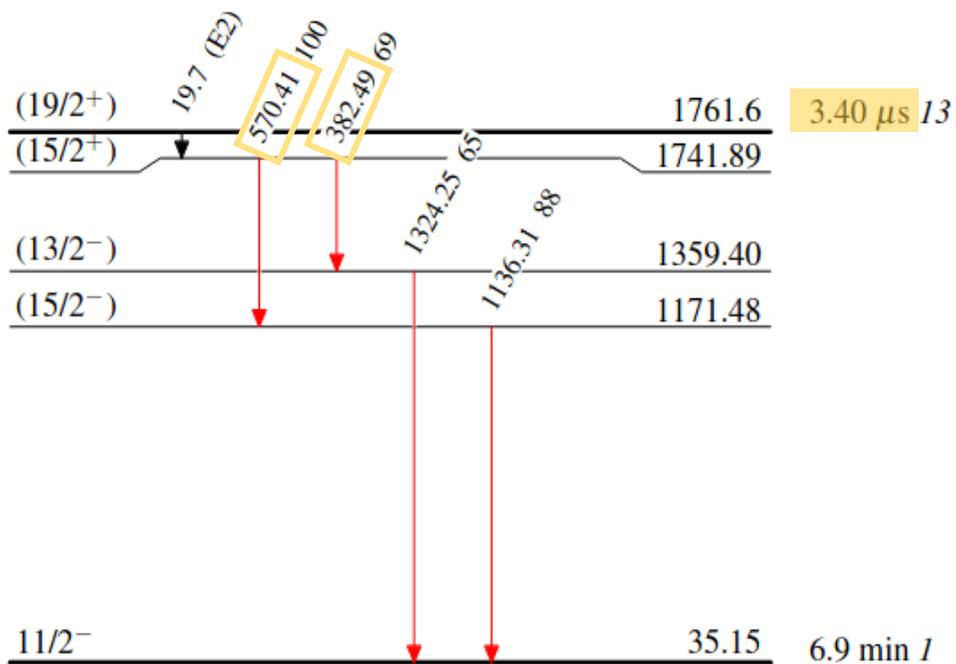
$^{129}\text{Sn}$  IT decay (3.40  $\mu\text{s}$ ) 2002Ge07,2008Lo07

Decay Scheme

Intensities: Relative  $I_\gamma$   
%IT=100.0

Legend

- $I_\gamma < 2\% \times I_\gamma^{\text{max}}$
- $I_\gamma < 10\% \times I_\gamma^{\text{max}}$
- $I_\gamma > 10\% \times I_\gamma^{\text{max}}$



$$t_{1/2} = 3.4 \pm 0.2 \mu\text{s (IAEA)}$$

VS

Fit experimental data with function  $f = A + B * (1 - e^{-\lambda t})$ :

382.5 keV:  $\lambda = 2.1553e-01 \pm 3.5996e-03$   $t_{1/2} = 3.22 \pm 0.05$

Quality of the fit:  $\chi^2 / \text{ndf} = 1.34$ , p-value =  $2.5e-01$

570.4 keV:  $\lambda = 2.1413e-01 \pm 8.4292e-03$   $t_{1/2} = 3.24 \pm 0.13$

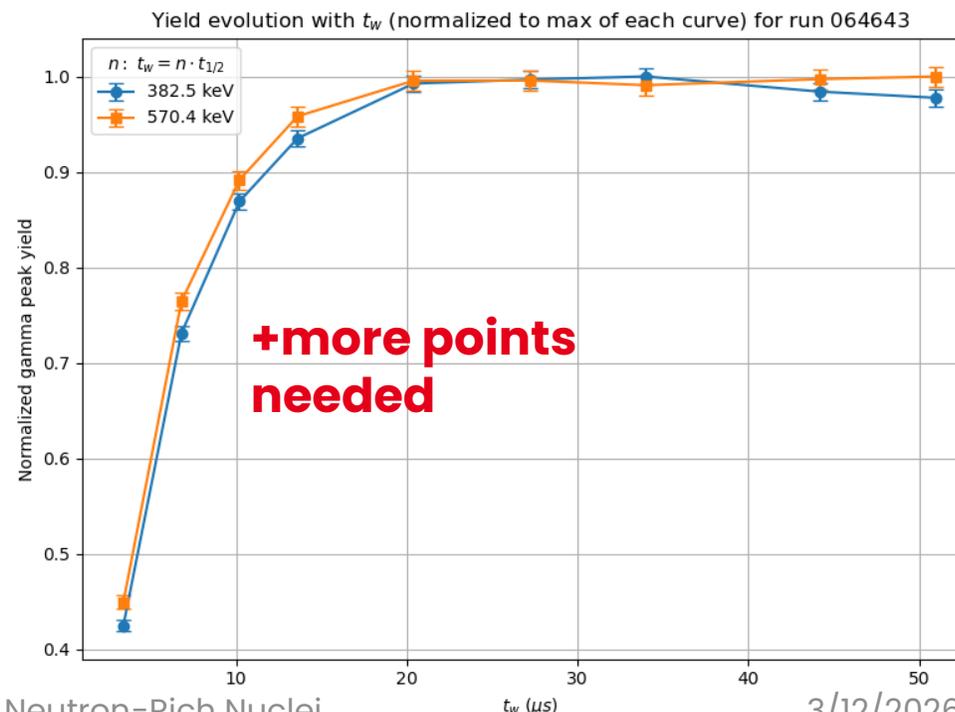
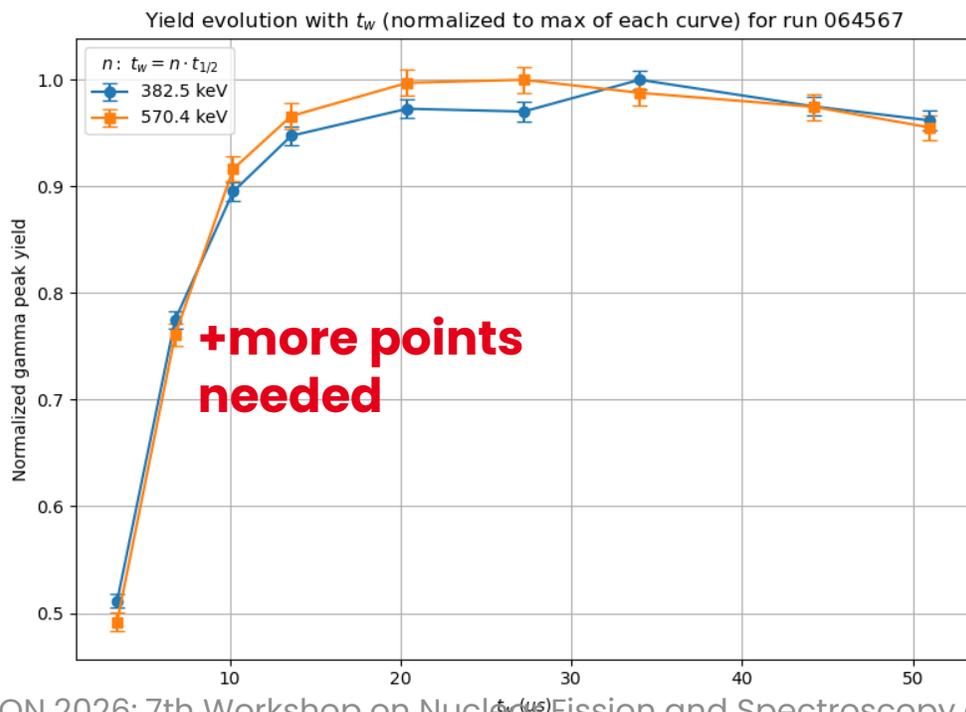
Quality of the fit:  $\chi^2 / \text{ndf} = 4.45$ , p-value =  $1.3e-03$

KE = 62

382.5 keV:  $\lambda = 1.8592e-01 \pm 8.2569e-03$   $t_{1/2} = 3.73 \pm 0.17$   
Quality of the fit:  $\chi^2 / \text{ndf} = 16.75$ , p-value =  $9.7e-14$

570.4 keV:  $\lambda = 2.0176e-01 \pm 1.0225e-02$   $t_{1/2} = 3.44 \pm 0.17$   
Quality of the fit:  $\chi^2 / \text{ndf} = 14.14$ , p-value =  $1.5e-11$

KE = 74



Correct central value but not always good fit

# $^{131}\text{Sb}$ – Study Case



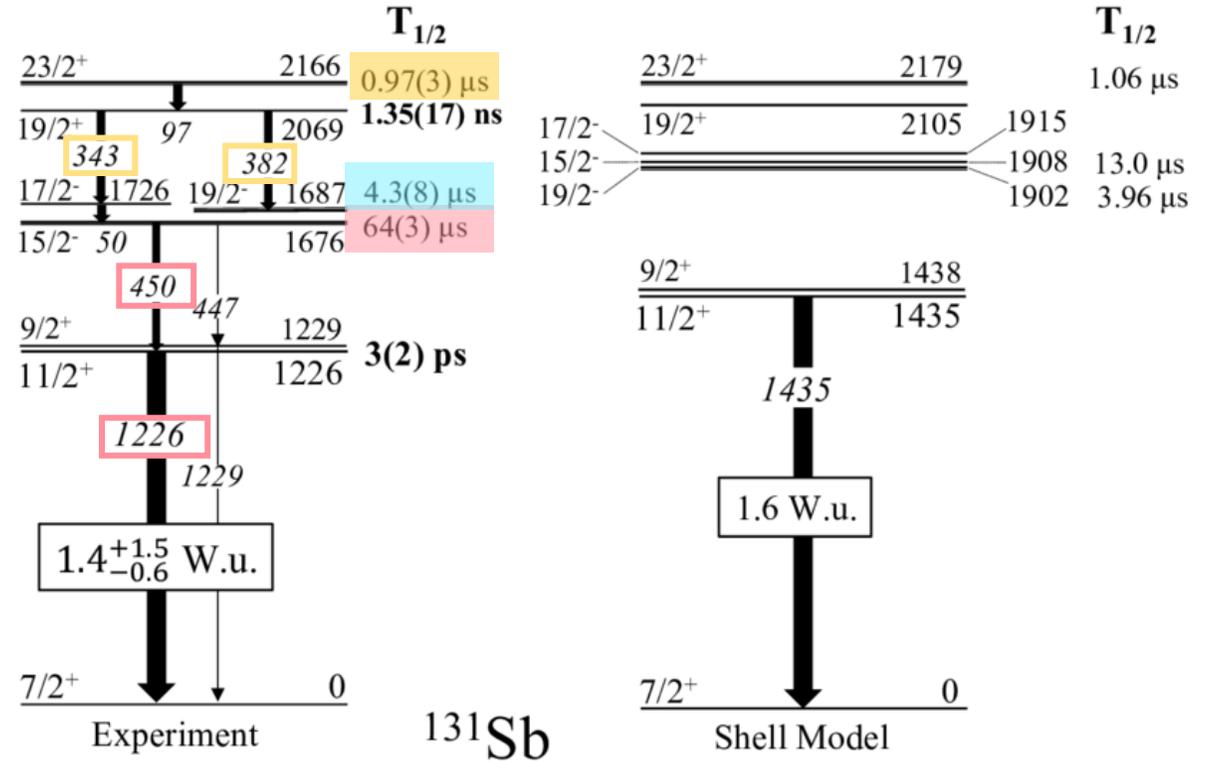
IDENTIFICATION CARD

$^{131}\text{Sb}$

Ground State:  $t_{1/2} = 23.03$  min

Metastable States: (IAEA, NNDS)

1.  $E_{\text{level}} = 2165.6$  keV,  $t_{1/2} = 1.1$   $\mu\text{s}$
2.  $E_{\text{level}} = 1687.2$  keV,  $t_{1/2} = 4.3$   $\mu\text{s}$
3.  $E_{\text{level}} = 1676.1$  keV,  $t_{1/2} = 91$   $\mu\text{s}$



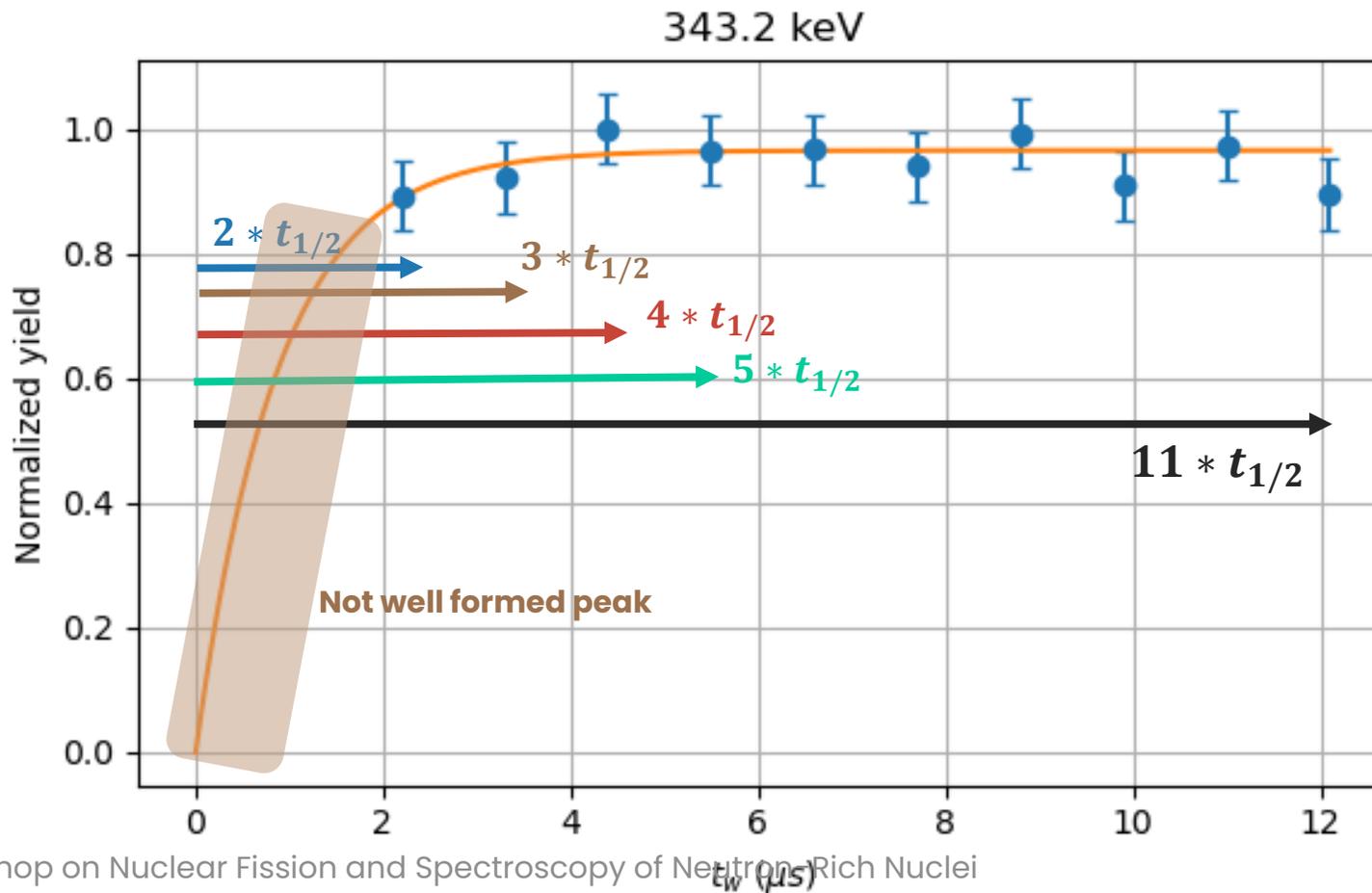
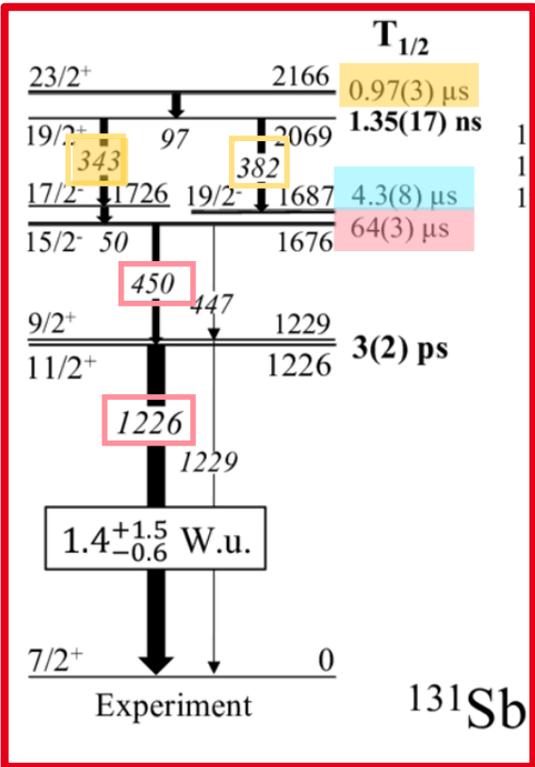
Lifetime measurements in  $^{131}\text{Sb}$  at LOHENGRIN  
S. Bottoni et al



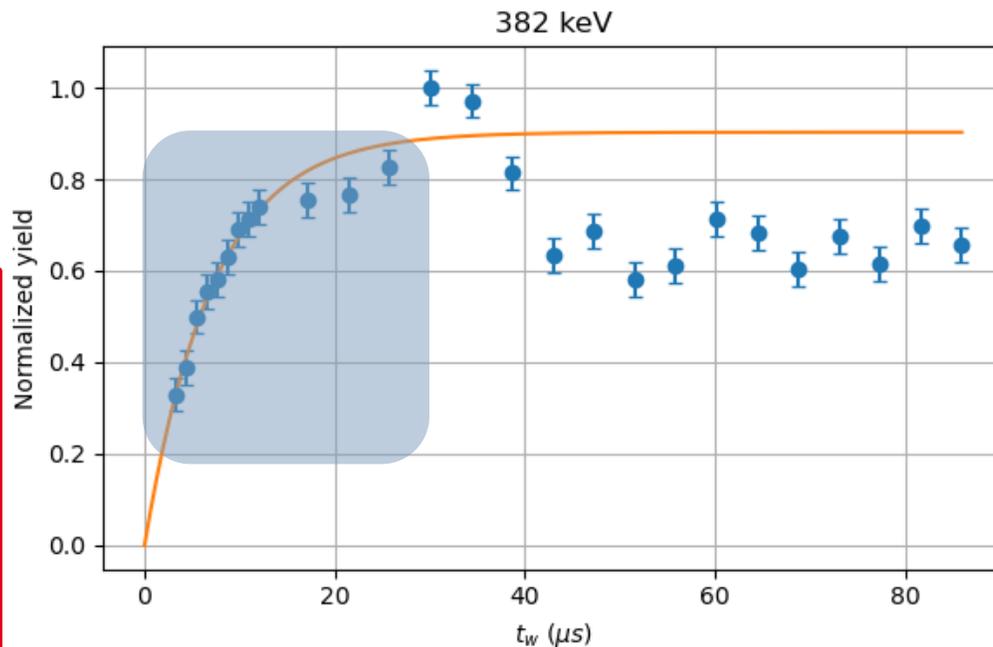
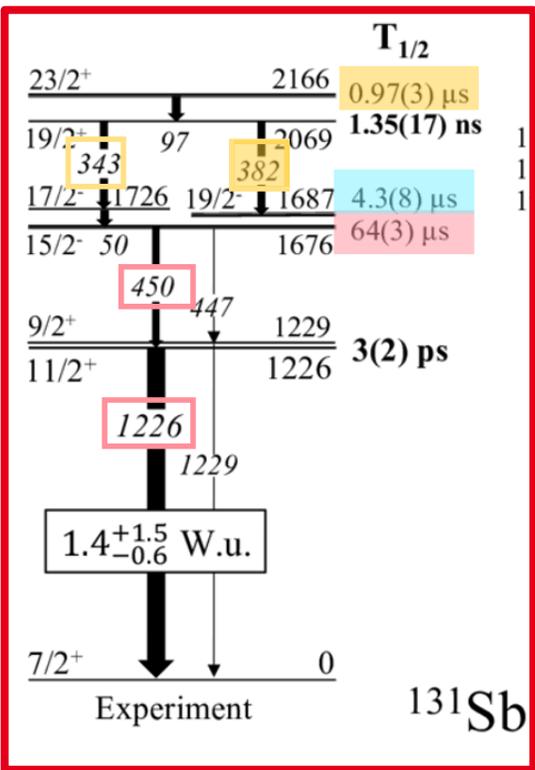
Python Script Weighted Fit:  
 $\lambda = 1.1524e+00 \pm 2.7269 \text{ 1}/\mu\text{s}$

**$t_{1/2} = 0.60 \pm 1.42 \mu\text{s}$**

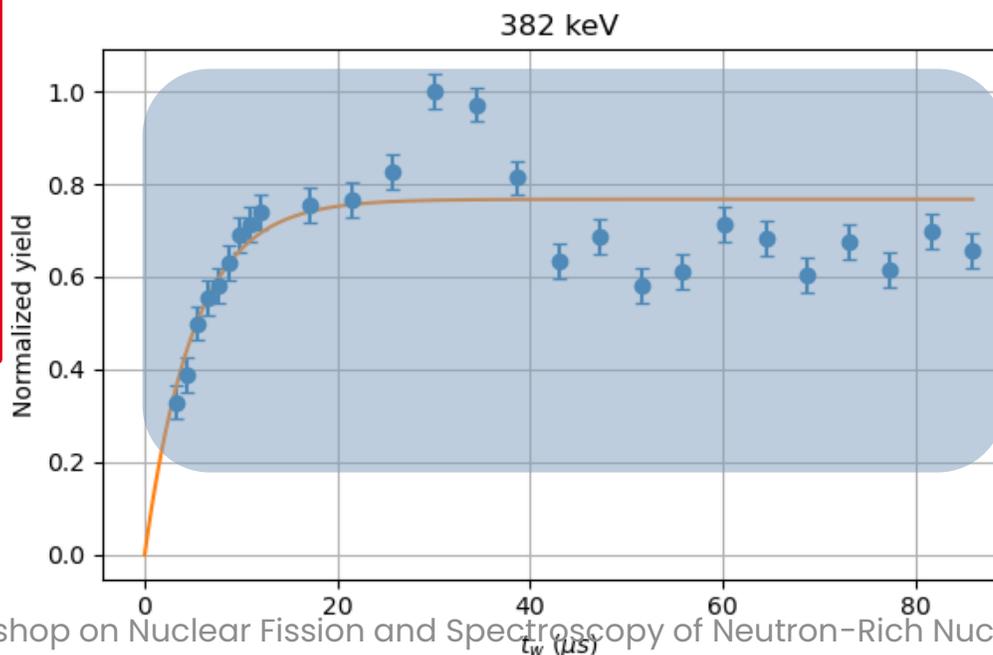
Quality of the fit:  $\chi^2 / \text{ndf} = 0.35, \text{p-value} = 9.1e-01$



# <sup>131</sup>Sb

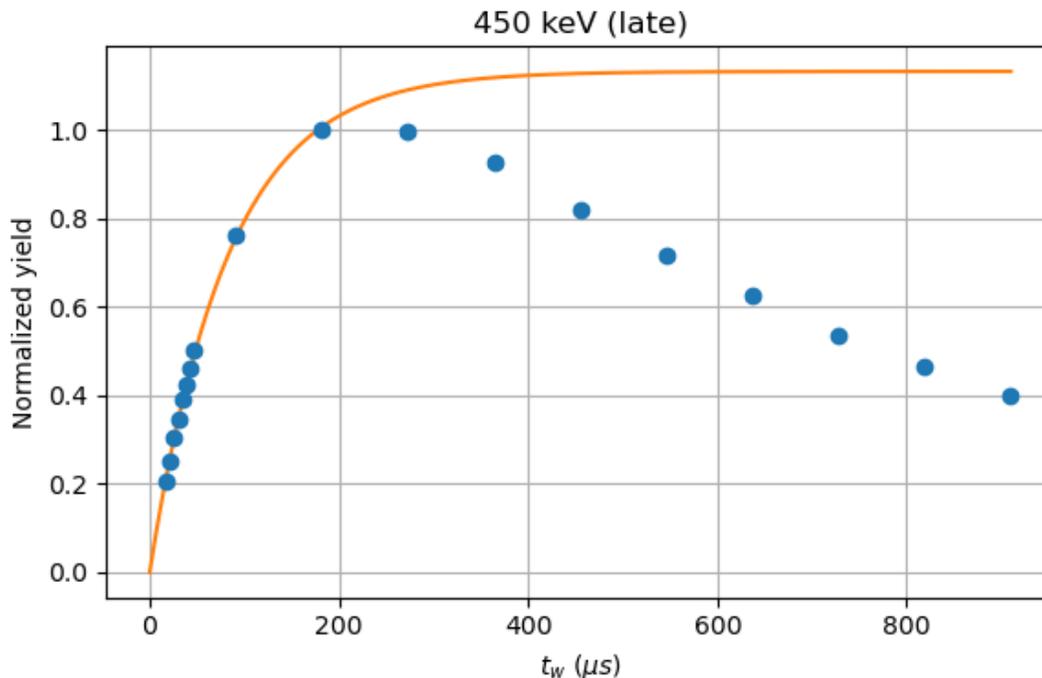
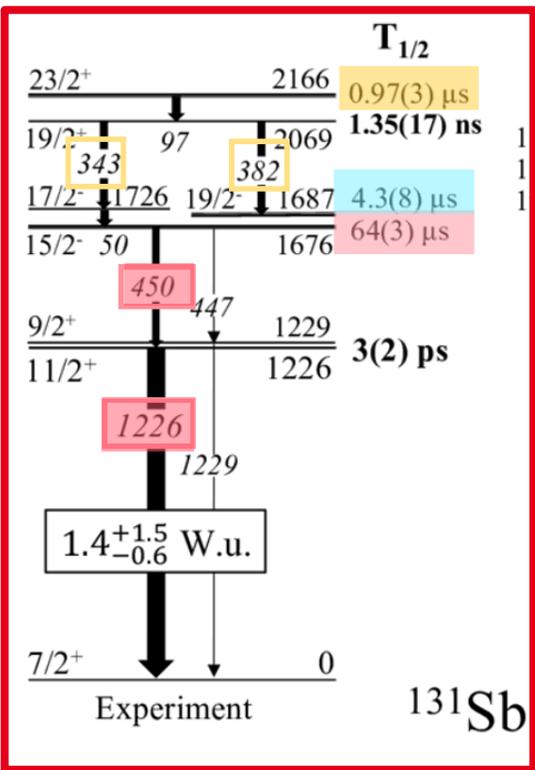


12 first points  
 Python Script Weighted Fit:  
 $\lambda = 1.6562e-01 \pm 3.2601e-02$  1/μs  
 $t_{1/2} = 4.19 \pm 0.82$  μs  
 Quality of the fit:  
 $\chi^2 / \text{ndf} = 0.46$ , p-value = 9.1e-01

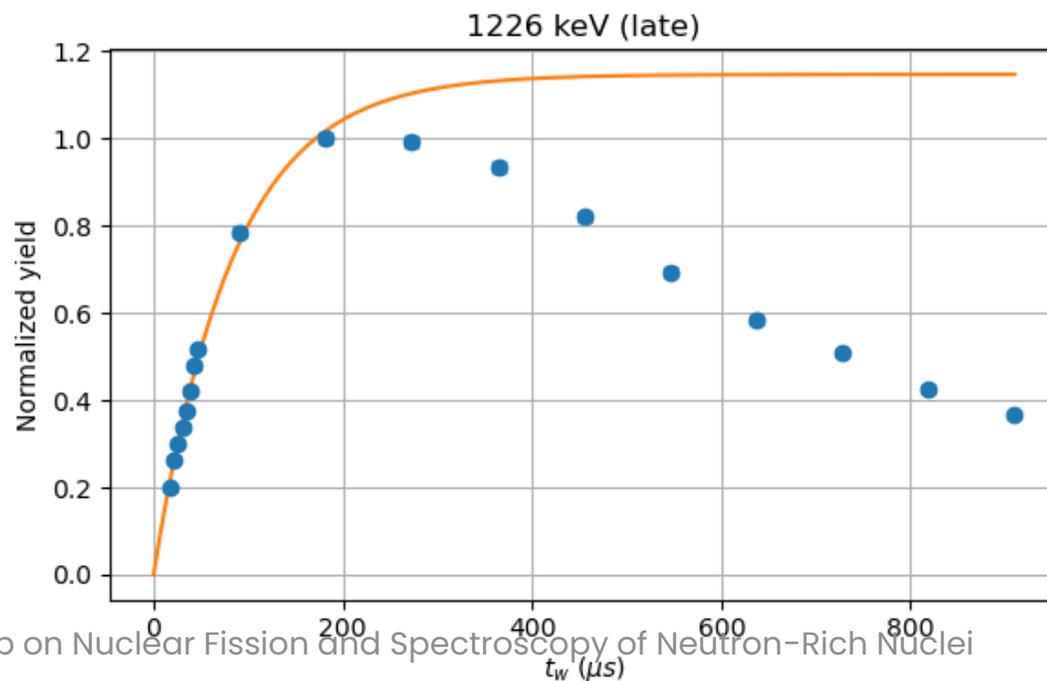


All points  
 Python Script Weighted Fit:  
 $\lambda = 1.9555e-01 \pm 2.9853e-02$  1/μs  
 $t_{1/2} = 3.54 \pm 0.54$  μs  
 Quality of the fit:  
 $\chi^2 / \text{ndf} = 7.88$ , p-value = 1e-015

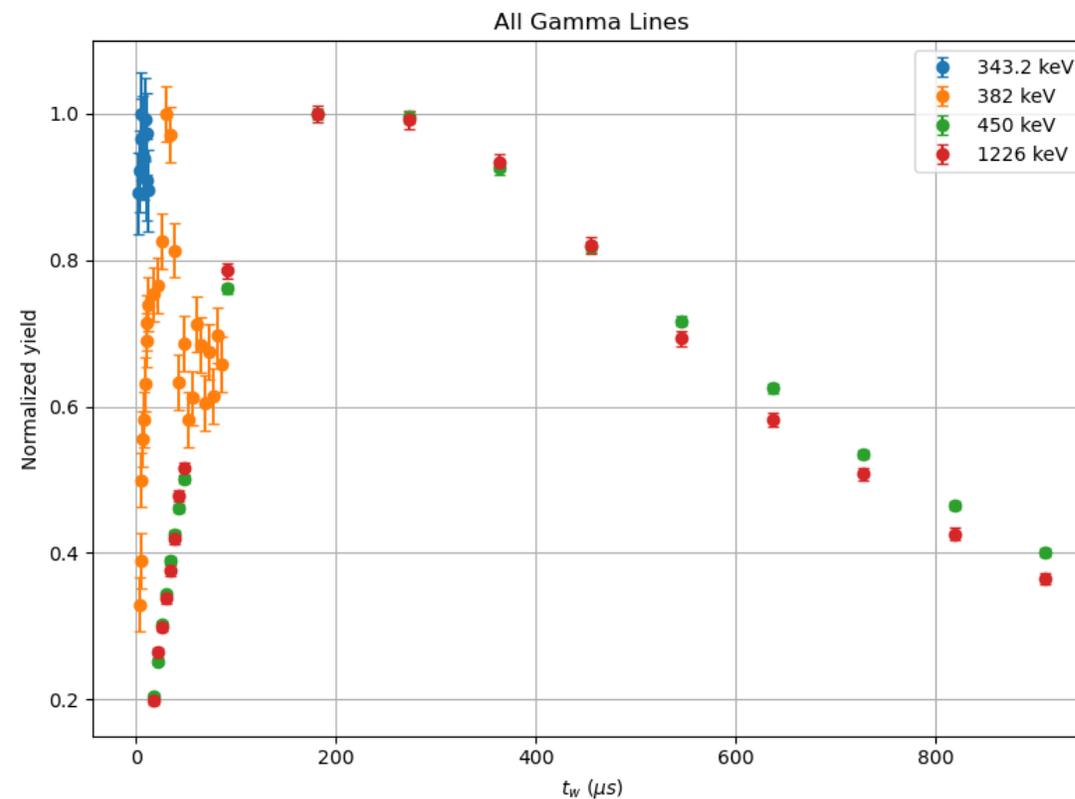
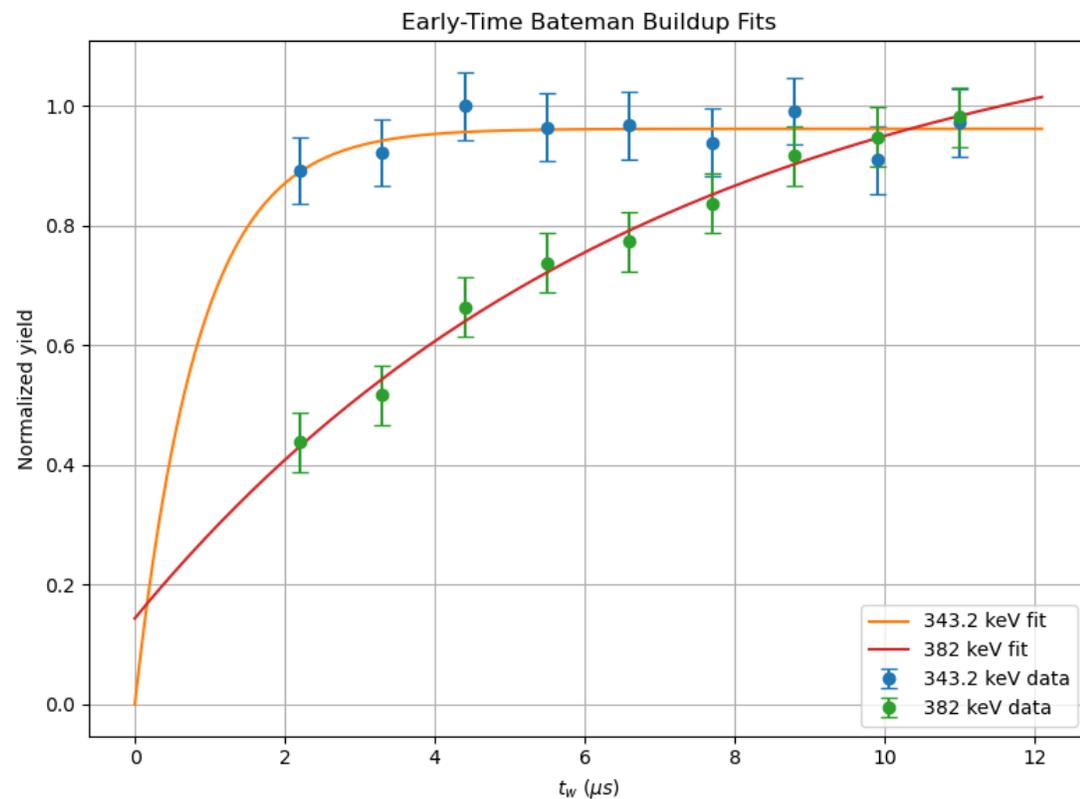
# <sup>131</sup>Sb



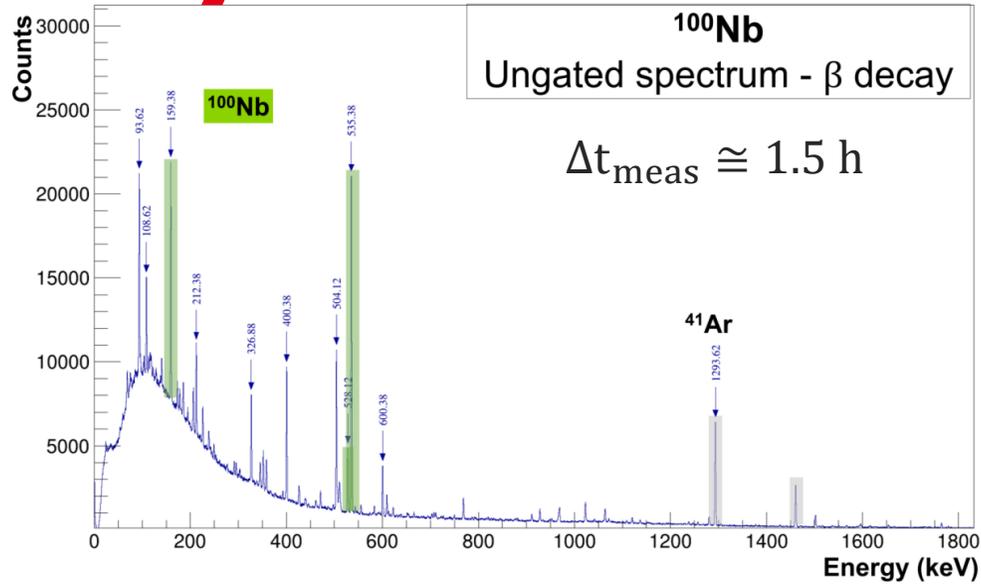
Python Script Weighted Fit:  
 $\lambda = 1.2091\text{e-}02 \pm 4.0769\text{e-}04 \text{ 1}/\mu\text{s}$   
 $t_{1/2} = \mathbf{57.33 \pm 1.93 \mu\text{s}}$   
Quality of the fit:  
 $\chi^2 / \text{ndf} = 1.60, \text{p-value} = 1.3\text{e-}01$



Python Script Weighted Fit:  
 $\lambda = 1.2000\text{e-}02 \pm 5.8706\text{e-}04 \text{ 1}/\mu\text{s}$   
 $t_{1/2} = \mathbf{57.76 \pm 2.83 \mu\text{s}}$   
Quality of the fit:  
 $\chi^2 / \text{ndf} = 4.1, \text{p-value} = 1.7\text{e-}04$

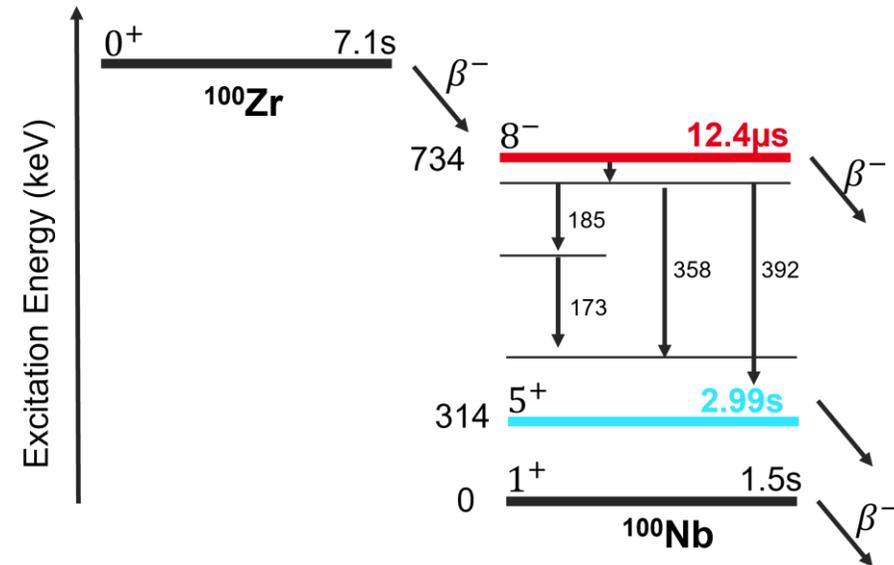


# Analysis - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$

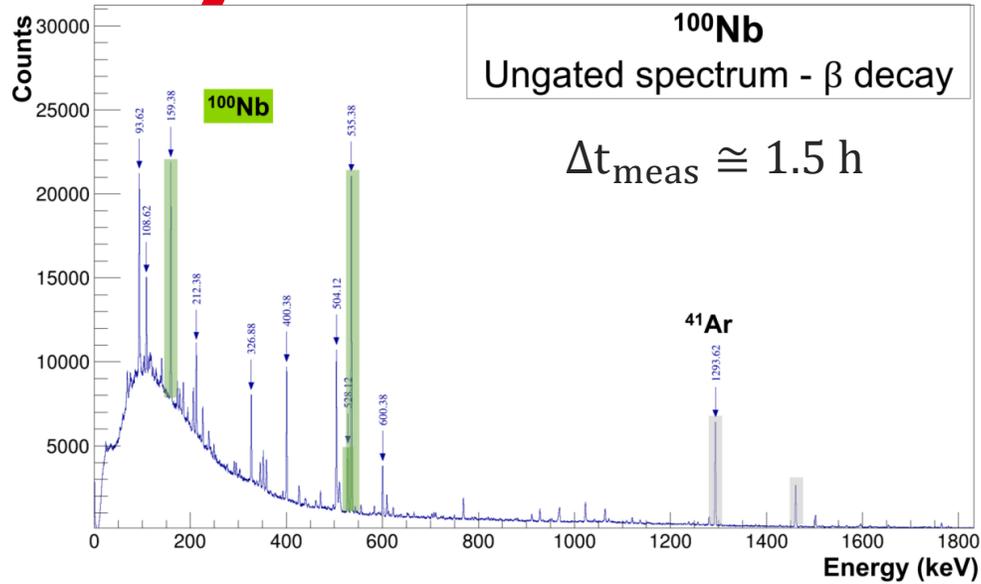


Ungated Spectrum

→ Ground state measurement →  $N_{\gamma}$

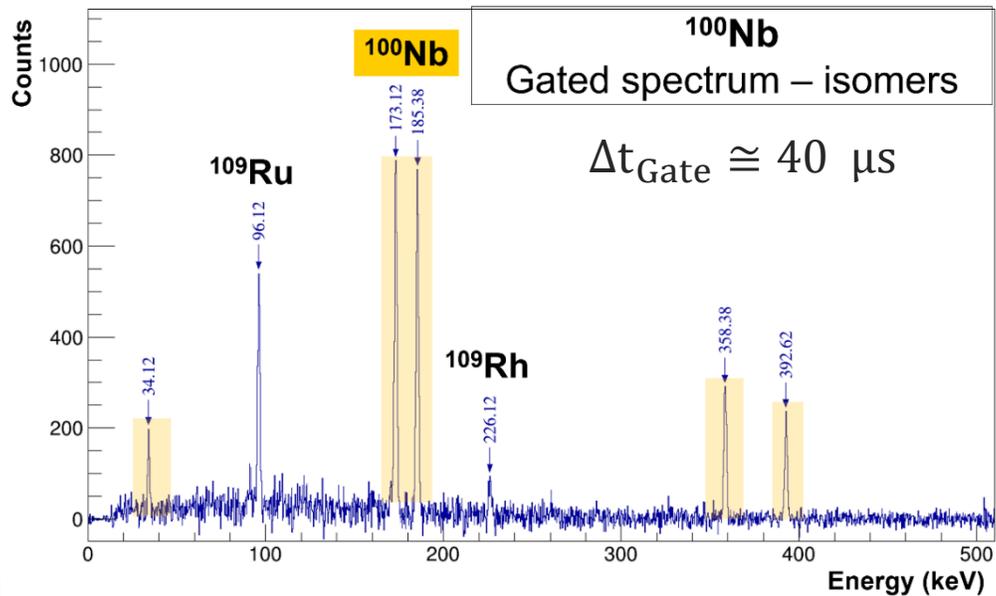


# Analysis - $^{100}\text{Nb}$ from $^{241}\text{Am}(2n_{\text{th}},f)$



Ungated Spectrum

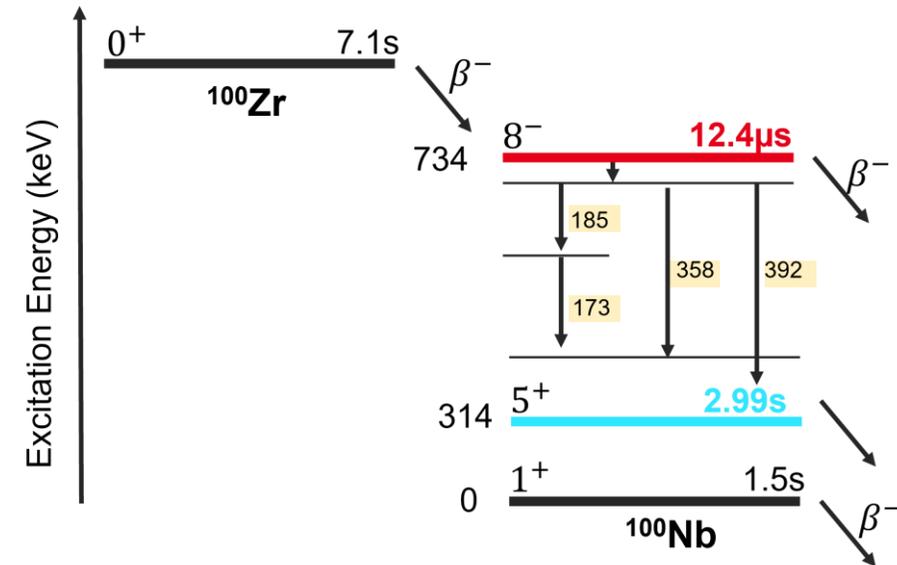
$\rightarrow \beta$ -decay states measurement  $\rightarrow N_{\gamma}, N_{\gamma}$

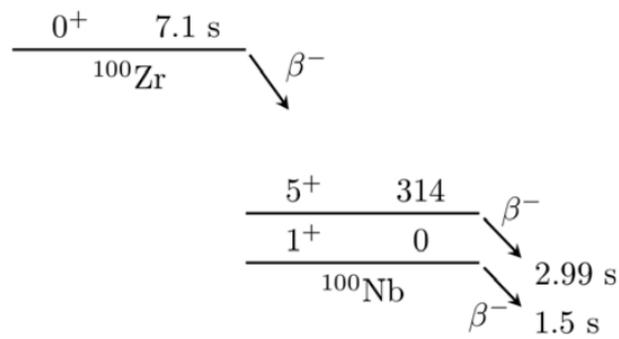
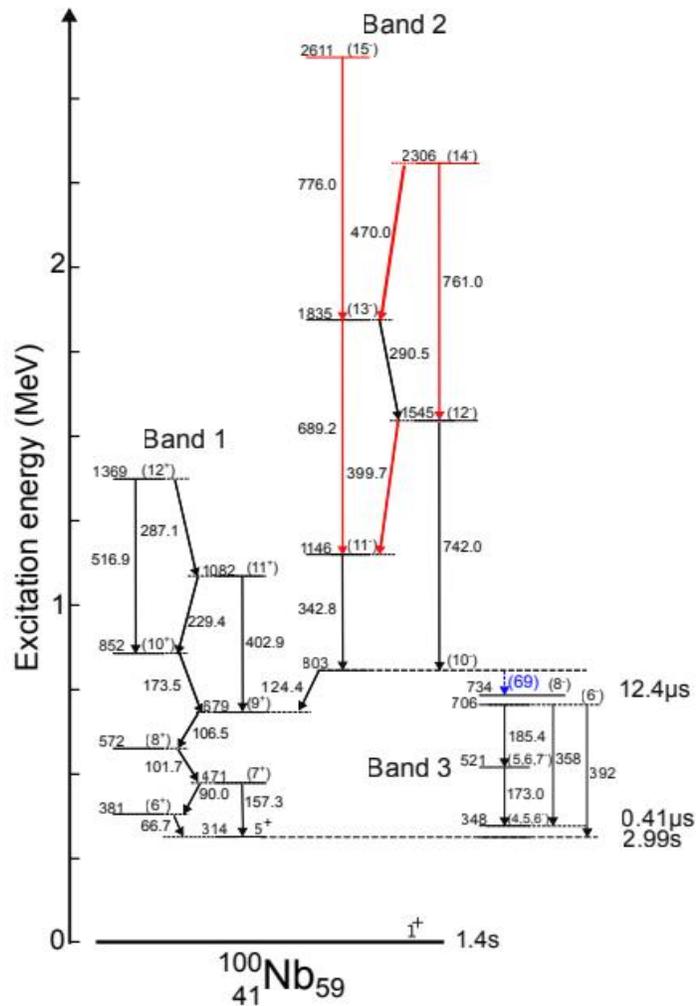


Gated Spectrum

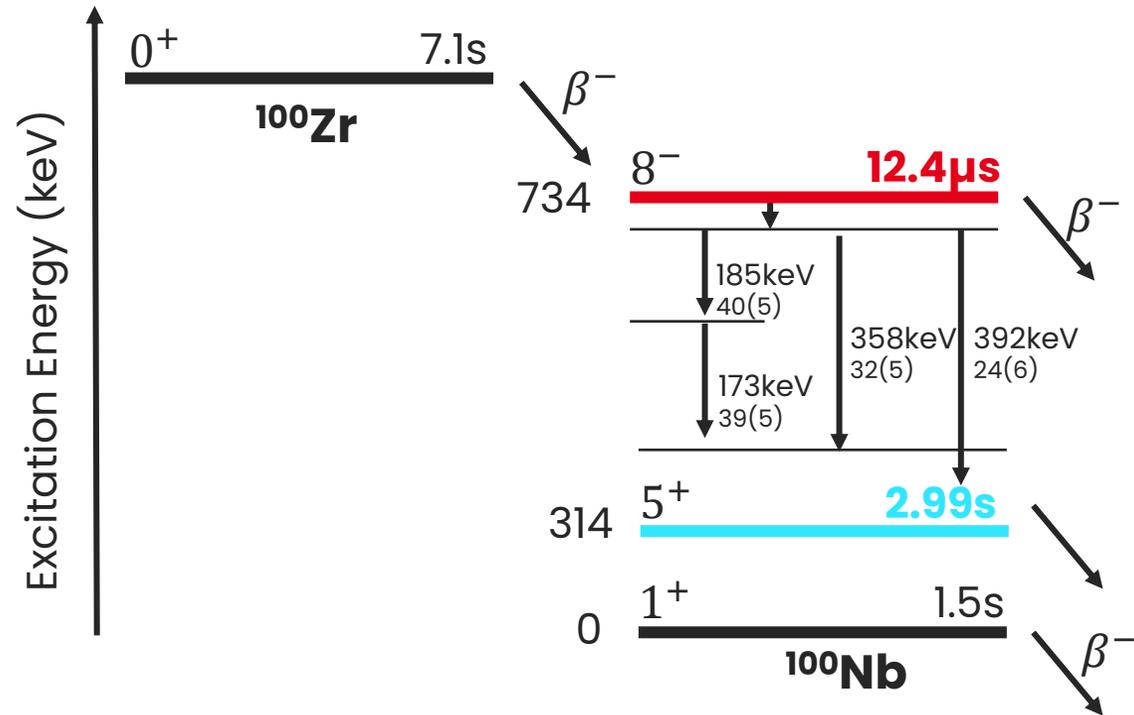
**Time coincidence**  
between ionization  
chamber and  
HPGe detectors

$\rightarrow \mu\text{s}$  Isomeric state  
measurement  $\rightarrow N_{\gamma}$

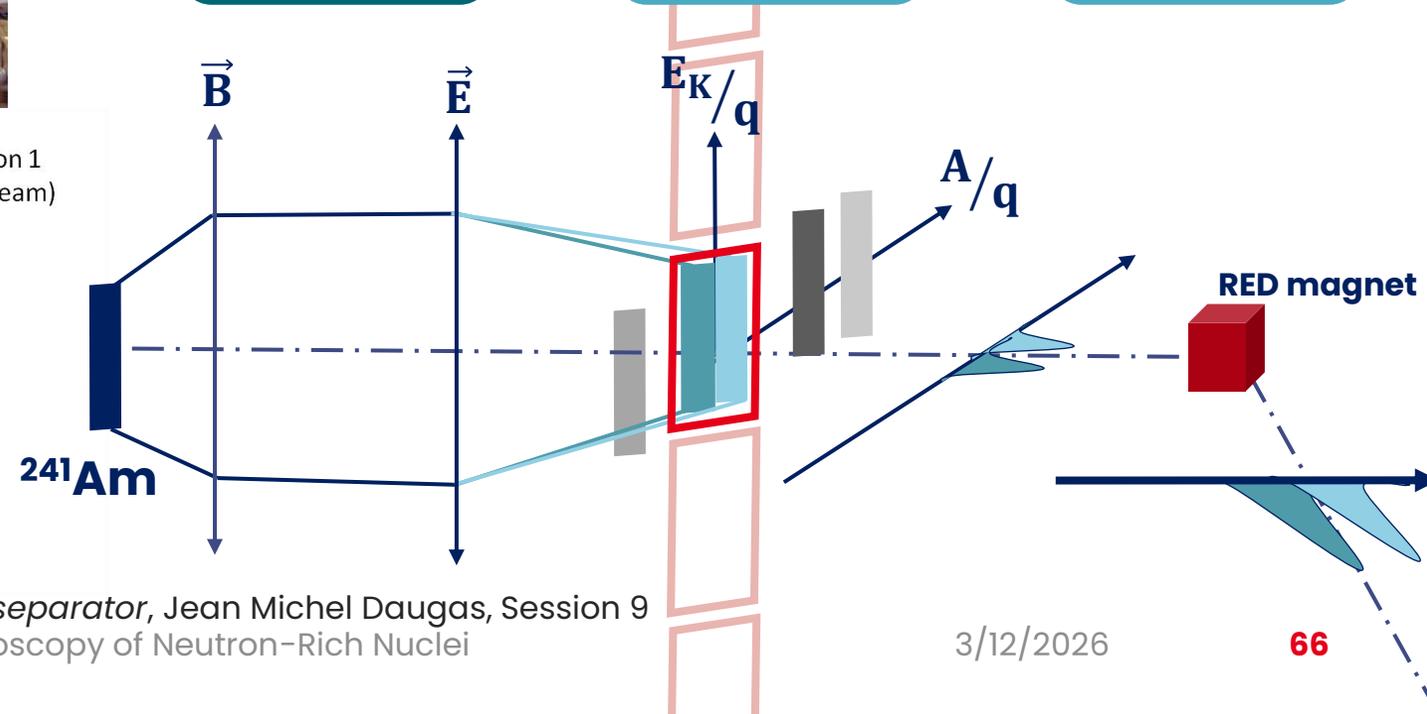
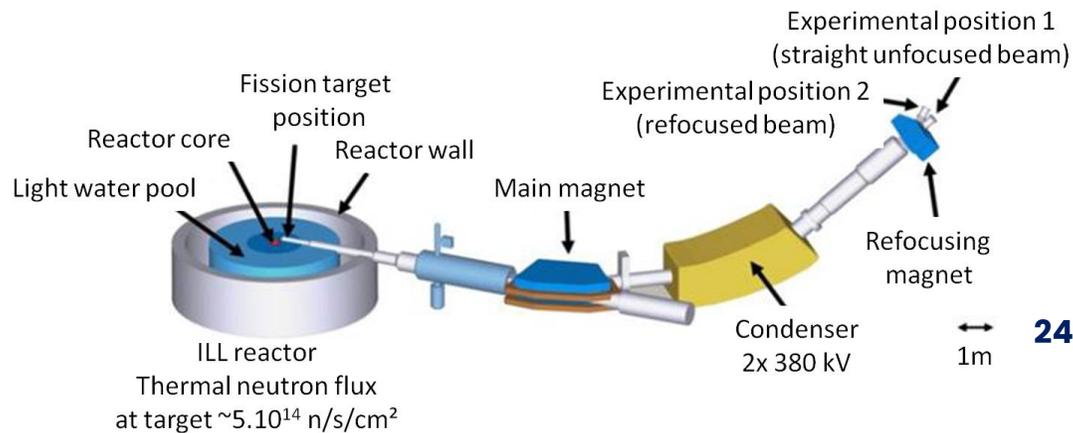
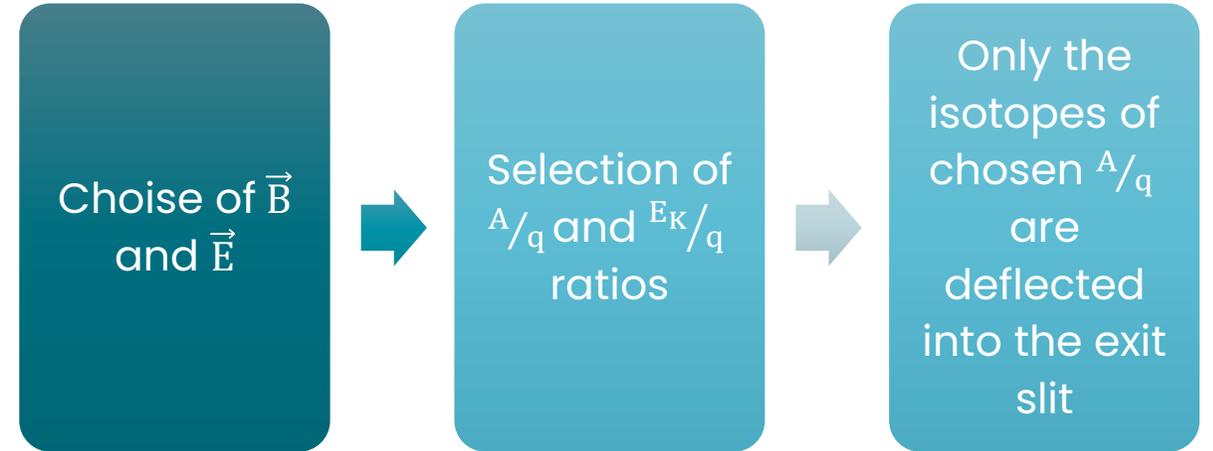
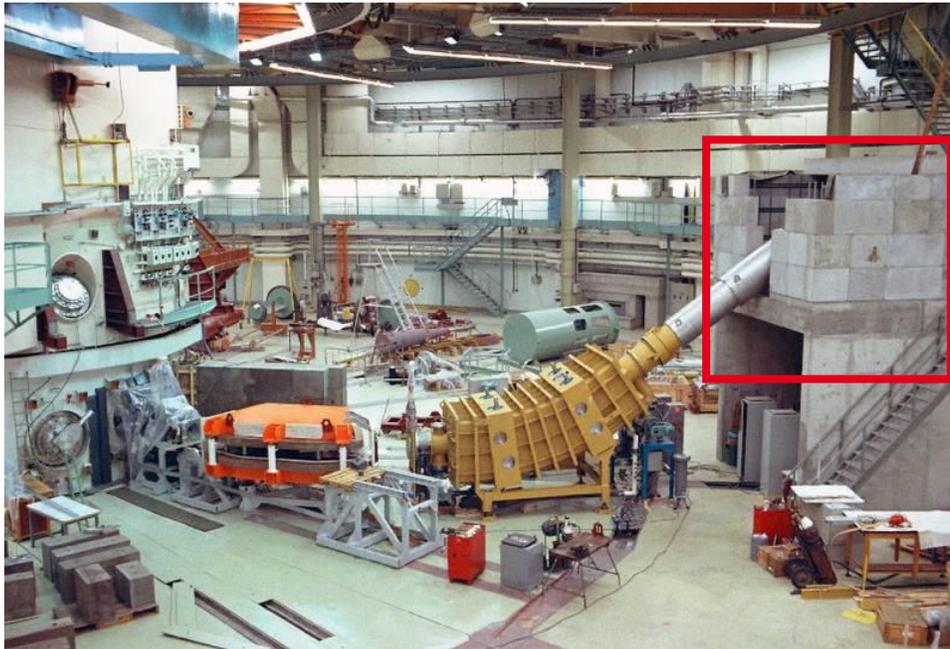




**Figure 1.** Scheme of the two isomers in  $^{100}\text{Nb}$ . Spin-parity and energy (in keV) of both isomeric states, as well as the half-life of their  $\beta$  decay are presented. The decay of the ground state of the grandparent,  $^{100}\text{Zr}$ , is also depicted.

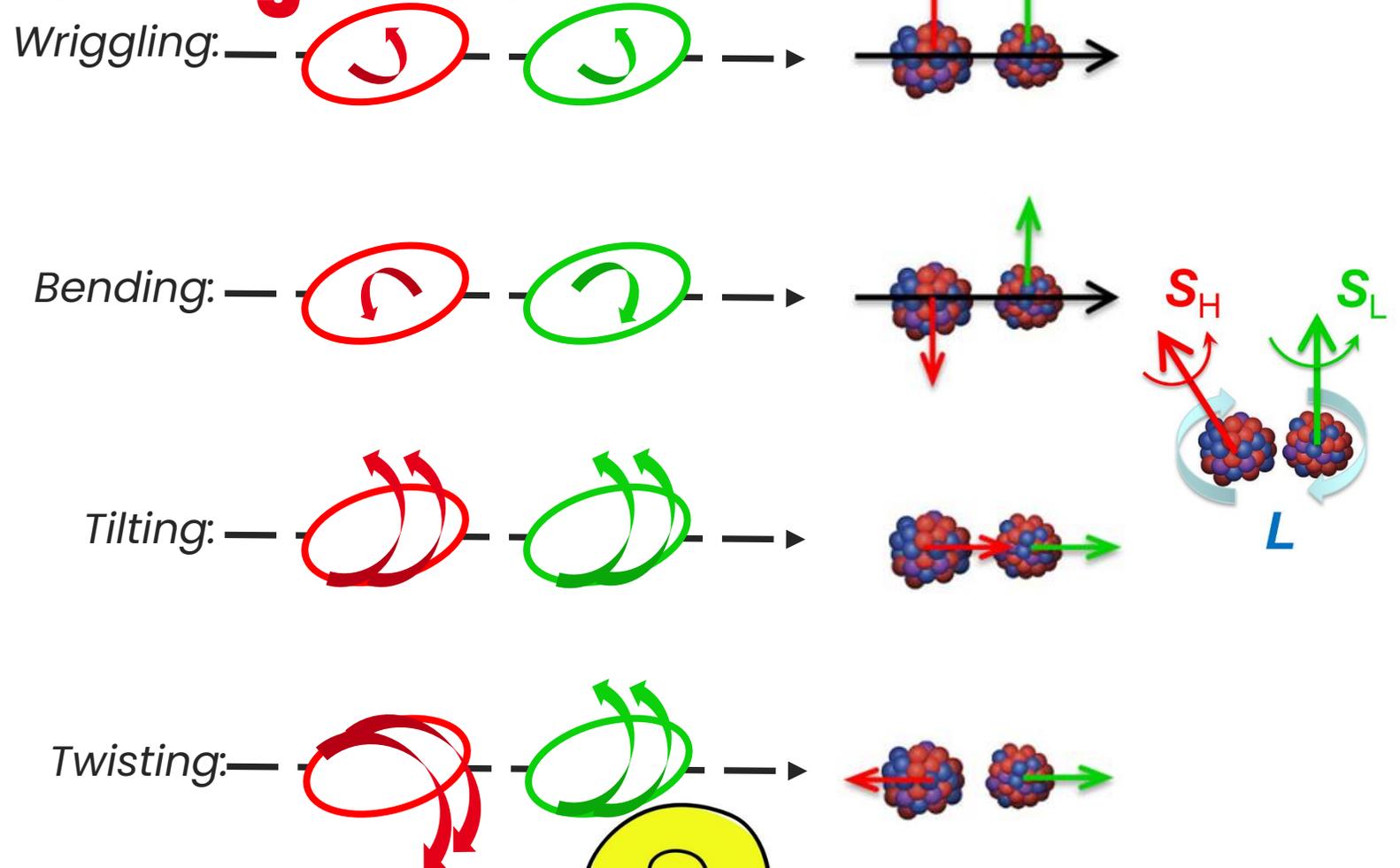
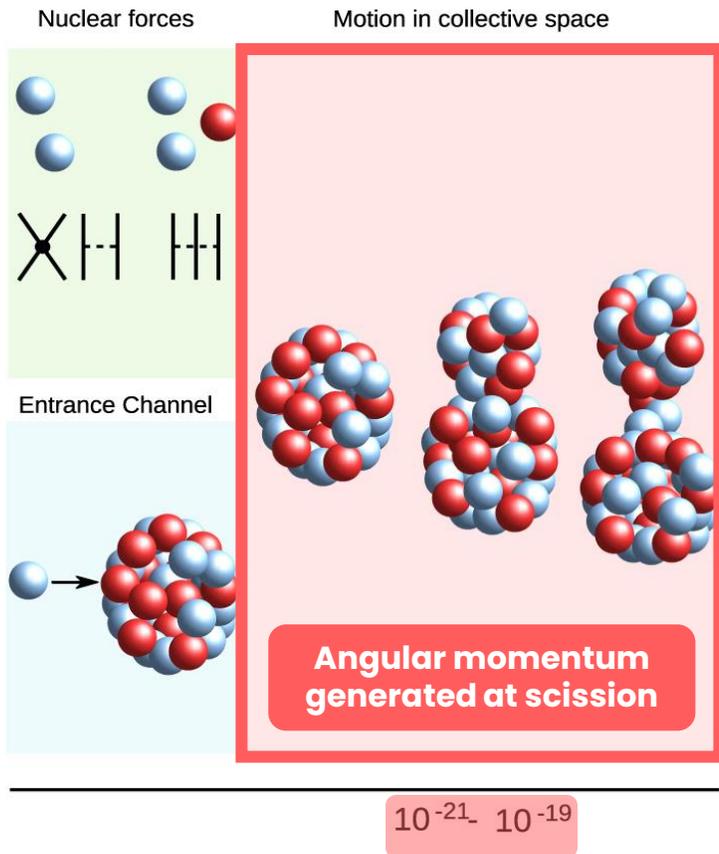


# Experimental setup of LOHENGRIN spectrometer in ILL



Study opportunities with the LOHENGRIN fission fragment separator, Jean Michel Daugas, Session 9  
 FISSION 2026: 7th Workshop on Nuclear Fission and Spectroscopy of Neutron-Rich Nuclei

# Motivation: Angular momentum generation mechanism of Fission Fragments



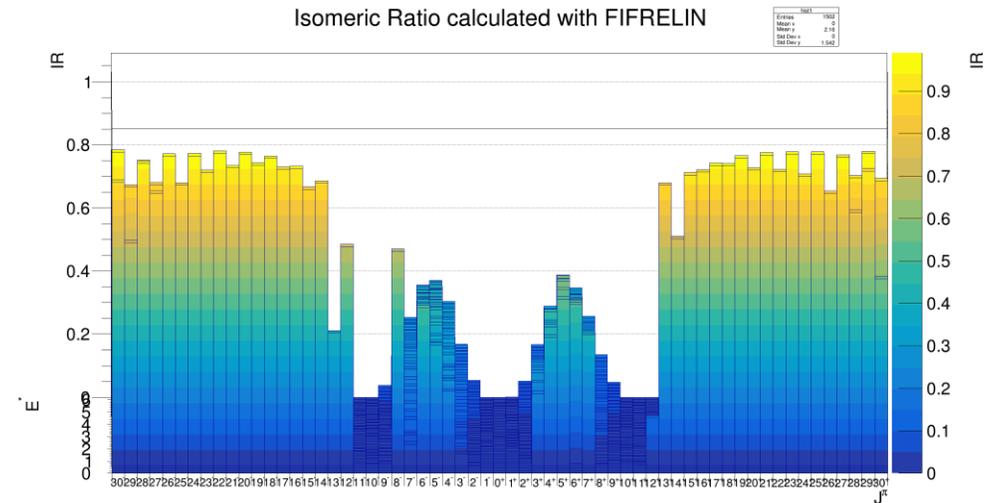
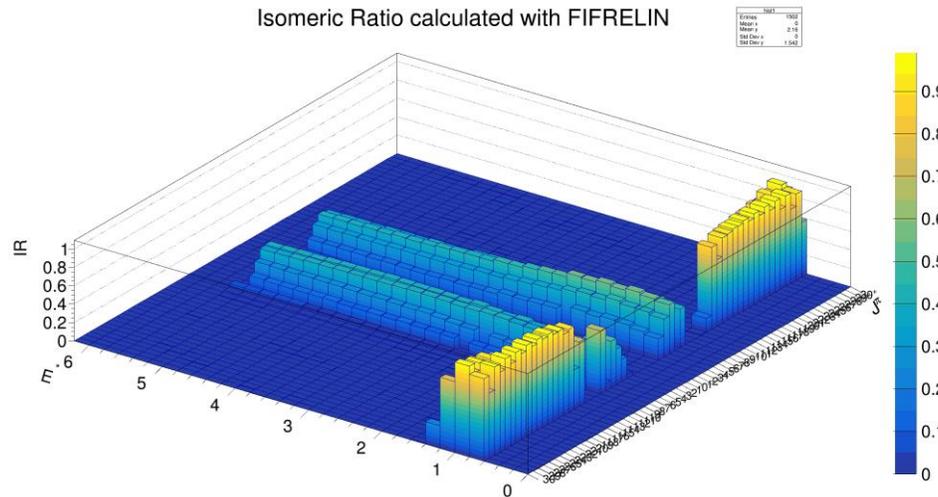
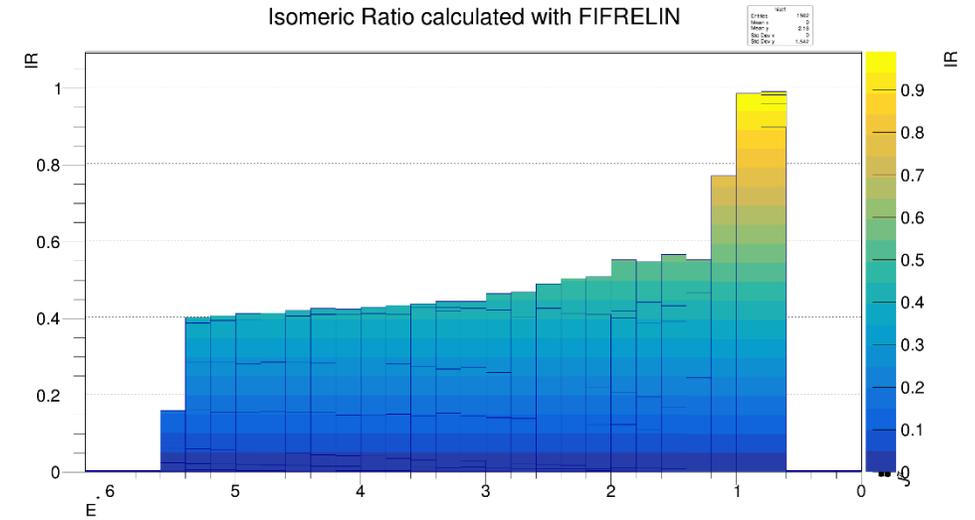
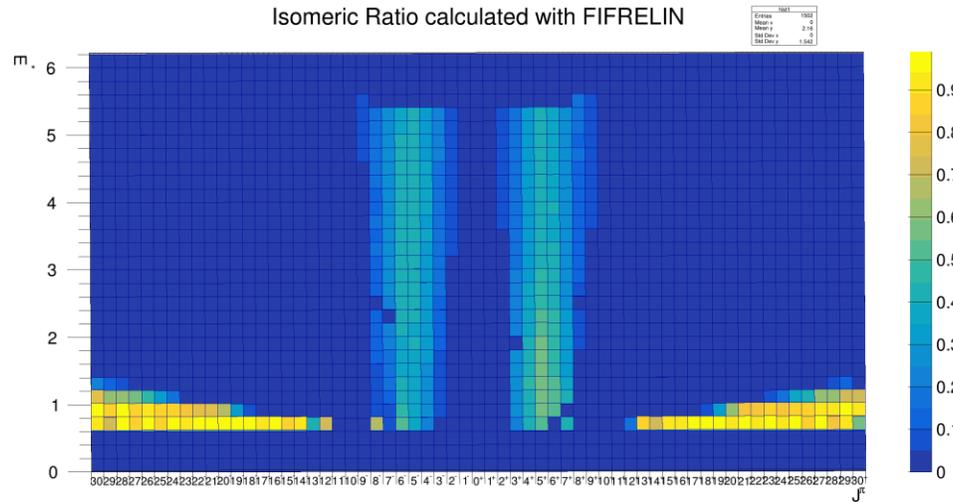
M. Bender et al., *Future of nuclear fission theory*, J. Phys. G: Nucl. Part. Phys. 47 113002 (2020)  
 J. Randrup et al., *Generation of Fragment Angular Momentum in Nuclear Fission*, EPJ WoC, 284,04004 (2023)

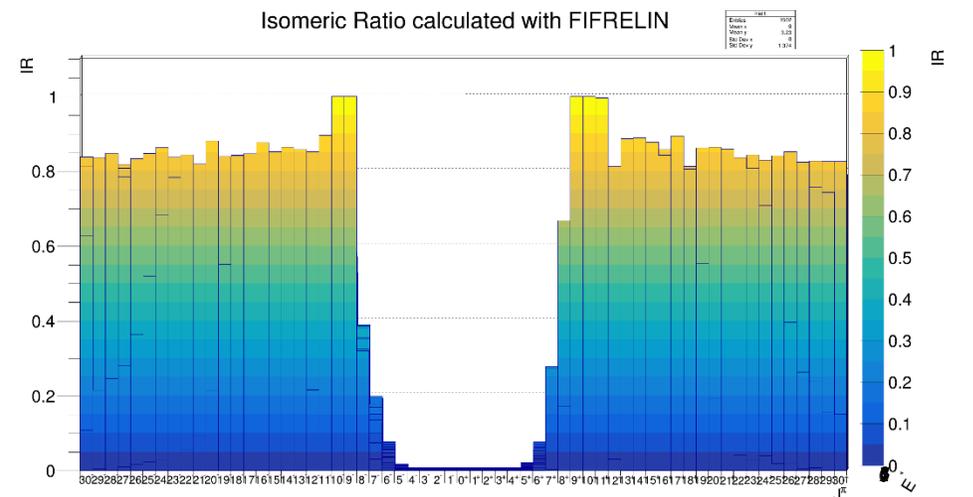
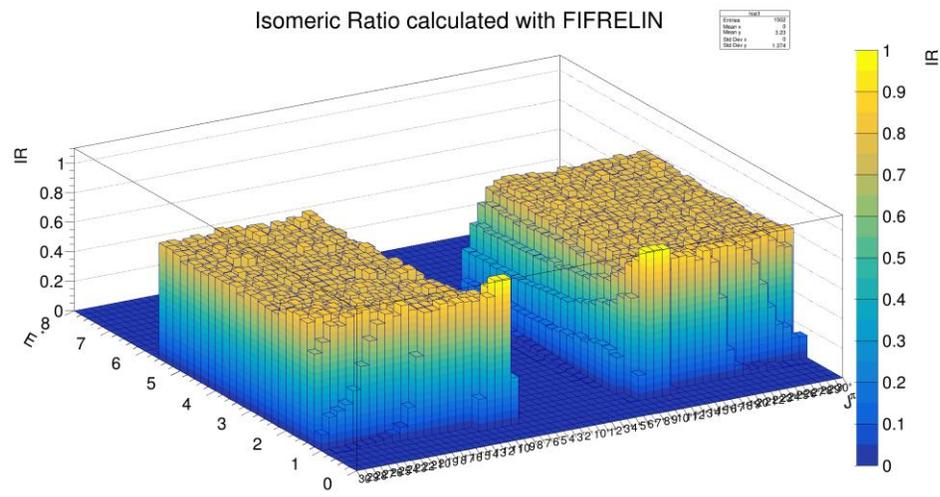
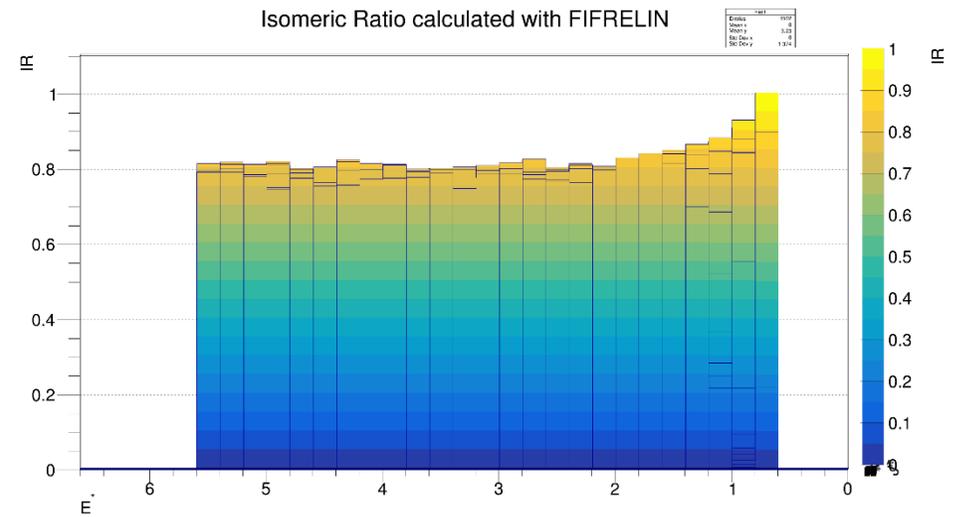
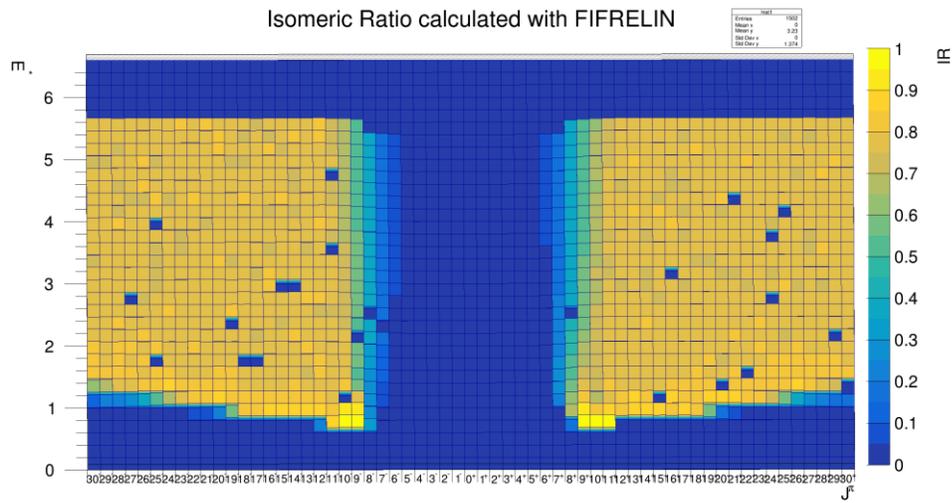
FISSION 2026: 7th Workshop on Nuclear Fission and Spectroscopy of Neutron-Rich Nuclei

3/12/2026

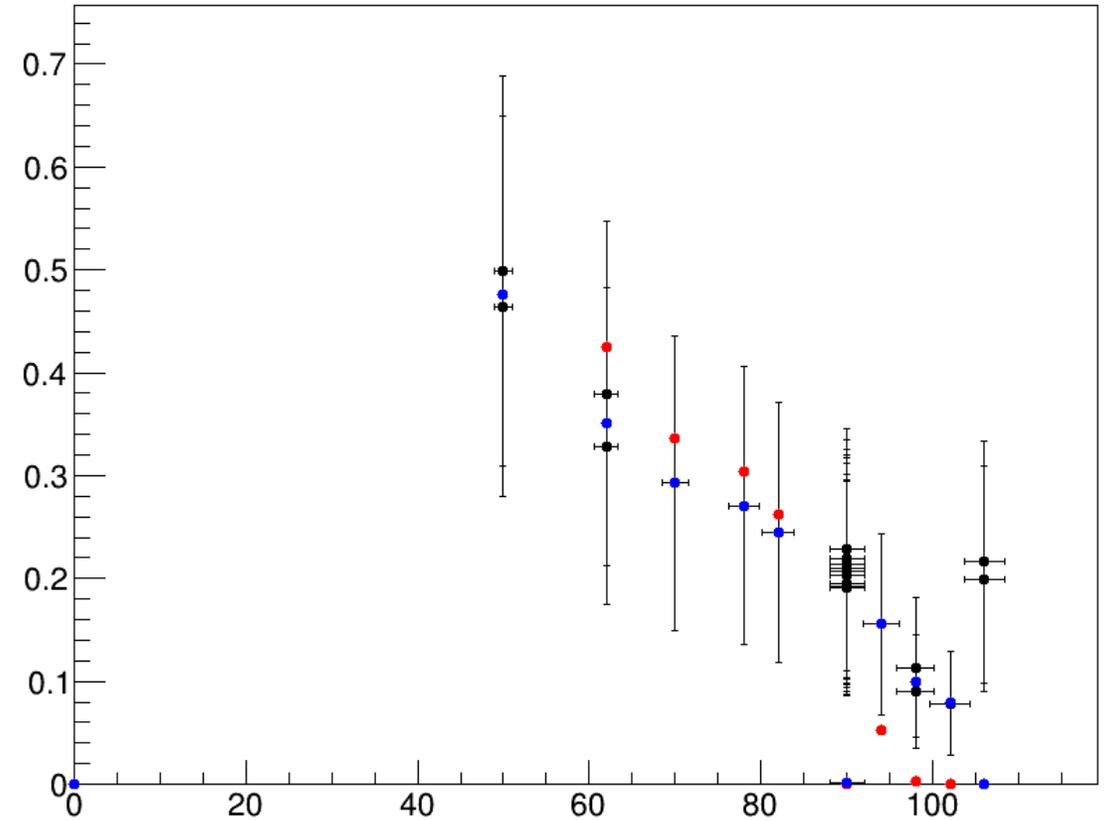
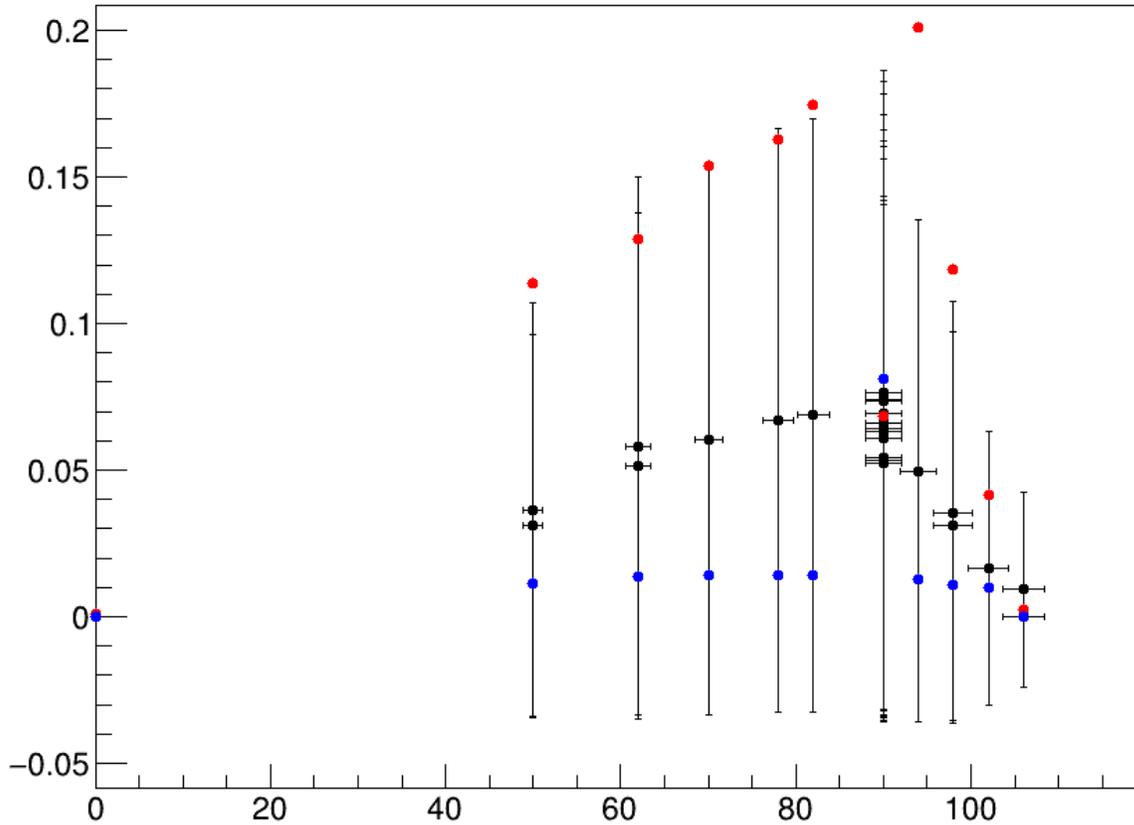
67

# Comparison with FIFRELIN

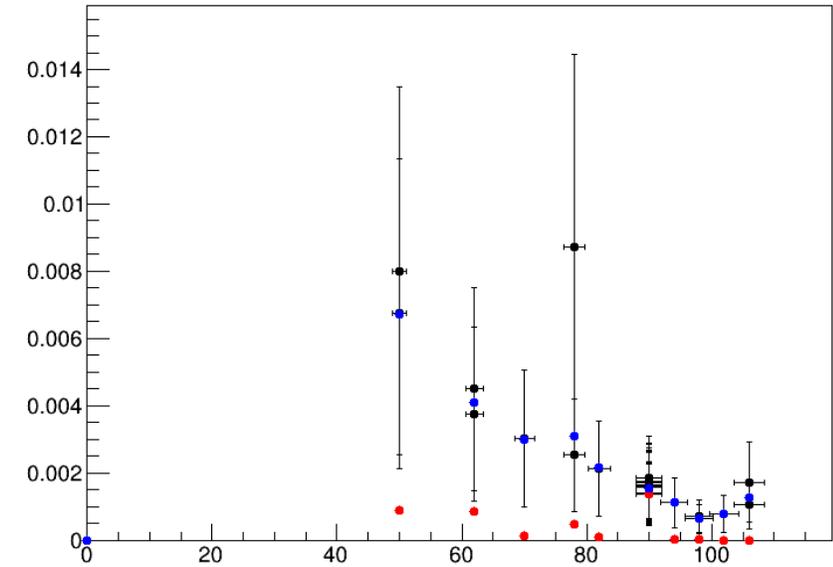
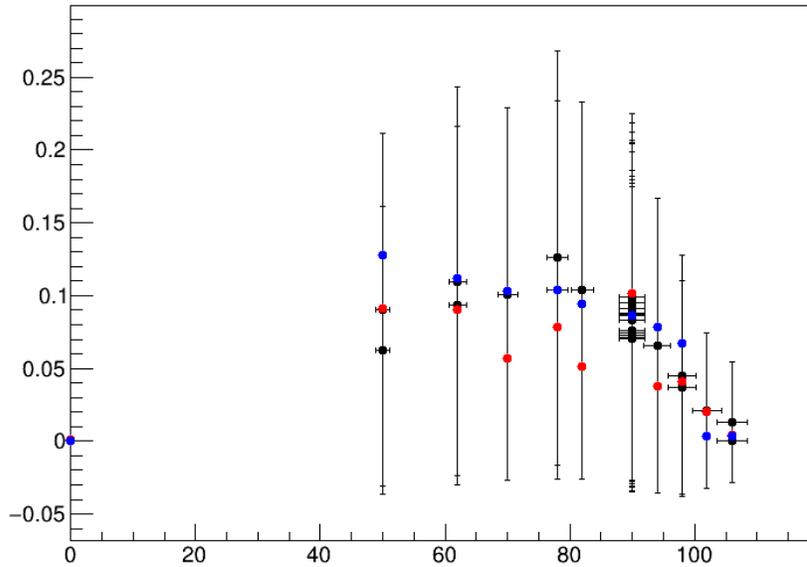




# Comparison with FIFRELIN

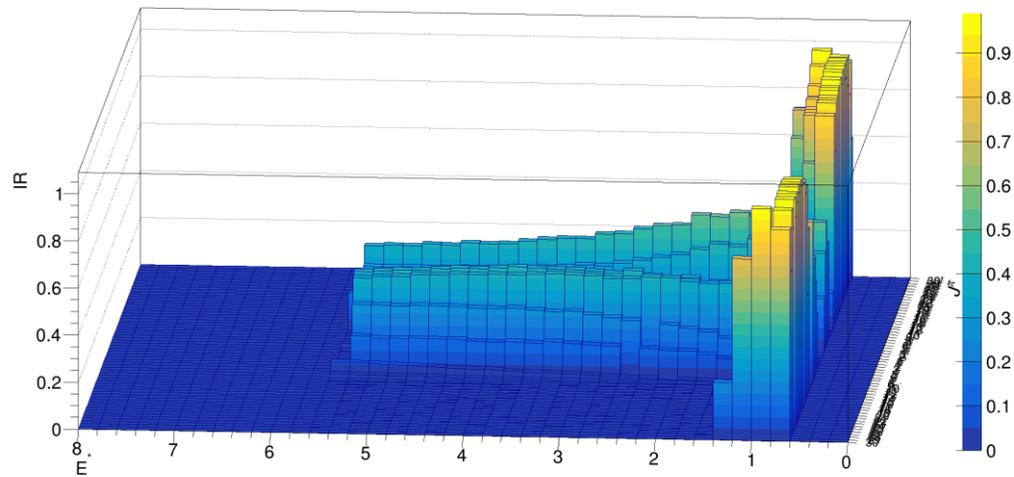


# Comparison with FIFRELIN



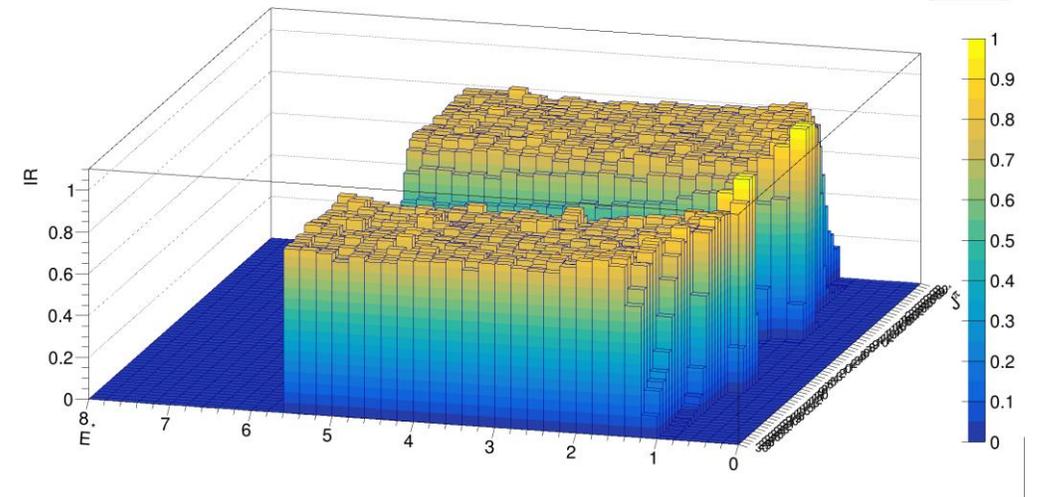
Isomeric Ratio calculated with FIFRELIN

|       |      |
|-------|------|
| Count | 100  |
| Mean  | 1.00 |
| Stdev | 1.00 |
| Min   | 0.00 |
| Max   | 2.00 |



Isomeric Ratio calculated with FIFRELIN

|       |       |
|-------|-------|
| Count | 1000  |
| Mean  | 0.0   |
| Stdev | 3.25  |
| Min   | -0.0  |
| Max   | 1.374 |



# Bayesian $J_{\text{cutoff}}$

